

ANISOTROPIC ELLIPTIC OPTICAL FIBERS

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## SUMMARY

The instability of the fabrication process of optical fibers introduces both ellipticities and stress anisotropy. These perturbations are the causes for birefringence in single-mode optical fibers which have been researched extensively. In this research, the mode propagations in optical fibers with anisotropic elliptical core have been investigated.

The exact characteristic equation for anisotropic elliptical optical fibers can be obtained for odd and even hybrid modes in terms of infinite determinants utilizing Mathieu and modified Mathieu functions. The exact characteristic equation is applicable to elliptical fibers that have any ellipticity. A simplified characteristic equation can then be obtained by applying the weakly guiding approximation such that the difference in the refractive indices of the core and the cladding is small. It has been shown that significant simplification can be achieved under this approximation.

The simplified characteristic equation is used to compute the propagation constants for the anisotropic elliptical fiber. The expression for the power carried by the fiber is also obtained and numerical results are presented. These results may be used to approximate a number of different shapes of fibers.

## 1. INTRODUCTION

The circular optical fiber is one of the most studied media for long distance communication. An optical fiber consists of a core of a dielectric material in which the refractive index is higher than the refractive index of the cladding. However, a cladded fiber is often times modeled as a dielectric rod when the cladding radius is large enough such that the guided mode fields will decay to insignificant values at the outer boundary of cladding. The theory of optical fibers of this type is well understood and has been described in detail in the previously published research[1,2].

In an effort to obtain a low-loss fiber, Monerie[3] carried out an experimental study of doubly clad fibers in which the refractive index of inner cladding was less than that of core and outer cladding. This study shows that the optimum doping levels in the core of doubly clad fibers are less than those required by dispersion-free singly clad fibers. This leads to a smaller propagation loss, since the scattering losses decreased with a decreasing dopant concentration in fibers.

It is interesting to note, however, that the instability in the fabrication process may introduce ellipticities in the optical fibers. This lack of circular symmetry is one of the causes for birefringence in single-mode optical fibers; such a birefringent fiber is also called a single-polarization single-mode fiber[4]. These birefringent fibers are important for systems utilizing such fibers as fiber optic sensors and for predicting the transmission bandwidth reduction caused by group-delay differences between orthogonally polarized modes.

The birefringence due to ellipticity has been studied experimentally[5,6] and the measured data have been compared with those

equation for an uniaxially anisotropic circular rod for hybrid modes of excitation. Analytic solutions for the circular fiber when both core and cladding consist of uniaxial material was presented by Tønning[32]. This study indicates that the cut-off frequency for the lowest-order mode is not affected by the cladding anisotropy.

When the circular cross-section of the fiber is deformed into a noncircularly symmetric profile, a single mode in a circular fiber may split into two modes with different polarizations and propagation velocities[33]. This has been experimentally verified by employing the near-field method[34] and the spectral polarization method[35]. Cozen and Dyott[36] obtained the cut-off frequency of the first higher order mode in an elliptical fiber from an approximate characteristic equation. However, the limitation of their results is described by Citerne[37] and Rengarajan[38]. The cut-off characteristic has also been obtained by solving the exact characteristic equation in terms of Mathieu functions and modified Mathieu functions[38,39] and by applying the mode-matching method[27] or the critical wavelength shift formula method[40]. However, there exists a disagreement in the interpretation of their results, especially in the region where the ratio between minor axis and major axis is small; Saad[41] presented possible reasons for these differences.

For most of the practical fibers used as optical communication lines, the simplification of the characteristic equation is possible by applying the weakly guiding approximation such that the difference in the refractive indices of the core and the cladding is small[26,33,42,43]. It is shown that the error introduced by this simplification is less than 10% even when the difference in the refractive indices is equal to two. The perturbation method



can also be applied to study the polarization effects in multi-moded fibers when the fiber is weakly guiding and/or weakly anisotropic[44]. It is also possible for multi-mode fibers to simplify the characteristic equation by applying a perturbation method based on the far-from-cut-off approximation as shown by Paul and Shevgaonkar[45] in a study of circular fiber with uniaxial anisotropy. This approximation is useful for the multi-mode propagation in optical waveguides since the lower order modes that carry most of the power could be considered in the far-from-cut-off region.

For the more general case of biaxial anisotropic waveguide, an analytic solution of the field equations is not possible even for waveguides with simple geometries. However, the fields and propagation constants can be obtained by applying the numerical techniques discussed above. The propagation constants can also be computed by using a coupled mode theory. This coupled mode approach has been applied for the study of mode propagation in rectangular guides[46] and cylindrical fibers[47].

As it has been shown through the previous discussion, the instability of the fabrication process of optical fibers introduces both ellipticities and stress anisotropy. Also, the results obtained for an elliptical optical fiber may be used to approximate a number of different shapes of fibers. It can take the shape of a circular fiber or that of a flat tape fiber depending upon the eccentricity of the elliptical fiber. Hence, it is proposed to investigate the mode propagation in elliptical optical fibers containing uniaxial anisotropic media. In this study, the fiber will be modeled as a dielectric elliptical rod, since the departure of the cladding's cross-section from circular form can be ignored in the case of large dimension of cladding radius. The exact characteristic equation for the anisotropic elliptical

fiber having any ellipticity will be obtained using the series of Mathieu and modified Mathieu functions. A simplified characteristic equation will then be obtained by applying the weakly guiding approximation and the computed results will be presented.

## 2. WAVE EQUATION IN ELLIPTICAL COORDINATES

In solving Maxwell's equations, the wave equation in the waveguide can be expressed in the orthogonal curvilinear coordinates  $(\mathcal{Y}_1, \mathcal{Y}_2, z)$  as

$$(2.1) \quad \begin{aligned} & (1/\ell_1^2)(\partial^2 E_z / \partial \ell_1^2) + (1/\ell_2^2)(\partial^2 E_z / \partial \mathcal{Y}_2^2) \\ & + (1/\ell_1 \ell_2) \{ (\partial \ell_2 / \partial \ell_1) (\partial E_z / \partial \mathcal{Y}_1) + (\partial \ell_1 / \partial \mathcal{Y}_2) (\partial E_z / \partial \mathcal{Y}_2) \} \\ & + k_1^2 E_z = 0 \end{aligned}$$

where  $k_1$  is a constant and  $\ell_1$  and  $\ell_2$  are multiplying factors depending upon the particular coordinates [48].  $\partial/\partial z$  is replaced by  $-i\beta$ . An identical equation can be obtained for  $H_z$ .

For the elliptical coordinates shown in Figure 1,

$$(2.2) \quad \mathcal{Y}_1 = \xi, \quad \mathcal{Y}_2 = \eta$$

and

$$(2.3) \quad \ell_1 = \ell_2 = q [(\cosh 2\xi - \cos 2\eta)/2]^{\frac{1}{2}}.$$

By substituting Eqs. (2.2) and (2.3) into Eq. (2.1), the following equation is obtained

$$(2.4) \quad \partial^2 E_z / \partial \xi^2 + \partial^2 E_z / \partial \eta^2 + 2k^2 (\cosh 2\xi - \cos 2\eta) E_z = 0$$

with  $2k = k_1 q$ . Then Eq. (2.4) is the two-dimensional wave equation in elliptical coordinates.

If we let the solution of Eq. (2.4) be  $E_z(\xi, \eta) = \Psi(\xi)\Phi(\eta)$ , Eq. (2.4) becomes

$$(2.5) \quad \Phi d^2 \Psi / d\xi^2 + \Psi d^2 \Phi / d\eta^2 + 2k^2 (\cosh 2\xi - \cos 2\eta) \Psi \Phi = 0$$

Dividing Eq. (2.5) by  $\Psi\Phi$  yields

$$(2.6) \quad (1/\Psi) d^2 \Psi / d\xi^2 + 2k^2 \cosh 2\xi + (1/\Phi) d^2 \Phi / d\eta^2 - 2k^2 \cos 2\eta = 0.$$

Since the equations in the parenthesis are independent to each other, we obtain

$$(2.7) \quad d^2 \Phi / d\eta^2 + (a - 2k^2 \cos 2\eta) \Phi = 0$$

and

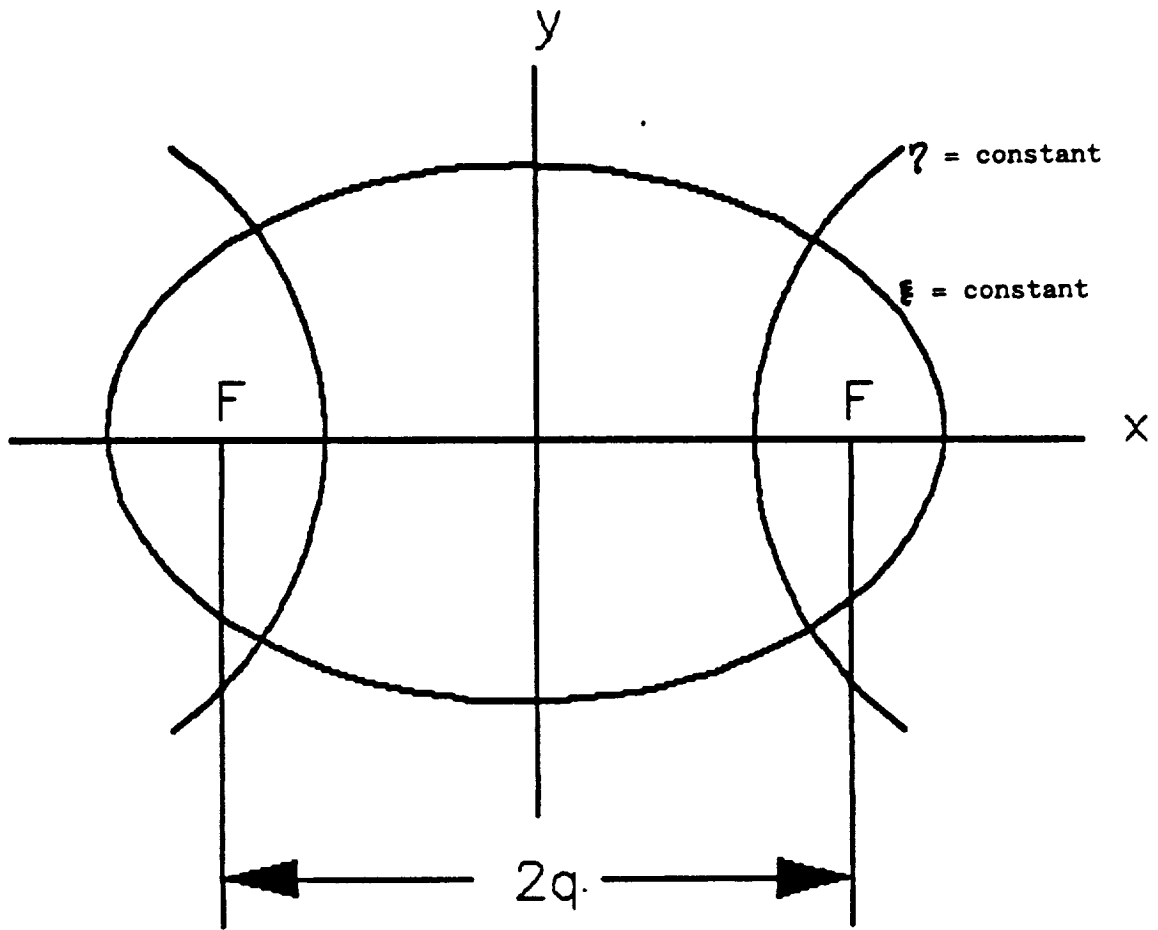


Figure 1. Elliptical coordinate system

$$(2.8) \quad d^2\psi/d\xi^2 - (a - 2k^2 \cosh 2\xi) \psi = 0$$

where  $a$  is the separation constant. Then the solutions for Eq.(2.4) are

$$(2.9) \quad E_z = \begin{cases} ce_m(\xi, k^2)ce_m(\eta, k^2) & (\text{even}) \\ se_m(\xi, k^2)se_m(\eta, k^2) & (\text{odd}) \end{cases}$$

for  $k^2 > 0$

$$(2.10) \quad E_z = \begin{cases} fek_m(\xi, k^2)ce_m(\eta, k^2) & (\text{even}) \\ gek_m(\xi, k^2)se_m(\eta, k^2) & (\text{odd}) \end{cases}$$

for  $k^2 < 0$ .

Similary, the solutions for  $H_z$  can be obtained using the method discussed above.

### 3. CHARACTERISTIC EQUATION

The geometry shown in Figure 1 consists of an uniaxially anisotropic elliptical rod with a permittivity tensor

$$(3.1) \quad \epsilon_1 = \begin{vmatrix} \epsilon_1 & 0 & 0 \\ 0 & \epsilon_1 & 0 \\ 0 & 0 & g\epsilon_1 \end{vmatrix}$$

which is embedded in a lossless dielectric medium of permittivity  $\epsilon_0$ . The anisotropic parameter  $g$  indicates the effect of anisotropic dielectric. The anisotropic parameter is unity for isotropic case.

It has been shown that in order to satisfy the boundary conditions completely both longitudinal electric and magnetic fields must be present, thus only hybrid type modes exist in elliptical fibers[2]. Furthermore, due to the asymmetry of the elliptical cylinder, two types of modes exist and they are designated as an even type modes and an odd type modes.

#### 3.1 EVEN MODES

Assuming the  $t$ - $z$  dependence of  $e^{i(\omega t - \beta z)}$  for all field components, where  $\beta$  is the propagation constant and  $\omega$  is the angular frequency, the axial components of the field for even modes are

$$(3.2) \quad \begin{aligned} E_{z1} &= \sum_{m=1}^{\infty} A_{1m} \text{se}_m(\xi, \gamma_{1e}^2) \text{se}_m(\eta, \gamma_{1e}^2) \\ H_{z1} &= \sum_{m=0}^{\infty} B_{1m} \text{ce}_m(\xi, \gamma_{1h}^2) \text{ce}_m(\eta, \gamma_{1h}^2) \end{aligned}$$

for  $0 \leq \xi \leq \xi_0$

$$(3.3) \quad \begin{aligned} E_{z2} &= \sum_{m=1}^{\infty} A_{2m} \text{gek}_m(\xi, \gamma_2^2) \text{se}_m(\eta, \gamma_2^2) \\ H_{z2} &= \sum_{m=0}^{\infty} B_{2m} \text{fek}_m(\xi, \gamma_2^2) \text{ce}_m(\eta, \gamma_2^2) \end{aligned}$$

for  $\xi_0 \leq \xi < \infty$

where  $A_{1m}$  and  $B_{1m}$ ,  $i = 1, 2$  are arbitrary constants, and

$$\begin{aligned} \gamma_{1e}^2 &= q^2/4(v^2\mu\epsilon_1 - \beta^2) \\ (3.4) \quad \gamma_{1h}^2 &= q^2/4(v^2\mu\epsilon_1 - \beta^2) \\ \gamma_2^2 &= q^2/4(v^2\mu\epsilon_0 - \beta^2) \end{aligned}$$

$q$  is the semifocal length of the ellipse and  $\mu$  is the permeability.

The transverse field components are

for  $0 \leq \xi \leq \xi_0$

$$\begin{aligned} (3.5) \quad E_{\xi 1} &= -1/(v^2\mu\epsilon_1 - \beta^2)L \\ &\quad \left( \beta \sum_{m=1}^{\infty} A_{1m} \text{se}_m'(\xi, \gamma_{1e}^2) \text{se}_m(\eta, \gamma_{1e}^2) \right. \\ &\quad \left. + v\mu \sum_{m=0}^{\infty} B_{1m} \text{ce}_m(\xi, \gamma_{1h}^2) \text{ce}_m'(\eta, \gamma_{1h}^2) \right) \end{aligned}$$

$$\begin{aligned} (3.6) \quad E_{\eta 1} &= -1/(v^2\mu\epsilon_1 - \beta^2)L \\ &\quad \left( \beta \sum_{m=1}^{\infty} A_{1m} \text{se}_m(\xi, \gamma_{1e}^2) \text{se}_m'(\eta, \gamma_{1e}^2) \right. \\ &\quad \left. - v\mu \sum_{m=0}^{\infty} B_{1m} \text{ce}_m'(\xi, \gamma_{1h}^2) \text{ce}_m(\eta, \gamma_{1h}^2) \right) \end{aligned}$$

$$\begin{aligned} (3.7) \quad H_{\xi 1} &= -1/(v^2\mu\epsilon_1 - \beta^2)L \\ &\quad \left( -v\epsilon_1 \sum_{m=1}^{\infty} A_{1m} \text{se}_m(\xi, \gamma_{1e}^2) \text{se}_m'(\eta, \gamma_{1e}^2) \right. \\ &\quad \left. + \beta \sum_{m=0}^{\infty} B_{1m} \text{ce}_m'(\xi, \gamma_{1h}^2) \text{ce}_m(\eta, \gamma_{1h}^2) \right) \end{aligned}$$

$$\begin{aligned} (3.8) \quad H_{\eta 1} &= -1/(v^2\mu\epsilon_1 - \beta^2)L \\ &\quad \left( v\epsilon_1 \sum_{m=1}^{\infty} A_{1m} \text{se}_m'(\xi, \gamma_{1e}^2) \text{se}_m(\eta, \gamma_{1e}^2) \right. \\ &\quad \left. + \beta \sum_{m=0}^{\infty} B_{1m} \text{ce}_m(\xi, \gamma_{1h}^2) \text{ce}_m'(\eta, \gamma_{1h}^2) \right) \end{aligned}$$

for  $\xi_0 \leq \xi < \infty$

$$\begin{aligned} (3.9) \quad E_{\xi 2} &= -1/(v^2\mu\epsilon_0 - \beta^2)L \\ &\quad \left( \beta \sum_{m=1}^{\infty} A_{2m} \text{Gek}_m'(\xi, \gamma_2^2) \text{se}_m(\eta, \gamma_2^2) \right. \\ &\quad \left. + v\mu \sum_{m=0}^{\infty} B_{2m} \text{Fek}_m(\xi, \gamma_2^2) \text{ce}_m'(\eta, \gamma_2^2) \right) \end{aligned}$$

$$\begin{aligned} (3.10) \quad E_{\eta 2} &= -1/(v^2\mu\epsilon_0 - \beta^2)L \\ &\quad \left( \beta \sum_{m=1}^{\infty} A_{2m} \text{Gek}_m(\xi, \gamma_2^2) \text{se}_m'(\eta, \gamma_2^2) \right. \\ &\quad \left. - v\mu \sum_{m=0}^{\infty} B_{2m} \text{Fek}_m'(\xi, \gamma_2^2) \text{ce}_m(\eta, \gamma_2^2) \right) \end{aligned}$$

$$(3.11) \quad H_{\xi 2} = -1/(\nu^2 \mu \epsilon_0 - \beta^2) L \left( -\nu \epsilon_0 \sum_{m=1}^{\infty} A_{2m} \text{Gek}_m(\xi, \gamma_2^2) \text{se}_m'(\eta, \gamma_2^2) + \beta \sum_{m=0}^{\infty} B_{2m} \text{Fek}_m'(\xi, \gamma_2^2) \text{ce}_m(\eta, \gamma_2^2) \right)$$

$$(3.12) \quad H_{\eta 2} = -1/(\nu^2 \mu \epsilon_0 - \beta^2) L \left( \nu \epsilon_0 \sum_{m=1}^{\infty} A_{2m} \text{Gek}_m'(\xi, \gamma_2^2) \text{se}_m(\eta, \gamma_2^2) + \beta \sum_{m=0}^{\infty} B_{2m} \text{Fek}_m(\xi, \gamma_2^2) \text{ce}_m'(\eta, \gamma_2^2) \right)$$

where

$$L = q((\cosh 2\xi - \cos 2\eta)/2)^{1/2}$$

and the derivative with respect to  $\xi$  or  $\eta$  is denoted by the prime.

The boundary conditions require that the tangential  $E$  and  $H$  fields be continuous at the dielectric discontinuities. Equating the tangential fields at the boundary surface,  $\xi = \xi_0$ , gives

$$(3.13) \quad \sum_{m=1}^{\infty} A_{1m} \text{se}_m(\xi_0) \text{se}_m(\eta) = \sum_{m=1}^{\infty} A_{2m} \text{Gek}_m(\xi_0) \text{se}_m^*(\eta)$$

$$(3.14) \quad \sum_{m=0}^{\infty} B_{1m} \text{ce}_m(\xi_0) \text{ce}_m(\eta) = \sum_{m=0}^{\infty} B_{2m} \text{Fek}_m(\xi_0) \text{ce}_m^*(\eta)$$

$$(3.15) \quad \begin{aligned} 1/(\nu^2 \mu \epsilon_1 - \beta^2) \left( \beta \sum_{m=1}^{\infty} A_{1m} \text{se}_m(\xi_0) \text{se}_m'(\eta) - \nu \mu \sum_{m=0}^{\infty} B_{1m} \text{ce}_m'(\xi_0) \text{ce}_m(\eta) \right) \\ = 1/(\nu^2 \mu \epsilon_0 - \beta^2) \left( \beta \sum_{m=1}^{\infty} A_{2m} \text{Gek}_m(\xi_0) \text{se}_m^{*'}(\eta) - \nu \mu \sum_{m=0}^{\infty} B_{2m} \text{Fek}_m'(\xi_0) \text{ce}_m^{*'}(\eta) \right) \end{aligned}$$

$$(3.16) \quad \begin{aligned} 1/(\nu^2 \mu \epsilon_1 - \beta^2) \left( \nu \epsilon_1 \sum_{m=1}^{\infty} A_{1m} \text{se}_m'(\xi_0) \text{se}_m(\eta) + \beta \sum_{m=0}^{\infty} B_{1m} \text{ce}_m(\xi_0) \text{ce}_m'(\eta) \right) \\ = 1/(\nu^2 \mu \epsilon_0 - \beta^2) \left( \nu \epsilon_0 \sum_{m=1}^{\infty} A_{2m} \text{Gek}_m'(\xi_0) \text{se}_m^*(\eta) + \beta \sum_{m=0}^{\infty} B_{2m} \text{Fek}_m(\xi_0) \text{ce}_m^{*'}(\eta) \right) \end{aligned}$$

where the following abbreviations have been used,

$$(3.17) \quad \text{se}_m(\xi_0) = \text{se}_m(\xi_0, \gamma_{1e}^2)$$

$$(3.18) \quad \text{se}_m(\eta) = \text{se}_m(\eta, \gamma_{1e}^2)$$

$$(3.19) \quad \text{ce}_m(\xi_0) = \text{ce}_m(\xi_0, \gamma_{1h}^2)$$

$$(3.20) \quad \text{ce}_m(\eta) = \text{ce}_m(\eta, \gamma_{1h}^2)$$



$$(3.21) \quad \text{Gek}_m(\xi_0) = \text{Gek}_m(\xi_0, \gamma_2^2)$$

$$(3.22) \quad \text{se}_m^*(\eta) = \text{se}_m(\eta, \gamma_2^2)$$

$$(3.23) \quad \text{Fek}_m(\xi_0) = \text{Fek}_m(\xi_0, \gamma_2^2)$$

$$(3.24) \quad \text{ce}_m^*(\eta) = \text{ce}_m(\eta, \gamma_2^2).$$

Multiplying both sides of Eqs. (3.13) and (3.16) by  $\text{se}_n(\eta)$  and Eqs. (3.14) and (3.15) by  $\text{ce}_n(\eta)$ , integrating with respect to  $\eta$  from 0 to  $2\pi$ , and applying the orthogonality relations of the angular Mathieu functions

$$(3.25) \quad \int_0^{2\pi} \text{ce}_m \text{ce}_n d\eta = 0 \quad \text{if } m \neq n$$

leads to

$$(3.26) \quad A_{1n} \text{se}_n = \sum_{m=1}^{\infty} A_{2m} \text{Gek}_m \beta_{m,n}$$

$$(3.27) \quad B_{1n} \text{ce}_n = \sum_{m=0}^{\infty} B_{2m} \text{Fek}_m \alpha_{m,n}$$

$$(3.28) \quad \beta \sum_{m=1}^{\infty} A_{1m} \text{se}_m \nu_{n,m} - \nu \mu B_{1n} \text{ce}_n =$$

$$\beta \gamma_{1h}^2 / \gamma_2^2 \sum_{m=1}^{\infty} A_{2m} \text{Gek}_m \nu_{n,m} - \nu \mu \gamma_{1h}^2 / \gamma_2^2 \sum_{m=0}^{\infty} B_{2m} \text{Fek}_m \alpha_{n,m}$$

$$(3.29) \quad \nu \epsilon_1 A_{1n} \text{se}_n + \beta \sum_{m=0}^{\infty} B_{1m} \text{ce}_m \psi_{n,m} =$$

$$\nu \epsilon_0 \gamma_{1h}^2 / \gamma_2^2 \sum_{m=1}^{\infty} A_{2m} \text{Gek}_m \beta_{n,m} + \beta \gamma_{1h}^2 / \gamma_2^2 \sum_{m=0}^{\infty} B_{2m} \text{Fek}_m \psi_{n,m}$$

The prime over the summation sign is used to indicate that either odd or even values of  $m$  are used accordingly as to whether  $n$  is odd or even.

$\alpha_{m,n}$ ,  $\beta_{m,n}$ ,  $\psi_{m,n}$  and  $\nu_{m,n}$  are given by the following

$$(3.30) \quad \alpha_{m,n} = \int_0^{2\pi} \text{ce}_m^*(\eta) \text{ce}_n(\eta) d\eta / \int_0^{2\pi} \text{ce}_n^2(\eta) d\eta$$

$$(3.31) \quad \beta_{m,n} = \int_0^{2\pi} \text{se}_m^*(\eta) \text{se}_n(\eta) d\eta / \int_0^{2\pi} \text{se}_n^2(\eta) d\eta$$

$$(3.32) \quad \psi_{m,n} = \int_0^{2\pi} \text{ce}_m^*(\eta) \text{se}_n(\eta) d\eta / \int_0^{2\pi} \text{se}_n^2(\eta) d\eta$$

$$(3.33) \quad \nu_{m,n} = \int_0^{2\pi} \text{se}_m^*(\eta) \text{ce}_n(\eta) d\eta / \int_0^{2\pi} \text{ce}_n^2(\eta) d\eta.$$

Making use of Eqs. (3.26) and (3.27), Eqs. (3.28) and (3.29) yields two sets of infinite homogeneous equations

$$(3.34) \quad \sum_{m=1}^{\infty} A_{2m} s_{m,n} + \sum_{m=0}^{\infty} B_{2m} t_{m,n} = 0$$

$$\sum_{m=1}^{\infty} A_{2m} g_{m,n} + \sum_{m=0}^{\infty} B_{2m} h_{m,n} = 0$$

where

$$(3.35) \quad g_{m,n} = -(1 - \gamma_{1h}^2/\gamma_2^2) \text{Gek}_m(\epsilon_0) \sum_{r=0}^{\infty} \beta_{m,r} \nu_{r,n}$$

$$(3.36) \quad h_{m,n} = v\mu_{m,n}/\beta [ \text{Fek}_m(\epsilon_0) \text{Ce}_m'(\epsilon_0)/\text{Ce}_m(\epsilon_0) - \text{Fek}_m'(\epsilon_0) \gamma_{1h}^2/\gamma_2^2 ]$$

$$(3.37) \quad s_{m,n} = v\beta_{m,n}/\beta [ \epsilon_1 \text{Gek}_m(\epsilon_0) \text{se}_m'(\epsilon_0)/\text{se}_m(\epsilon_0) - \epsilon_0 \text{Gek}_m'(\epsilon_0) \gamma_{1h}^2/\gamma_2^2 ]$$

$$(3.38) \quad t_{m,n} = (1 - \gamma_{1h}^2/\gamma_2^2) \text{Fek}_m(\epsilon_0) \sum_{r=0}^{\infty} \alpha_{m,r} \psi_{r,n}$$

For a nontrivial solution, the infinite determinant of Eq.(3.34) must vanish.

The propagation constant  $\beta$  can then be determined from the roots of this infinite determinant.

The infinite determinant for odd values of  $m$  and  $n$  is

$$(3.39) \quad \begin{vmatrix} s_{11} & t_{11} & s_{31} & t_{31} & - & - \\ g_{11} & h_{11} & g_{31} & h_{31} & - & - \\ s_{13} & t_{13} & s_{33} & t_{33} & - & - \\ g_{13} & h_{13} & g_{33} & h_{33} & - & - \\ : & : & : & : & - & - \\ : & : & : & : & - & - \end{vmatrix} = 0$$

and for even values of  $m$  and  $n$  is

$$(3.40) \quad \begin{vmatrix} h_{00} & g_{20} & h_{20} & g_{40} & - & - \\ t_{02} & s_{22} & t_{22} & s_{42} & - & - \\ h_{02} & g_{22} & h_{22} & g_{42} & - & - \\ t_{04} & s_{24} & t_{24} & s_{44} & - & - \\ : & : & : & : & - & - \\ : & : & : & : & - & - \end{vmatrix} = 0$$

### 3.2 ODD MODES

The axial components of the field for odd modes are

$$(3.41) \quad \begin{aligned} E_{z1} &= \sum_{m=0}^{\infty} A_{1m} Ce_m(\xi, \gamma_{1e}^2) ce_m(\eta, \gamma_{1e}^2) \\ H_{z1} &= \sum_{m=1}^{\infty} B_{1m} se_m(\xi, \gamma_{1h}^2) se_m(\eta, \gamma_{1h}^2) \end{aligned}$$

for  $0 \leq \xi \leq \xi_0$

$$(3.42) \quad \begin{aligned} E_{z2} &= \sum_{m=1}^{\infty} A_{2m} Fek_m(\xi, \gamma_2^2) ce_m(\eta, \gamma_2^2) \\ H_{z2} &= \sum_{m=1}^{\infty} B_{2m} Gek_m(\xi, \gamma_2^2) se_m(\eta, \gamma_2^2) \end{aligned}$$

for  $\xi_0 \leq \xi < \infty$

where  $A_{1m}$  and  $B_{1m}$ ,  $i = 1, 2$  are arbitrary constants, and  $\gamma_{1e}^2$ ,  $\gamma_{1h}^2$ , and  $\gamma_2^2$  are given in Eq.(3.4).

The transverse field components are

for  $0 \leq \xi \leq \xi_0$

$$(3.43) \quad \begin{aligned} E_{\xi 1} &= -1/(\nu^2 \mu \epsilon_1 - \beta^2) L \\ &\quad \left( \beta \sum_{m=0}^{\infty} A_{1m} Ce_m'(\xi, \gamma_{1e}^2) ce_m(\eta, \gamma_{1e}^2) \right. \\ &\quad \left. + \nu \mu \sum_{m=1}^{\infty} B_{1m} se_m(\xi, \gamma_{1h}^2) se_m'(\eta, \gamma_{1h}^2) \right) \end{aligned}$$

$$(3.44) \quad \begin{aligned} E_{\eta 1} &= -1/(\nu^2 \mu \epsilon_1 - \beta^2) L \\ &\quad \left( \beta \sum_{m=0}^{\infty} A_{1m} Ce_m(\xi, \gamma_{1e}^2) ce_m'(\eta, \gamma_{1e}^2) \right. \\ &\quad \left. - \nu \mu \sum_{m=1}^{\infty} B_{1m} se_m'(\xi, \gamma_{1h}^2) se_m(\eta, \gamma_{1h}^2) \right) \end{aligned}$$

$$(3.45) \quad \begin{aligned} H_{\xi 1} &= -1/(\nu^2 \mu \epsilon_1 - \beta^2) L \\ &\quad \left( -\nu \epsilon_1 \sum_{m=0}^{\infty} A_{1m} Ce_m(\xi, \gamma_{1e}^2) ce_m'(\eta, \gamma_{1e}^2) \right. \\ &\quad \left. + \beta \sum_{m=1}^{\infty} B_{1m} se_m'(\xi, \gamma_{1h}^2) se_m(\eta, \gamma_{1h}^2) \right) \end{aligned}$$

$$(3.46) \quad \begin{aligned} H_{\eta 1} &= -1/(\nu^2 \mu \epsilon_1 - \beta^2) L \\ &\quad \left( \nu \epsilon_1 \sum_{m=0}^{\infty} A_{1m} Ce_m'(\xi, \gamma_{1e}^2) ce_m(\eta, \gamma_{1e}^2) \right. \\ &\quad \left. + \beta \sum_{m=1}^{\infty} B_{1m} se_m(\xi, \gamma_{1h}^2) se_m'(\eta, \gamma_{1h}^2) \right) \end{aligned}$$

for  $\xi_0 \leq \xi < \infty$

$$(3.47) \quad E\xi_2 = -1/(\nu^2\mu\epsilon_0 - \beta^2)L \\ \left[ \beta \sum_{m=0}^{\infty} A_{2m} Fek_m'(\xi, \gamma_2^2) ce_m(\gamma, \gamma_2^2) \right. \\ \left. + \nu\mu \sum_{m=1}^{\infty} B_{2m} Gek_m(\xi, \gamma_2^2) se_m'(\gamma, \gamma_2^2) \right]$$

$$(3.48) \quad E\eta_2 = -1/(\nu^2\mu\epsilon_0 - \beta^2)L \\ \left[ \beta \sum_{m=0}^{\infty} A_{2m} Fek_m(\xi, \gamma_2^2) ce_m'(\gamma, \gamma_2^2) \right. \\ \left. - \nu\mu \sum_{m=1}^{\infty} B_{2m} Gek_m'(\xi, \gamma_2^2) se_m(\gamma, \gamma_2^2) \right]$$

$$(3.49) \quad H\xi_2 = -1/(\nu^2\mu\epsilon_0 - \beta^2)L \\ \left[ -\nu\epsilon_0 \sum_{m=0}^{\infty} A_{2m} Fek_m(\xi, \gamma_2^2) ce_m'(\gamma, \gamma_2^2) \right. \\ \left. + \beta \sum_{m=1}^{\infty} B_{2m} Gek_m'(\xi, \gamma_2^2) se_m(\gamma, \gamma_2^2) \right]$$

$$(3.50) \quad H\eta_2 = -1/(\nu^2\mu\epsilon_0 - \beta^2)L \\ \left[ \nu\epsilon_0 \sum_{m=0}^{\infty} A_{2m} Fek_m'(\xi, \gamma_2^2) ce_m(\gamma, \gamma_2^2) \right. \\ \left. + \beta \sum_{m=1}^{\infty} B_{2m} Gek_m(\xi, \gamma_2^2) se_m'(\gamma, \gamma_2^2) \right]$$

The derivative with respect to  $\xi$  or  $\gamma$  is denoted by the prime.

Equating the tangential fields at the boundary surface,  $\xi = \xi_0$ , gives

$$(3.51) \quad \sum_{m=0}^{\infty} A_{1m} Ce_m(\xi_0) ce_m(\gamma) = \sum_{m=0}^{\infty} A_{2m} Fek_m(\xi_0) ce_m^*(\gamma)$$

$$(3.52) \quad \sum_{m=1}^{\infty} B_{1m} Se_m(\xi_0) se_m(\gamma) = \sum_{m=1}^{\infty} B_{2m} Gek_m(\xi_0) se_m^*(\gamma)$$

$$(3.53) \quad 1/(\nu^2\mu\epsilon_1 - \beta^2) \left[ \beta \sum_{m=0}^{\infty} A_{1m} Ce_m(\xi_0) ce_m'(\gamma) \right. \\ \left. - \nu\mu \sum_{m=1}^{\infty} B_{1m} Se_m'(\xi_0) se_m(\gamma) \right] \\ = 1/(\nu^2\mu\epsilon_0 - \beta^2) \left[ \beta \sum_{m=0}^{\infty} A_{2m} Fek_m(\xi_0) ce_m^{*'}(\gamma) \right. \\ \left. - \nu\mu \sum_{m=1}^{\infty} B_{2m} Gek_m'(\xi_0) se_m^{*'}(\gamma) \right]$$

$$(3.54) \quad 1/(\nu^2\mu\epsilon_1 - \beta^2) \left[ \nu\epsilon_1 \sum_{m=0}^{\infty} A_{1m} Ce_m'(\xi_0) ce_m(\gamma) \right. \\ \left. + \beta \sum_{m=1}^{\infty} B_{1m} Se_m(\xi_0) se_m'(\gamma) \right] \\ = 1/(\nu^2\mu\epsilon_0 - \beta^2) \left[ \nu\epsilon_0 \sum_{m=0}^{\infty} A_{2m} Fek_m'(\xi_0) ce_m^{*'}(\gamma) \right. \\ \left. + \beta \sum_{m=1}^{\infty} B_{2m} Gek_m(\xi_0) se_m^{*'}(\gamma) \right]$$

The abbreviations

$$(3.55) \quad Ce_m(\xi_0) = Ce_m(\xi_0, \gamma_{1e}^2)$$

$$(3.56) \quad ce_m(\eta) = ce_m(\eta, \gamma_{1e}^2)$$

$$(3.57) \quad se_m(\xi_0) = se_m(\xi_0, \gamma_{1h}^2)$$

$$(3.58) \quad se_m(\eta) = se_m(\eta, \gamma_{1h}^2)$$

$$(3.59) \quad Fek_m(\xi_0) = Fek_m(\xi_0, \gamma_2^2)$$

$$(3.60) \quad ce_m^*(\eta) = ce_m(\eta, \gamma_2^2)$$

$$(3.61) \quad Gek_m(\xi_0) = Gek_m(\xi_0, \gamma_2^2)$$

$$(3.62) \quad se_m^*(\eta) = se_m(\eta, \gamma_2^2)$$

have been used.

Multiplying both sides of Eqs. (3.51) and (3.54) by  $ce_n(\eta)$  and Eqs. (3.52) and (3.53) by  $se_n(\eta)$ , integrating with respect to  $\eta$  from 0 to  $2\pi$ , and applying the orthogonality relations of the angular Mathieu functions leads to

$$(3.63) \quad A_{1n}Ce_n = \sum_{m=0}^{\infty} A_{2m}Fek_m\beta_{m,n}$$

$$(3.64) \quad B_{1n}se_n = \sum_{m=1}^{\infty} B_{2m}Gek_m\alpha_{m,n}$$

$$(3.65) \quad \beta \sum_{m=0}^{\infty} A_{1m}Ce_m\psi_{n,m} - \nu\mu B_{1n}se_n' = \\ \beta\gamma_{1h}^2/\gamma_2^2 \sum_{m=0}^{\infty} A_{2m}Fek_m\psi_{n,m} - \nu\mu\gamma_{1h}^2/\gamma_2^2 \sum_{m=1}^{\infty} B_{2m}Gek_m'\alpha_{n,m}$$

$$(3.66) \quad \nu\epsilon_1 A_{1n}Ce_n + \beta \sum_{m=1}^{\infty} B_{1m}se_m\psi_{n,m} = \\ \nu\epsilon_0\gamma_{1h}^2/\gamma_2^2 \sum_{m=0}^{\infty} A_{2m}Fek_m'\beta_{n,m} + \beta\gamma_{1h}^2/\gamma_2^2 \sum_{m=1}^{\infty} B_{2m}Gek_m\psi_{n,m}$$

The prime over the summation sign is used to indicate that either odd or even values of  $m$  are used accordingly as to whether  $n$  is odd or even.

$\alpha_{m,n}$ ,  $\beta_{m,n}$ ,  $\psi_{m,n}$  and  $\mu_{m,n}$  are given by the following

$$(3.67) \quad \alpha_{m,n} = \int_0^{2\pi} se_m^*(\eta)se_n(\eta) d\eta / \int_0^{2\pi} se_n^2(\eta) d\eta$$

$$(3.68) \quad \beta_{m,n} = \int_0^{2\pi} ce_m^*(\eta)ce_n(\eta) d\eta / \int_0^{2\pi} ce_n^2(\eta) d\eta$$

$$(3.69) \quad \psi_{m,n} = \int_0^{2\pi} se_m^*(\eta)ce_n(\eta) d\eta / \int_0^{2\pi} ce_n^2(\eta) d\eta$$

$$(3.70) \quad \mu_{m,n} = \int_0^{2\pi} ce_m^*(\eta)se_n(\eta) d\eta / \int_0^{2\pi} se_n^2(\eta) d\eta$$

Making use of Eqs. (3.63) and (3.64), Eqs. (3.65) and (3.66) yields two sets of infinite homogeneous equations

$$(3.71) \quad \sum_{m=0}^{\infty} A_{2m} s_{m,n} + \sum_{m=1}^{\infty} B_{2m} t_{m,n} = 0$$

$$\sum_{m=0}^{\infty} A_{2m} g_{m,n} + \sum_{m=1}^{\infty} B_{2m} h_{m,n} = 0$$

where

$$(3.72) \quad g_{m,n} = -(1 - \gamma_{1h}^2/\gamma_2^2) \text{Fek}_m(\epsilon_0) \sum_{r=0}^{\infty} \beta_{m,r} \nu_{r,n}$$

$$(3.73) \quad h_{m,n} = \nu \mu \alpha_{m,n} / \beta \{ \text{Gek}_m(\epsilon_0) \text{se}_m'(\epsilon_0) / \text{se}_m(\epsilon_0) - \text{Gek}_m'(\epsilon_0) \gamma_{1h}^2/\gamma_2^2 \}$$

$$(3.74) \quad s_{m,n} = \nu \beta_{m,n} / \beta \{ \epsilon_1 \text{Fek}_m(\epsilon_0) \text{ce}_m'(\epsilon_0) / \text{ce}_m(\epsilon_0) - \epsilon_0 \text{Fek}_m'(\epsilon_0) \gamma_{1h}^2/\gamma_2^2 \}$$

$$(3.75) \quad t_{m,n} = (1 - \gamma_{1h}^2/\gamma_2^2) \text{Gek}_m(\epsilon_0) \sum_{r=1}^{\infty} \alpha_{m,r} \psi_{r,n}$$

For a nontrivial solution, the infinite determinant of Eq. (3.71) must vanish. The propagation constant  $\beta$  can then be determined from the roots of this infinite determinant.

The infinite determinant for odd values of  $m$  and  $n$  is

$$(3.76) \quad \begin{vmatrix} s_{11} & t_{11} & s_{31} & t_{31} & - & - \\ g_{11} & h_{11} & g_{31} & h_{31} & - & - \\ s_{13} & t_{13} & s_{33} & t_{33} & - & - \\ g_{13} & h_{13} & g_{33} & h_{33} & - & - \\ : & : & : & : & - & - \\ : & : & : & : & - & - \end{vmatrix} = 0$$

and for even values of  $m$  and  $n$  is

$$(3.77) \quad \begin{vmatrix} s_{00} & s_{20} & t_{20} & s_{40} & - & - \\ s_{02} & s_{22} & t_{22} & s_{42} & - & - \\ g_{02} & g_{22} & h_{22} & g_{42} & - & - \\ s_{04} & s_{24} & t_{24} & s_{44} & - & - \\ : & : & : & : & - & - \\ : & : & : & : & - & - \end{vmatrix} = 0$$

#### 4. WEAKLY GUIDING APPROXIMATION

The exact characteristic equations obtained in Chapter 3 are valid for an anisotropic elliptical fiber with any eccentricities. These equations are also applicable to the fiber with any refractive index differences between the core and cladding material. However, for most of practical fibers, the difference in the refractive indices of the core and the cladding is typically very small. The simplified characteristic equations can be obtained under this condition which is known as the weakly guiding approximation.

Applying the weakly guiding approximation results in the following

$$(4.1) \quad \gamma_2^2 = \gamma_{1h}^2 + kv^2\mu\epsilon_1 (1 - \epsilon_0/\epsilon_1) \approx \gamma_{1h}^2$$

$$(4.2) \quad 1 - \gamma_{1h}^2/\gamma_2^2 \approx 0.$$

##### 4.1 EVEN MODES

Applying Eqs.(4.1) and (4.2) into Eqs.(3.20) and (3.30) yields the equations

$$(4.3) \quad ce_n(\eta, \gamma_2^2) \approx ce_n(\eta, \gamma_{1h}^2)$$

$$(4.4) \quad a_{n,n} \approx \int_0^{2\pi} ce_n(\eta, \gamma_{1h}^2) ce_n(\eta, \gamma_{1h}^2) d\eta / \int_0^{2\pi} ce_n^2(\eta, \gamma_{1h}^2) d\eta \\ = \Delta_{n,n}$$

where  $\Delta_{n,n}$  is the Kronecker delta which is zero when  $m \neq n$  and is unity when  $m = n$ .

Substituting Equations (4.1) - (4.4) into Equations (3.35) - (3.38), the infinite determinants for even modes become

$$(4.5) \quad \prod_m (g_{m,m} t_{m,m} - h_{m,m} s_{m,m}) = 0$$

or

$$(4.6) \quad g_{m,m} t_{m,m} - h_{m,m} s_{m,m} = 0$$

for  $m = 0, 1, 2, \dots$



By substituting Equations (3.35) - (3.38) into Eq.(4.6), the following equation is obtained

$$\begin{aligned}
 & -(1 - \gamma_{1h}^2/\gamma_2^2)(\epsilon_1/\epsilon_0 - \gamma_{1h}^2/\gamma_2^2) \left( \sum_{n=1}^{\infty} \beta_{m,n} \nu_{n,m} \sum_{n=0}^{\infty} \alpha_{m,n} \psi_{n,m} / \alpha_{m,n} \beta_{m,n} \right) \\
 (4.7) \quad & = [C_{em}'(\xi_0)/C_{em}(\xi_0) - (\gamma_{1h}^2/\gamma_2^2) Fek_m'(\xi_0)/Fek_m(\xi_0)] \\
 & [(\epsilon_1/\epsilon_0) Se_m'(\xi_0)/Se_m(\xi_0) - (\gamma_{1h}^2/\gamma_2^2) Gek_m'(\xi_0)/Gek_m(\xi_0)].
 \end{aligned}$$

This is the simplified characteristic equation for even modes compared to the infinite determinants as given in Eqs.(3.39) and (3.40). When the elliptical rod degenerates to a circular rod, the simplified characteristic equation becomes that of the anisotropic circular fiber.

#### 4.2 ODD MODES

Applying Eqs.(4.1) and (4.2) in Eqs.(3.58) and (3.67) yields the equations

$$(4.8) \quad se_n(\eta, \gamma_2^2) = se_n(\eta, \gamma_{1h}^2)$$

$$\begin{aligned}
 (4.9) \quad \alpha_{m,n} &= \int_0^{2\pi} se_m(\eta, \gamma_{1h}^2) se_n(\eta, \gamma_{1h}^2) d\eta / \int_0^{2\pi} se_m^2(\eta, \gamma_{1h}^2) d\eta \\
 &= \Delta_{m,n}
 \end{aligned}$$

where  $\Delta_{m,n}$  is the Kronecker delta which is zero when  $m \neq n$  and is unity when  $m = n$ .

Substituting Equations (4.1) - (4.2) and (4.8) - (4.9) into Equations (3.72) - (3.75), the infinite determinants for odd modes become

$$(4.10) \quad \prod_m (g_{m,m} t_{m,m} - h_{m,m} s_{m,m}) = 0$$

or

$$(4.11) \quad g_{m,m} t_{m,m} - h_{m,m} s_{m,m} = 0$$

for  $m = 0, 1, 2, \dots$

By substituting Equations (3.72) - (3.75) into Eq.(4.11), the following equation is obtained

$$\begin{aligned}
 & -(1 - \gamma_{1h}^2/\gamma_2^2)(\epsilon_1/\epsilon_0 - \gamma_{1h}^2/\gamma_2^2) \left( \sum_{n=0}^{\infty} \beta_{m,n} \nu_{n,m} \sum_{n=1}^{\infty} \alpha_{m,n} \psi_{n,m}/\alpha_{m,n} \beta_{m,n} \right) \\
 (4.12) \quad & = [se_m'(\xi_0)/se_m(\xi_0) - (\gamma_{1h}^2/\gamma_2^2)Gek_m'(\xi_0)/Gek_m(\xi_0)] \\
 & [(\epsilon_1/\epsilon_0)Ce_m'(\xi_0)/Ce_m(\xi_0) - (\gamma_{1h}^2/\gamma_2^2)Fek_m'(\xi_0)/Fek_m(\xi_0)].
 \end{aligned}$$

This is the simplified characteristic equation for odd modes compared to the infinite determinants given in Eqs.(3.76) and (3.77). When the elliptical rod degenerates to a circular rod, the simplified characteristic equation becomes that of the anisotropic circular fiber.

## 5. NUMERICAL RESULTS FOR PROPAGATION CONSTANTS

### 5.1 ISOTROPIC ELLIPTICAL FIBERS

When the anisotropic parameter  $g$  in Eq.(3.1) is equal to unity, the simplified characteristic equations in Eq.(4.7) and Eq.(4.12) become that of an isotropic elliptic guide. In Figures 2 through 5, the normalized guide wavelength  $\lambda/\lambda_0$  for the isotropic elliptical fibers is plotted as a function of the normalized cross-section area and normalized major axis for  $\epsilon_1/\epsilon_0 = 2.5$  and for the various values of  $\epsilon_0$ . These results are compared with those given by Yeh[2] which are indicated by symbols; the results are in close agreement.

For the  $eHE_{11}$  mode, it can be seen in Figure 2 that the normalized guide wavelength is almost equal to unity for the very small value of the cross-section area. This indicates that the geometry of the waveguide has no effect on the normalized guide wavelength when the wavelength is much larger than the physical dimension of the core of fibers. For a fixed value of cross-section area, the normalized guide wavelength is smaller for larger the value of  $\epsilon_0$ . This indicates that more energy is carried inside of the circular core than the elliptical core. As the normalized cross-section area becomes larger, the difference in the normalized guide wavelengths for varying  $\epsilon_0$  becomes small again. This is since most of the energy is carried inside of the core and the geometry of waveguide has no effect on the normalized guide wavelength.

However, as observed in Figure 3, the  $oHE_{11}$  mode is different from the  $eHE_{11}$  mode in that the difference in the normalized guide wavelengths for varying  $\epsilon_0$  is smaller than that of  $eHE_{11}$  when the value of normalized cross-section area is fixed. This small difference is due to the fact that the electric lines are being compressed such that the field density is more

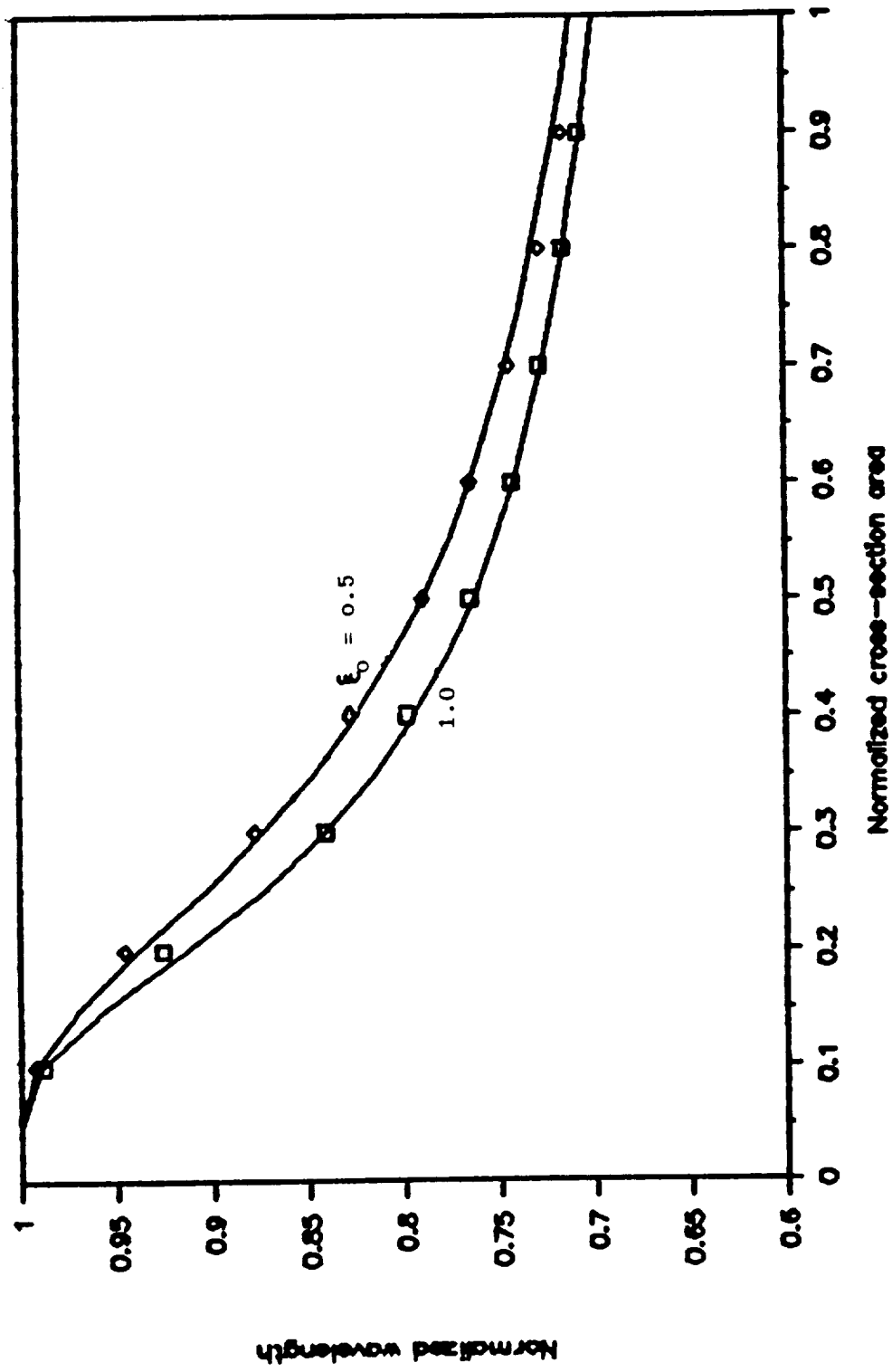


Figure 2. Normalized wavelength for isotropic fiber as a function of normalized cross-section area for even modes. Symbols are from Yeh[2].

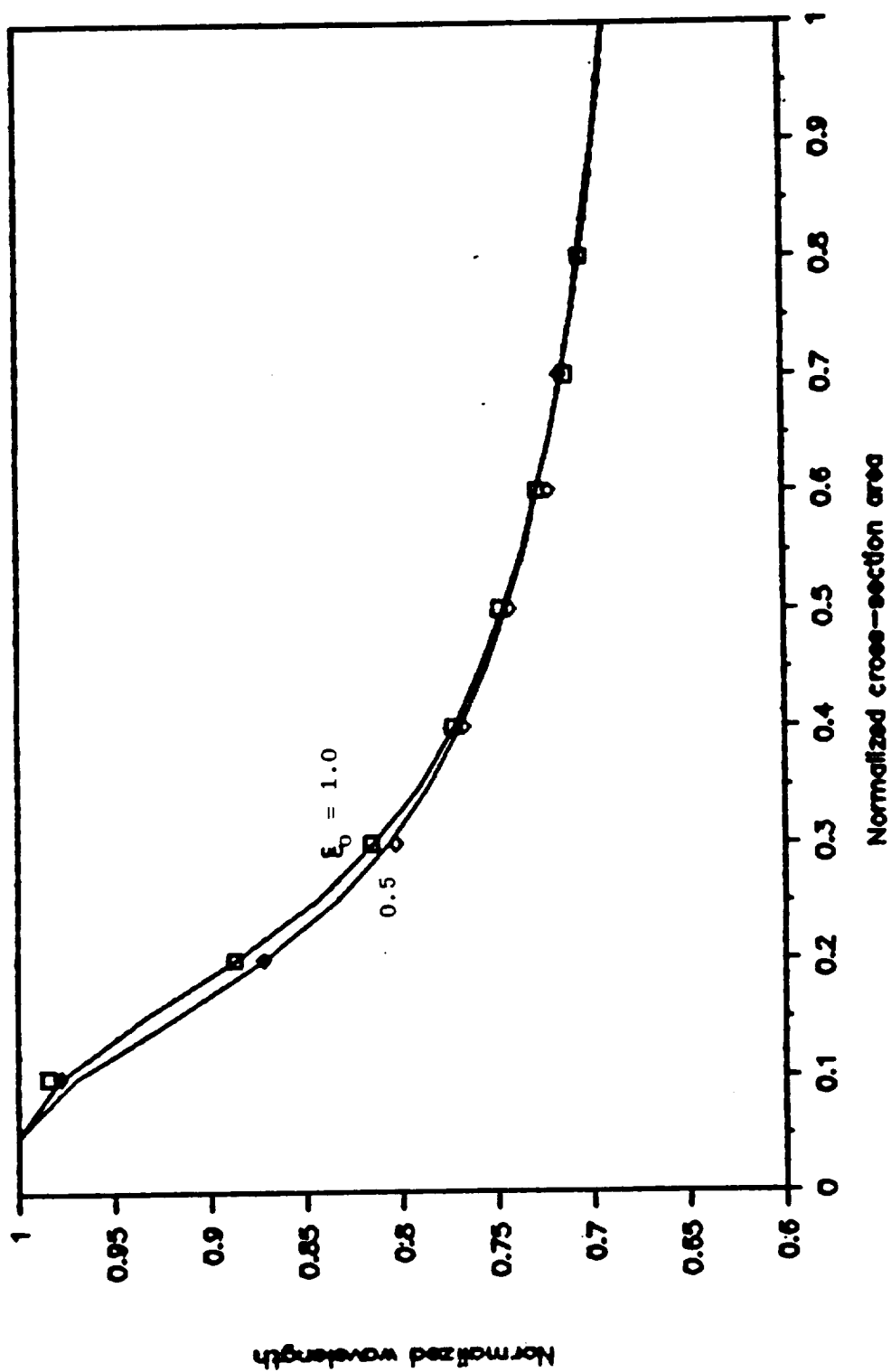


Figure 3. Normalized wavelength for isotropic fiber as a function of normalized cross-section area for odd modes. Symbols are from Yeh[2].

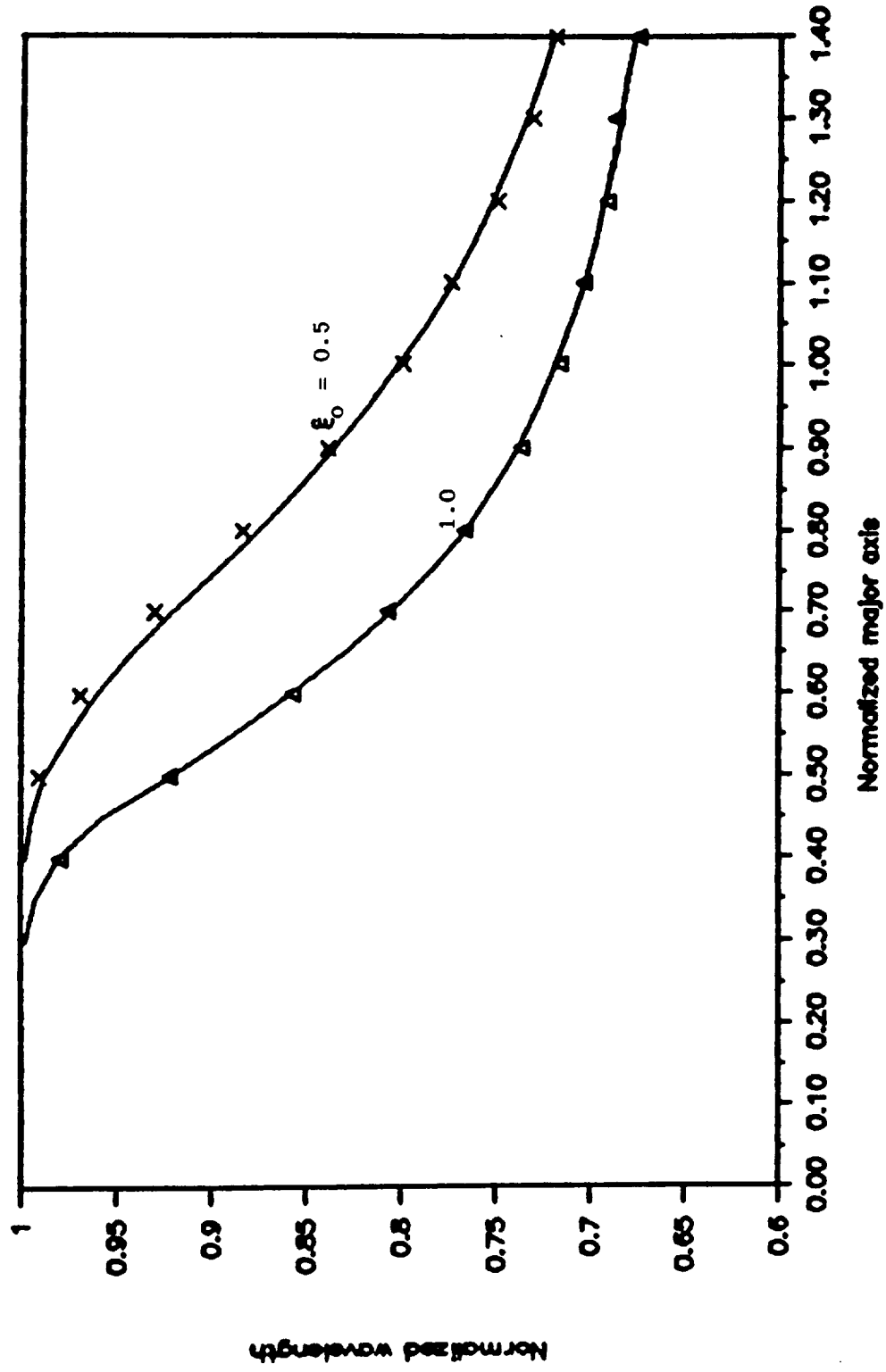


Figure 4. Normalized wavelength for isotropic fiber as a function of normalized major axis for even modes. Symbols are from Yeh[2].

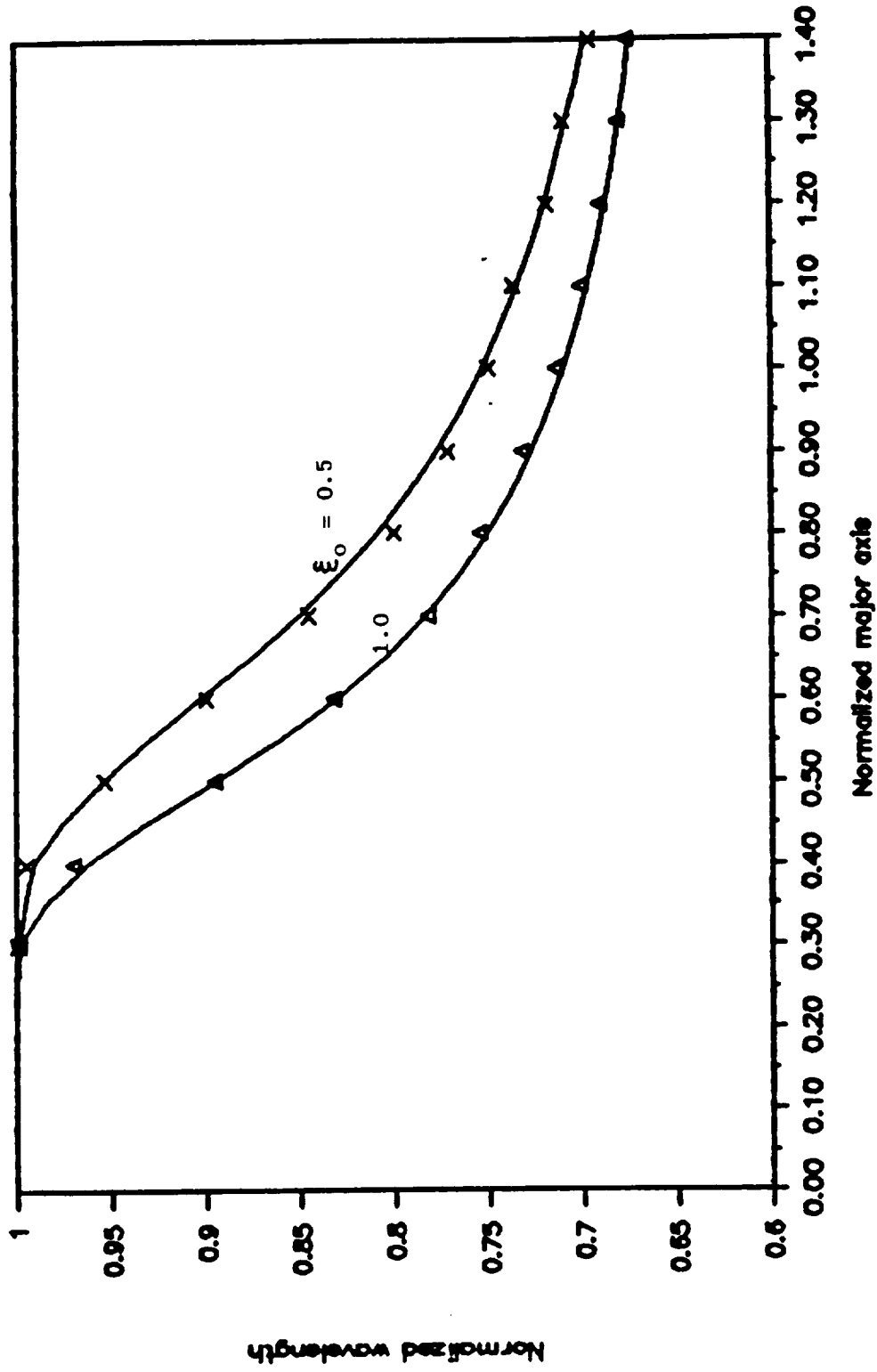


Figure 5. Normalized wavelength for isotropic fiber as a function of normalized major axis for odd modes. Symbols are from Yeh[21].

concentrated inside the waveguide. For a fixed value of cross-section area, more energy is carried inside of the elliptical core than the circular core since the normalized guide wavelength is smaller for smaller the value of  $\epsilon_0$ .

In Figures 4 and 5, the normalized guide wavelength is plotted against the normalized major axis for various values of  $\epsilon_0$  and for  $\epsilon_1/\epsilon_0 = 2.5$ . In these figures, the difference in the normalized guide wavelengths for varying  $\epsilon_0$  is larger than those in Figures 2 and 3 since there is more binding dielectric material in a circular core than in a flatter rod (i.e. smaller  $\epsilon_0$ ) when the value of normalized major axis is fixed.

## 5.2 ANISOTROPIC ELLIPTICAL FIBERS

In Figures 6 and 7, the normalized guide wavelength  $\lambda/\lambda_0$  for an anisotropic elliptical fiber is plotted as a function of the normalized cross-section area for various values of anisotropy and for  $\epsilon_1/\epsilon_0 = 2.5$  and  $\epsilon_0 = 0.5$ . These figures indicate that the geometry of the waveguide and anisotropy of the core have no effect when the wavelength is much larger than the physical dimension of the core of fibers which indicate that most of the energy is carried outside of the core. For a fixed value of cross-section area, the normalized guide wavelength is smaller for larger the value of anisotropy. This condition indicates that the field intensity is more concentrated in the core, thus indicating that more energy is carried inside of the core. As the normalized cross-section area becomes larger, the difference in the normalized guide wavelengths for the varying anisotropy becomes smaller again. This indicates that the geometry and anisotropy of waveguide have a smaller effect on the normalized guide wavelength.



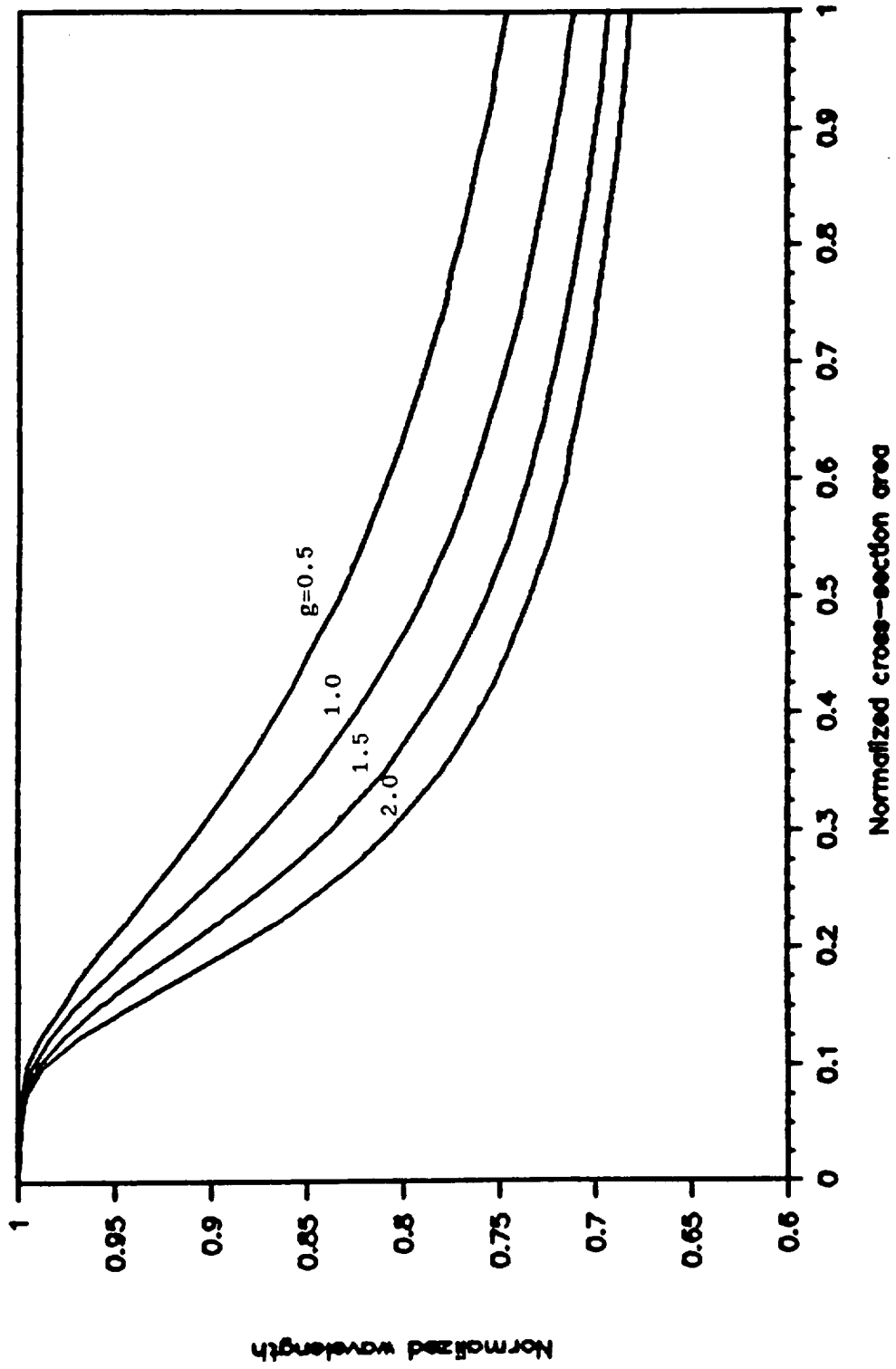


Figure 6. Normalized wavelength for anisotropic fiber as a function of normalized cross-section area for even modes.

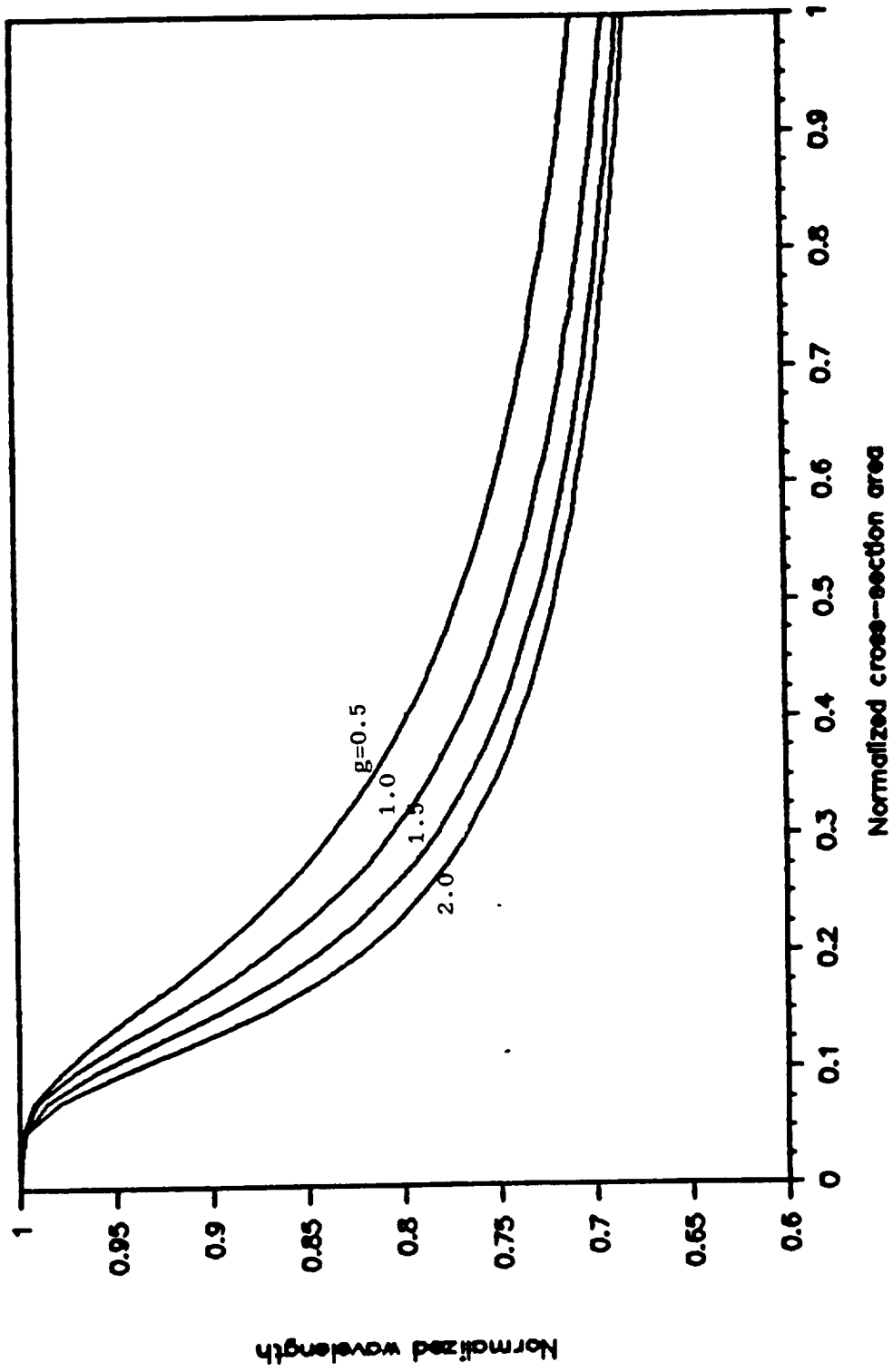


Figure 7. Normalized wavelength for anisotropic fiber as a function of normalized cross-section area for odd modes.

The normalized guide wavelength in Figures 8 and 9 is plotted against the normalized major axis for the various values of anisotropy and for  $\epsilon_1/\epsilon_0 = 2.5$  and  $\epsilon_0 = 0.5$ . The effect of anisotropy on the normalized guide wavelength is similar to those in Figures 6 and 7.

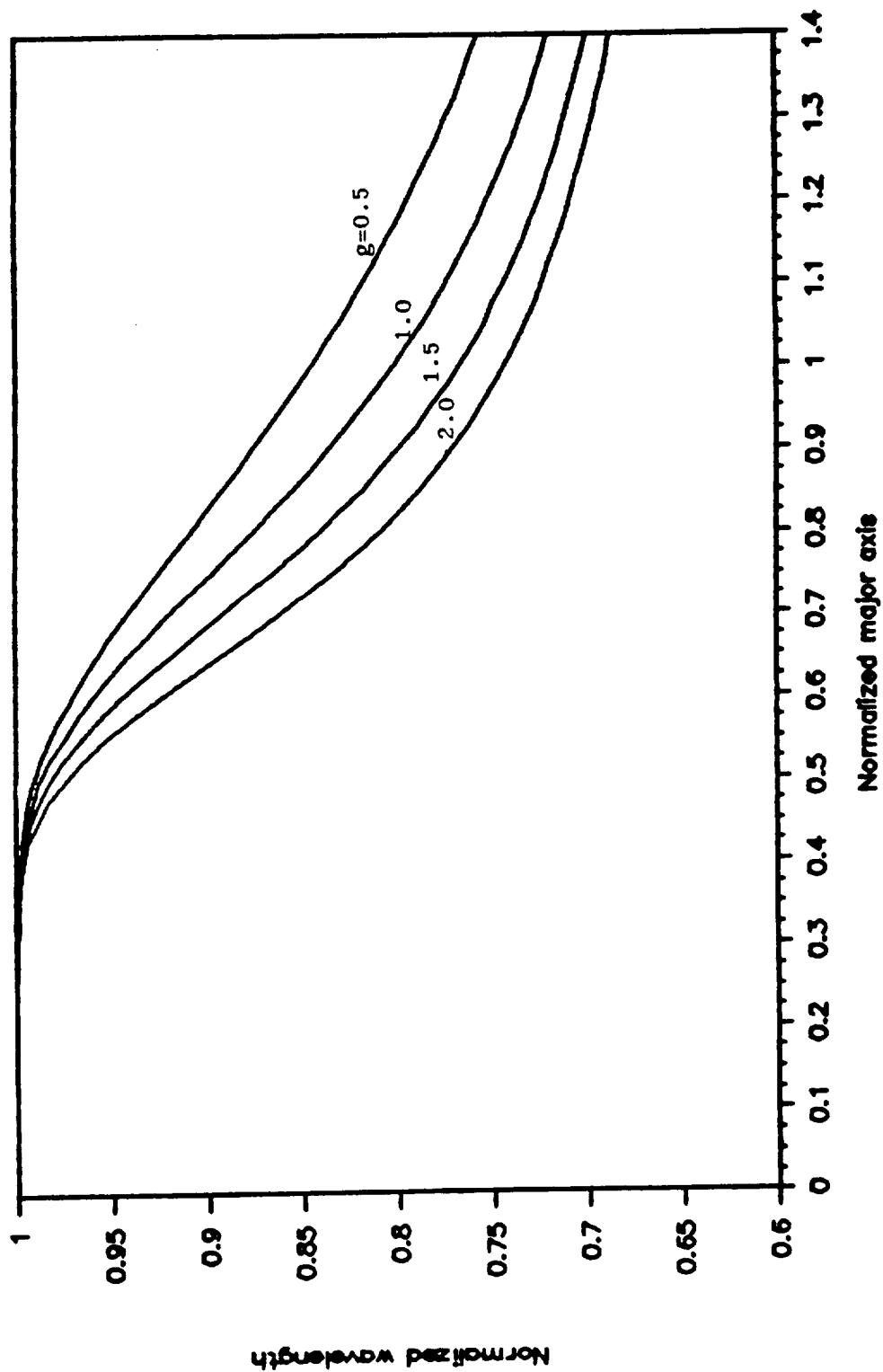


Figure 8. Normalized wavelength for anisotropic fiber as a function of normalized major axis for even modes.

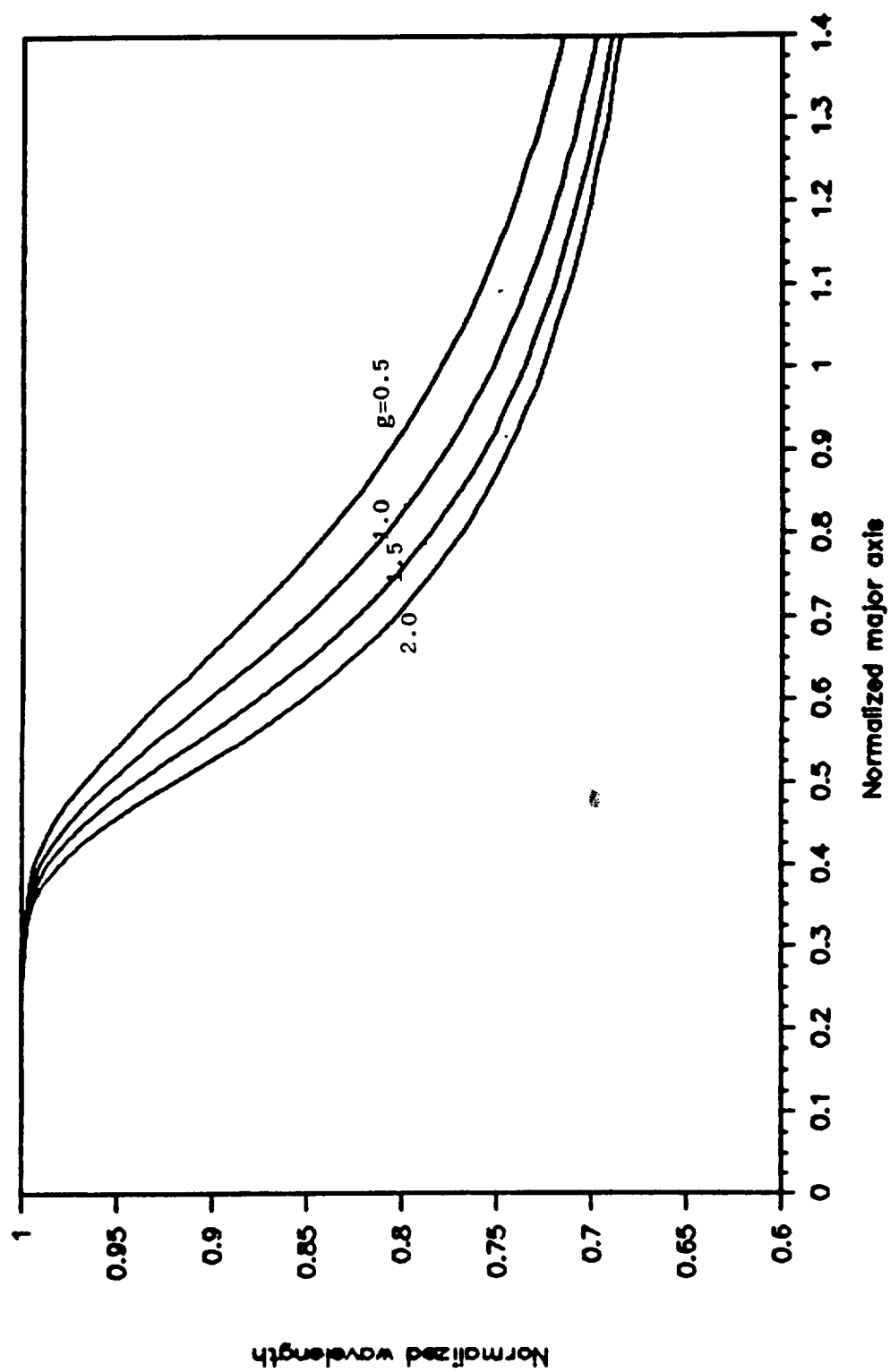


Figure 9. Normalized wavelength for anisotropic fiber as a function of normalized major axis for odd modes.

## 6. POWER CONSIDERATIONS

The power along the  $z$  axis in the medium  $i$  of the fiber may be obtained by integrating the poynting vector over the surface area,

$$(6.1) \quad \begin{aligned} P_i &= 1/2 \int_S (\vec{E}_t \times \vec{H}_t^*) \cdot \hat{z} \, ds \\ &= 1/2 \int_{\xi=0}^{\xi_0} \int_0^{2\pi} (E_{\xi 1} H_{\eta 1}^* - E_{\eta 1} H_{\xi 1}^*) L^2 \, d\eta \, d\xi \end{aligned}$$

where

$$L = q[(\cosh 2\xi - \cos 2\eta)/2]^{1/2}$$

and  $\xi_0 = 0$  and  $\xi_2 = \infty$ .

### 6.1 EVEN MODES

Substituting Equations (3.5) through (3.8) into Eq.(6.1) and integrating over the core area yields

$$(6.2) \quad \begin{aligned} P_{core} &= (1/2\gamma_1^4) \int_0^{\xi_0} \{ \beta v \epsilon_1 \sum_{m=1}^{\infty} A_{1m}^2 \{ \kappa s e_m'^2 + s_{mn} s e_m^2 \} \\ &\quad + \beta v \mu \sum_{m=0}^{\infty} B_{1m}^2 \{ \kappa c e_m'^2 + c_{mn} c e_m^2 \} \} d\xi \\ &\quad + (\beta^2 + v^2 \mu \epsilon_1) / 2\gamma_1^4 \sum_{m=0}^{\infty} \sum_{n=1}^{\infty} B_{1m} A_{1n} T_{mn} \{ c e_m s e_n \}_{\xi_0}^{\xi_0} \\ &\quad + \beta v \mu / 2\gamma_1^4 \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} B_{1m} B_{1n} C_{mn} \{ c e_m' c e_n - c e_m c e_n' \}_{\xi_0}^{\xi_0} (a_m - a_n)^{-1} \\ &\quad + \beta v \epsilon_1 / 2\gamma_1^4 \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{1m} A_{1n} S_{mn} \{ s e_m' s e_n - s e_m s e_n' \}_{\xi_0}^{\xi_0} (b_m - b_n)^{-1} \end{aligned}$$

Similary, the power carried in the cladding is

$$(6.3) \quad \begin{aligned} P_{clad} &= (1/2\gamma_2^4) \int_{\xi_0}^{\infty} \{ \beta v \epsilon_0 \sum_{m=1}^{\infty} A_{2m}^2 \{ \kappa G e_k_m'^2 + s_{mn}^* G e_k_m^2 \} \\ &\quad + \beta v \mu \sum_{m=0}^{\infty} B_{2m}^2 \{ \kappa F e_k_m'^2 + c_{mn}^* F e_k_m^2 \} \} d\xi \\ &\quad + (\beta^2 + v^2 \mu \epsilon_0) / 2\gamma_2^4 \sum_{m=0}^{\infty} \sum_{n=1}^{\infty} B_{2m} A_{2n} T_{mn}^* \{ F e_k_m G e_k_n \}_{\xi_0}^{\infty} \\ &\quad + \beta v \mu / 2\gamma_2^4 \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} B_{2m} B_{2n} C_{mn}^* \{ F e_k_m' F e_k_n \\ &\quad \quad \quad - F e_k_m F e_k_n' \}_{\xi_0}^{\infty} (a_m - a_n)^{-1} \\ &\quad + \beta v \epsilon_0 / 2\gamma_2^4 \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{2m} A_{2n} S_{mn}^* \{ G e_k_m' G e_k_n \\ &\quad \quad \quad - G e_k_m G e_k_n' \}_{\xi_0}^{\infty} (b_m - b_n)^{-1} \end{aligned}$$

In Eqs.(6.2) and (6.3),  $a_m$  and  $b_m$  are the characteristic values of the even and odd Mathieu functions of order  $m$ , respectively. The prime over the summation sign is used to indicate that  $m = n$  is excluded. Also,  $\gamma_1^2 = 4\gamma_{1h}^2/q^2$  and  $\gamma_2^2 = 4\gamma_2^2/q^2$  are used.

The following abbreviations have been used,

$$(6.4) \quad C_{mn} = \int_0^{2\pi} ce_m'(\zeta) ce_n'(\zeta) d\zeta$$

$$(6.5) \quad S_{mn} = \int_0^{2\pi} se_m'(\zeta) se_n'(\zeta) d\zeta$$

$$(6.6) \quad T_{mn} = \int_0^{2\pi} ce_m'(\zeta) se_n(\zeta) d\zeta$$

$$(6.7) \quad C_{mn}^* = \int_0^{2\pi} ce_m^*(\zeta) ce_n^*(\zeta) d\zeta$$

$$(6.8) \quad S_{mn}^* = \int_0^{2\pi} se_m^*(\zeta) se_n^*(\zeta) d\zeta$$

$$(6.9) \quad T_{mn}^* = \int_0^{2\pi} ce_m^*(\zeta) se_n^*(\zeta) d\zeta.$$

The power distribution characteristics for the  $eHE_{11}$  mode is given in Figure 10. The fractional power carried by the core and cladding is plotted against the normalized major axis for the various values of anisotropy and for  $\epsilon_1/\epsilon_0 = 2.5$  and  $\epsilon_0 = 1.0$ . Most of the power is carried in the cladding near the cut-off and in the core far from the cut-off. For a fixed value of the normalized major axis, the more energy is concentrated inside of the core for larger the value of anisotropy and far from the cut-off.

## 6.2 ODD MODES

Substituting Equations (3.43) through (3.46) into Eq.(6.1) and integrating over the core area yields

$$(6.10) \quad P_{core} = (1/2\gamma_1^4) \int_0^{2\pi} \left[ \beta v \epsilon_1 \sum_{m=1}^{\infty} A_{1m}^2 ( \pi ce_m'^2 + C_{mn} ce_m'^2 ) \right. \\ \left. + \beta v \mu \sum_{n=1}^{\infty} B_{1n}^2 ( \pi se_n'^2 + S_{nn} se_n'^2 ) \right] d\zeta \\ + ( \beta^2 + v^2 \mu \epsilon_1 ) / 2\gamma_1^4 \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} B_{1m} A_{1n} T_{mn} [ se_m ce_n ]_0^{2\pi} \\ + \beta v \mu / 2\gamma_1^4 \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} B_{1m} B_{1n} S_{mn} [ se_m' se_n' - se_m se_n' ]_0^{2\pi} (b_m - b_n)^{-1} \\ + \beta v \epsilon_1 / 2\gamma_1^4 \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{1m} A_{1n} C_{mn} [ ce_m' ce_n' - ce_m ce_n' ]_0^{2\pi} (a_m - a_n)^{-1}$$

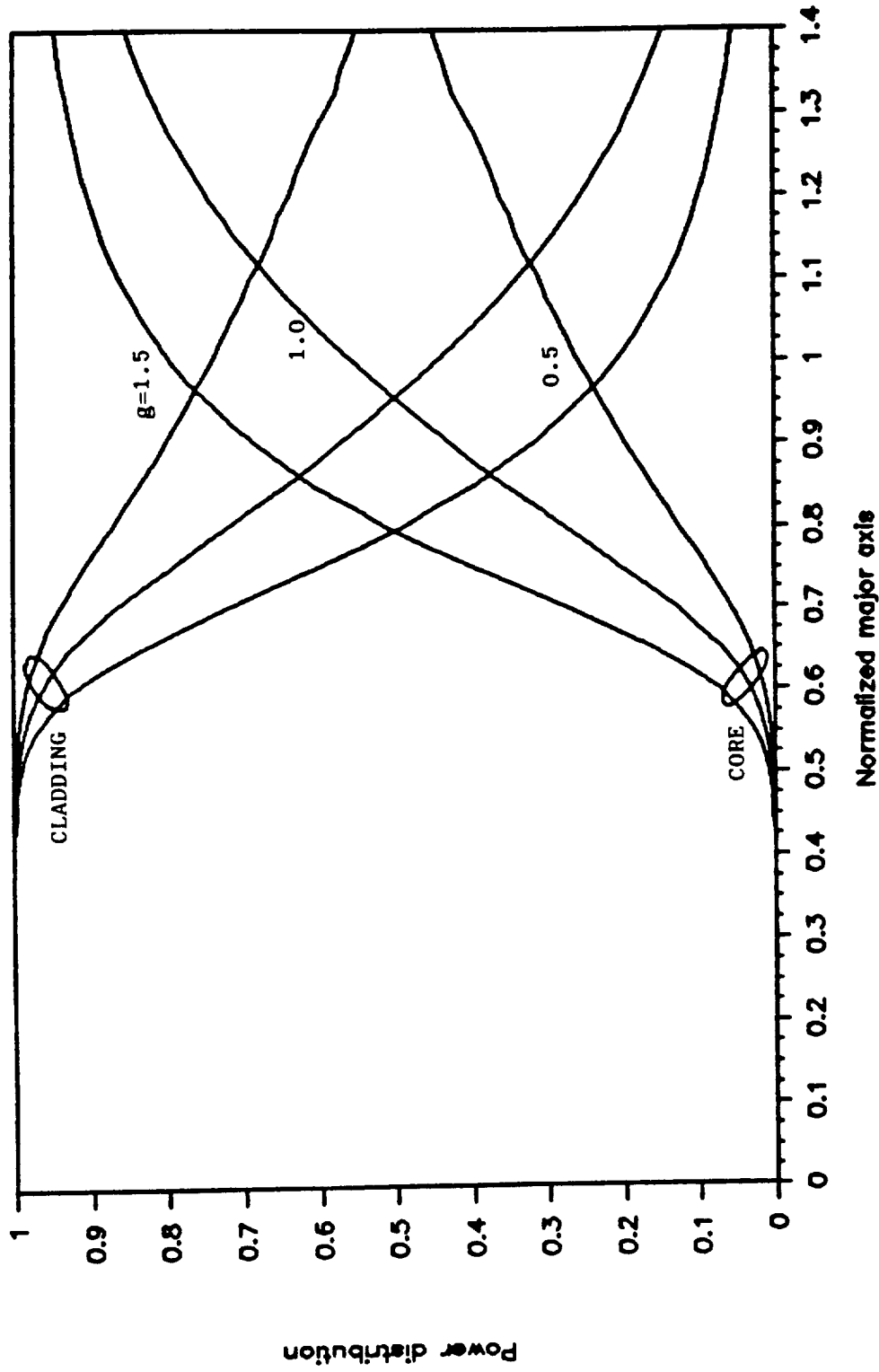


Figure 10. Power distribution characteristics for elliptical fiber as a function of normalized major axis for even modes.



Similary, the power carried in the cladding is

$$\begin{aligned}
 (6.11) \quad P_{clad} = & (1/2\gamma_2^4) \int_{\xi_0}^{\infty} \left[ \beta v \epsilon_0 \sum_{m=0}^{\infty} A_{2m}^2 \{ \pi P_{ek_m}'^2 + C_{mm}^* P_{ek_m}^2 \} \right. \\
 & + \beta v \mu \sum B_{2m}^2 \{ \pi G_{ek_m}'^2 + S_{mm}^* G_{ek_m}^2 \} \} d\xi \\
 & + (\beta^2 + v^2 \mu \epsilon_0) / 2\gamma_2^4 \sum_{m=1}^{\infty} \sum_{n=0}^{\infty} B_{2m} A_{2n} T_{mn}^* \{ G_{ek_m} P_{ek_n} \}_{\xi_0}^{\infty} \\
 & + \beta v \mu / 2\gamma_2^4 \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} B_{2m} B_{2n} S_{mn}^* \{ G_{ek_m}' G_{ek_n} \\
 & - G_{ek_m} G_{ek_n}' \}_{\xi_0}^{\infty} (b_m - b_n)^{-1} \\
 & + \beta v \epsilon_0 / 2\gamma_2^4 \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{2m} A_{2n} C_{mn}^* \{ P_{ek_m}' P_{ek_n} \\
 & - P_{ek_m} P_{ek_n}' \}_{\xi_0}^{\infty} (a_m - a_n)^{-1}
 \end{aligned}$$

In Eqs.(6.10) and (6.11),  $a_m$  and  $b_m$  are the characteristic values of the even and odd Mathieu functions of order  $m$ , respectively. The prime over the summation sign is used to indicate that  $m = n$  is excluded. Also,  $\gamma_1^2 = 4\gamma_{1h}^2/q^2$  and  $\gamma_2^2 = 4\gamma_2^2/q^2$  are used.

The following abbreviations have been used,

$$(6.12) \quad C_{mn} = \int_0^{2\pi} ce_m'(\eta) ce_n'(\eta) d\eta$$

$$(6.13) \quad S_{mn} = \int_0^{2\pi} se_m'(\eta) se_n'(\eta) d\eta$$

$$(6.14) \quad T_{mn} = \int_0^{2\pi} ce_m(\eta) se_n'(\eta) d\eta$$

$$(6.15) \quad C_{mn}^* = \int_0^{2\pi} ce_m^*(\eta) ce_n^*(\eta) d\eta$$

$$(6.16) \quad S_{mn}^* = \int_0^{2\pi} se_m^*(\eta) se_n^*(\eta) d\eta$$

$$(6.17) \quad T_{mn}^* = \int_0^{2\pi} ce_m^*(\eta) se_n^*(\eta) d\eta.$$

The power distribution characteristics for the  $oHE_{11}$  mode is given in Figure 11. The fractional power carried by the core and cladding is plotted against the normalized major axis for the various values of anisotropy and for  $\epsilon_1/\epsilon_0 = 2.5$  and  $\xi_0 = 1.0$ . Most of the power is carried in the cladding near the cut-off and in the core far from cut-off. However, the difference in the power distribution for  $oHE_{11}$  for the varying anisotropy is smaller than that of  $eHE_{11}$  for a fixed value of normalized major axis.

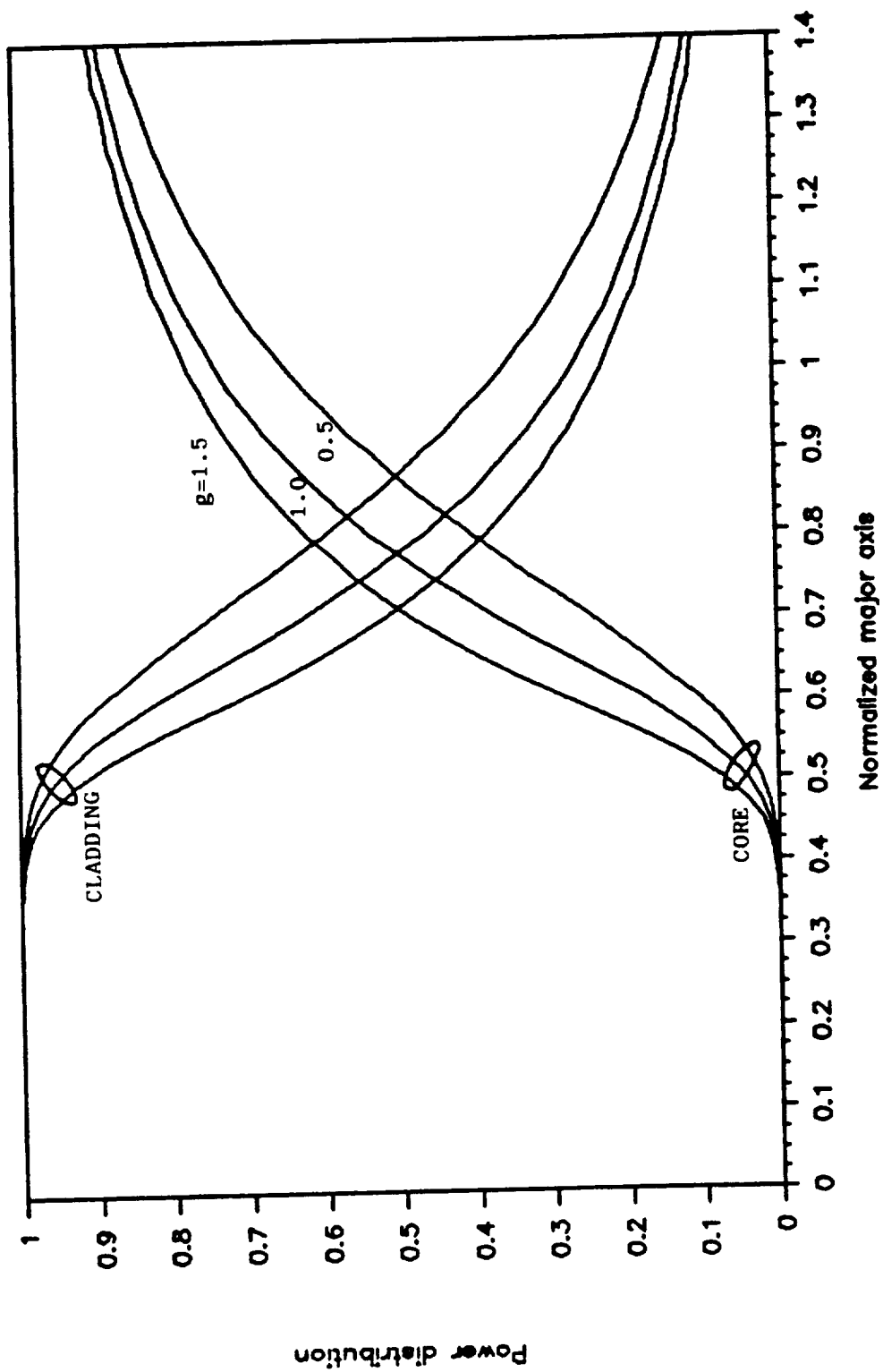


Figure 11. Power distribution characteristics for elliptical fiber as a function of normalized major axis for odd modes.

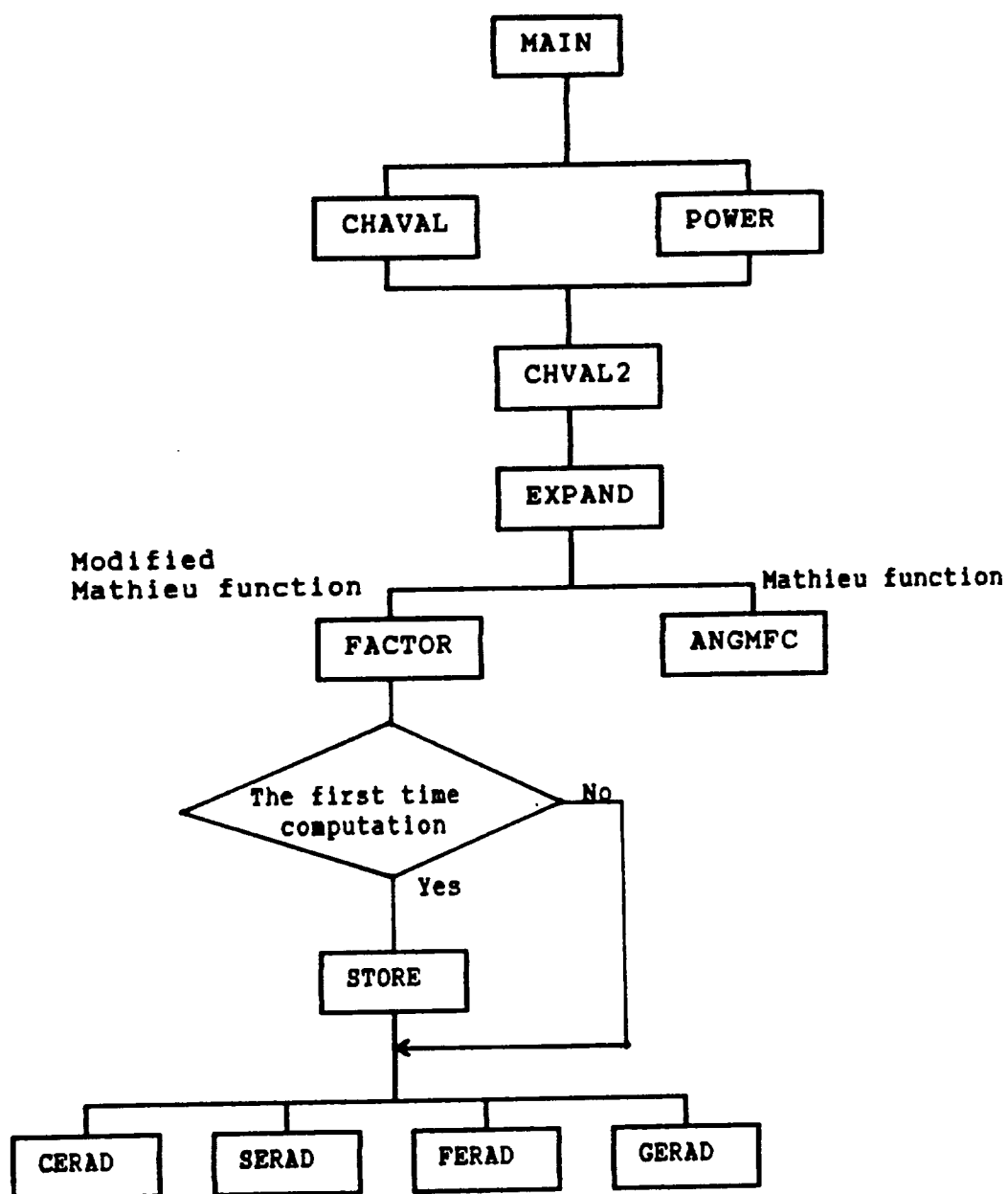
## 7. DESCRIPTION OF COMPUTER PROGRAMS

In this chapter, the computer programs for the anisotropic elliptical fiber will be considered. These programs are written in the language of FORTRAN IV and in order to present the complete computer programs the subroutines to calculate the Mathieu and modified Mathieu functions developed by Rengarajan and Lewis[49] will be included. The theory and notations used in the computer programs are the same as those employed by McLachlan[48].

The normalized propagation constant and power distribution characteristics as a function of the normalized cross section area or major axis for the given value of  $\epsilon$  and anisotropy have been determined by utilizing these programs. These computer programs consist of a main program and user called subroutines: CHAVAL and POWER. These subroutines CHAVAL and POWER call nine subroutines in order to compute the Mathieu and modified Mathieu functions.

In subroutine CHAVAL, an initial guess for the given mode is chosen and used to evaluate either Eq.(4.7) or Eq.(4.12). Next, Muller's method is used iteratively to determine the normalized wavelength that will minimize the function; an error criterion has been used to terminate the iteration. In subroutine POWER, the power distribution characteristics are calculated using the normalized wavelength obtained in subroutine CHAVAL. The algorithm was run on an CYBER 990 using only a single processor.

The sequence of the called subroutines is illustrated in the following flow chart.



## 8. CONCLUSION

The exact characteristic equation for anisotropic elliptical optical fibers is obtained for the odd and even hybrid modes in terms of infinite determinants employing Mathieu and modified Mathieu functions. The exact characteristic equation is applicable to elliptical fibers with any ellipticity. A simplified characteristic equation can then be obtained by applying the weakly guiding approximation such that the difference in the refractive indices of the core and the cladding is small. Under this approximation, it can be shown that significant simplification can be achieved.

The simplified characteristic equation is used to compute the normalized wavelength for an anisotropic elliptical fiber. When the anisotropy parameter is equal to unity, the characteristic equation becomes that of isotropic fiber. The results are compared to the previous research and they are in close agreement. For a fixed value of the normalized cross-section area or major axis, the normalized wavelength  $\lambda/\lambda_0$  is small for larger the value of anisotropy. This condition indicates that more energy is carried inside of the fiber. However, the geometry and anisotropy of fiber have a smaller effect when the normalized cross-section area or major axis is very small or very large.

An exact solution for the wave equation can not be determined when the thermoelastic stress causes a transverse anisotropy over the core of fibers. One possibility is that the propagation characteristics in the biaxial anisotropic fibers could be obtained by applying the numerical or approximation techniques given in Chapter 1 and this could be a subject for further study.

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**APPENDIX**

The following is a listing of the computer programs, MAIN, CHAVAL, POWER, CHVAL2, EXPAND, ANGMFC, FACTOR, STORE, CERAD, SERAD, FERAD AND GERAD written in FORTRAN IV language.

```

C
C   THIS PROGRAM CALCULATES THE NORMALIZED WAVELENGTH AND
C   POWER DISTRIBUTION FOR ELLIPTICAL FIBER AS FUNCTION OF
C   NORMALIZED CROSS-SECTION AREA OR MAJOR AXIS.
C
C   PROGRAM MAIN
C   IMPLICIT DOUBLE PRECISION (A-H,O-Z)
C   DIMENSION RES(56), X(56)
C   OPEN (6,FILE='OUTPUT')
C
C   NEVOD : = 1 FOR ODD MODE,
C           = 2 FOR EVEN MODE.
C   KASE  : = 1 FOR NORMALIZED CROSS-SECTIONAL AREA,
C           = 2 FOR NORMALIZED MAJOR AXIS.
C   ETA   : INDEPENDENT VARIABLE IN MATHIEU FUNCTION
C   PHI   : INDEPENDENT VARIABLE IN MODIFIED MATHIEU FUNCTION
C   MODE  : WAVE MODE NUMBER
C   P     : EP1/EPO
C   G     : ANISOTROPY EPZ/EPX
C
C
C   NEVOD=1
C   KASE=2
C   MODE=1
C   P=2.500
C   G=1.500
C   PHI=1.000
C   GO TO (11,12), NEVOD
11  WRITE(6,101) MODE
C   GO TO 13
12  WRITE(6,105) MODE
13  WRITE(6,102) P
C   WRITE(6,103) G
C   WRITE(6,104) PHI
C
C   CALCULATE THE NORMALIZED WAVELENGTH
C
C   CALL CHARVAL (NEVOD,P,G,MODE,PHI,KASE,RES)
C
C   DO 14 I=1,56
14  X(I)=RES(I)
C
C   CALCULATE POWER DISTRIBUTION FOR THE GIVEN MODE
C
C   CALL POWER (NEVOD,P,G,MODE,PHI,KASE,X)
C
101  FORMAT('1', 'THIS IS RESULT FOR ODD MODE, M =', I2)
102  FORMAT(1X, 'RATIO OF CORE AND CLADDING PERMITTIVITY =',
*      D12.5)
103  FORMAT(1X, 'ANISOTROPY =', D12.5)
104  FORMAT(1X, 'VARIABLE IN MATHIEU FUNCTION =', D12.5)

```

```

105  FORMAT('1', 'THIS IS THE RESULT FOR EVEN MODE, M=', I2)
      STOP
      END
      SUBROUTINE CHARVAL (NEVOD,P,G,MODE,PHI,KASE,RES)

C
C
C      PURPOSE : CALCULATE THE NORMALIZED WAVELENGTH FOR
C                  ELLIPTICAL GUIDE.
C
C      INPUT   : NEVOD -(INTEGER) SPECIFIES = 1 FOR ODD MODE
C                                     = 2 FOR EVEN MODE
C
C                  P -(DOUBLE PRECISION) IS THE RATIO BETWEEN CORE
C                  AND CLADDING PERMEABILITY.
C
C                  G -(DOUBLE PRECISION) IS THE ANISOTROPY, EZ/EX.
C
C                  MODE -(INTEGER) IS THE MODE OF CHARACTERISTIC
C                  EQUATION.
C
C                  PHI -(DOUBLE PRECISION) IS INDEPENDENT VARIABLE
C                  IN MODIFIED MATHIEU FUNCTION.
C
C                  KASE -(INTEGER) = 1 FOR NORMALIZED CROSS-SECTION
C                                     AREA,
C                                     = 2 FOR NORMALIZED MAJOR AXIS.
C
C      OUTPUT  : RES -(DOUBLE PRECISION) CONTAINS THE NORMALIZED
C                  WAVELENGTH.
C
C
C      IMPLICIT DOUBLE PRECISION (A-H,O-Z)
C      DIMENSION CHV1(23), CHV2(23), AB(25), QV(4), ETA(21),
C      *          SE1(21),SE1D(21),SEO(21),CE1(21),CE1D(21),
C      *          CEC(21),CE01(21),SEC1(21),CEDSE(21),SEDCE(21),
C      *          SE1SQ(21),CE1SQ(21),CE0M(21),SE0M(21),CE1M(21),
C      *          SE1M(21),RES(56)
C      DATA PI/3.141592653589793DC/

C
C      ETA : INDEPENDENT VARIABLE IN MATHIEU FUNCTION
C      XI  : INDEPENDENT VARIABLE IN MODIFIED MATHIEU FUNCTION
C      IORDER : WAVE MODE NUMBER
C      P   : EP1/EP0
C      G   : ANISOTROPY EP2/EPX
C
C      M=20
C      LAST=40
C      IF(KASE.EQ.2) LAST=56
C      XI=PHI
C      DT=DTANH(XI)
C      DCOS2=DCOSH(XI)*DCOSH(XI)
C      X2=PI*DT*DCOS2
C      WRITE(6,210) X2
210  FORMAT(1X, D12.5)
C
C      SUBDIVIDE ETA FOR INTEGRATION, ALFA, BETA, GAMMA AND ANU.
C
C      H=2.000*PI/20.000
C      ETA(1)=0.000
C      DO 11 I=2,21

```

```

11  ETA(1)=ETA(1)+(1-1)*H
C
C  CSA - NORMALIZED CROSS SECTION AREA OR MAJOR AXIS
C
      CSA=0.000
      XINC=2.5D-2
      VARX=0.99999999900
      DO 70 K=1, LAST
      CSA=CSA+XINC
      GO TO (12,13), KASE
12  WRITE(6,201) CSA
201  FORMAT(1X, 'NORMALIZED CROSS SECTION AREA = ', D12.5)
C
      X1=(PI*PI*CSA)/(4.000*DT*DCOS2)
      GO TO 14
13  WRITE(6,202) CSA
202  FORMAT(1X, 'NORMALIZED MAJOR AXIS = ', D12.5)
C
      X1=(PI*PI*CSA**2)/(4.000*DCOS2)
C
C  FIRST GUESS OF VAR = LANDA/LANDAO
C
14  IF(K.EQ.1) VAR=VARX
      NC=1
      NS=0
10  VAR2=1.0D0/(VAR*VAR)
C
C  EVALUATE Q : INDEPENDANT VARIABLE
C  GAMMAE = QV(1), GAMMAH = QV(2), GAMMA2 = QV(3)
C
      QV(1)=X1*(P*G-G*VAR2)
      QV(2)=X1*(P-VAR2)
      QV(3)=X1*(1.000-VAR2)
      QV(4)=X1*(1.000-VAR2)
C
C  CALCULATE MATHIEU AND MODIFIED MATHIEU FUNCTIONS
C
      DO 50 KQ=1,4
      GO TO (15,16), NEVOD
15  IEVOD=1
      IF(MOD(KQ,2).EQ.1) IEVOD=2
      IORDER=MODE
      IF(MOD(KQ,2).EQ.1) IORDER=MODE
      Q=QV(KQ)
      CALL CHVAL2(M,Q,CHV1,CHV2,J)
      CV=CHV1(IORDER)
      IF(MOD(KQ,2).EQ.1) CV=CHV2(IORDER+1)
      GO TO 17
16  IEVOD=1
      IF(MOD(KQ,2).EQ.0) IEVOD=2
      IORDER=MODE

```

```

IF(MOD(KQ,2).EQ.0) IORDER=MODE
Q=QV(KQ)
CALL CHVAL2(M,Q,CHV1,CHV2,J)
CV=CHV1(IORDER)
IF(MOD(KQ,2).EQ.0) CV=CHV2(IORDER+1)

C
C   OBTAIN EXPANDING COEFFICIENT, ABXX
C
C   CALL EXPAND(Q,IEVOD,IORDER,CV,3,AB,N)
17
C
C   CALCULATE MATHIEU FUNCTIONS AND DERIVATIVES,
C   ORDER = MODE
C
KQEO=KQ
IF(MOD(NEVOD,2).EQ.1) KQEO=KQ+4
GO TO(21,22,23,24,22,21,24,23), KQEO
21 DO 41 I=1,21
   SE1(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),C,AB,N)
41 SE1D(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),1,AB,N)
   GO TO 45
23 DO 42 I=1,21
   SE0(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),0,AB,N)
42 GO TO 45
22 DO 43 I=1,21
   CE1(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),0,AB,N)
43 CE1D(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),1,AB,N)
   GO TO 45
24 DO 44 I=1,21
   CE0(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),0,AB,N)
44
C   NORMALIZATION FACTOR FOR MODIFIED MATHIEU FUNCTION
C
C   CALL FACTOR(IEVOD,IORDER,Q,AB,N,PS)
45
C   COMPUTE AND STORE THE VALUES OF BESSEL FUNCTIONS
C
C   CALL STORE(Q,XI,N)
C
C   CALCULATE MODIFIED MATHIEU FUNCTIONS
C
GO TO (51,52,53,54,52,51,54,53),KQEO
51 SE=SERAD(Q,IORDER,0,PS,AB,N)
   SED=SERAD(Q,IORDER,1,PS,AB,N)
   GO TO 50
53 GE=GERAD(Q,IORDER,0,PS,AB,N)
   GED=GERAD(Q,IORDER,1,PS,AB,N)
   GO TO 50
52 CE=CERAD(Q,IORDER,0,PS,AB,N)
   CED=CERAD(Q,IORDER,1,PS,AB,N)
   GO TO 50
54 FE=FERAD(Q,IORDER,0,PS,AB,N)

```

```

FED=FERAD(Q,IORDER,1,PS,AB,N)
50 CONTINUE
C
C CALCULATE M TH TERM( = MODE) OF
C ALFA, BETA, GAMMA AND NU.
C
WRITE(6,107) SE,SED,GE,GED,CE,CED,FE,FED
DO 56 I=1,21
CE0M(I)=CE0(I)
SE0M(I)=SE0(I)
CE1M(I)=CE1(I)
SE1M(I)=SE1(I)
CE01(I)=CE0(I)*CE1(I)
SE01(I)=SE0(I)*SE1(I)
SE1SQ(I)=SE1(I)*SE1(I)
CE1SQ(I)=CE1(I)*CE1(I)
CEDSE(I)=CED(I)*SE1(I)
56 SEDCE(I)=SED(I)*CE1(I)
S1=SIMPSN(CE01,20,H)
S2=SIMPSN(SE01,20,H)
S3=SIMPSN(CEDSE,20,H)
S4=SIMPSN(SEDCE,20,H)
S5=SIMPSN(SE1SQ,20,H)
S6=SIMPSN(CE1SQ,20,H)
S5M=S5
S6M=S6
GO TO (57,58), NEVDD
57 ALFAM=S2/S5
BETAM=S1/S6
GAMMAM=S4/S6
ANUM=S3/S5
GO TO 59
58 ALFAM=S1/S6
BETAM=S2/S5
GAMMAM=S3/S6
ANUM=S4/S6
59 XMSQD=ALFAM*BETAM
RHC=QV(2)/QV(3)
WRITE(6,103) ALFAM,BETAM,GAMMAM,ANUM,RHC
103 FORMAT(1X,5D12.5)
C
C CALCULATE MATHIEU FUNCTION INTEGRALS
C ORDER = N, N+2 - -
C
ALFA=0.000
BETA=0.000
GAMMA=0.000
ANU=0.000
XMSQN1=0.000
XMSQN2=0.000
DO 90 IM=1,4

```

```

IORDER=2*IM-1
IF(MOD(MODE,2).EQ.0) IORDER=2*IM-2
IF(IORDER.NE.MODE) GO TO 61
ALFA=ALFAM
BETA=BETAM
GAMMA=GAMMAM
ANU=ANUM
XMSQN1=XMSQN1+BETA*ANU
XMSQN2=XMSQN2+ALFA*GAMMA
XMSQ=(XMSQN1*XMSQN2)/XMSQO
GO TO 92
61 DO 80 KQ=1,4
GO TO (62,63), NEVOD
62 IEVOD=1
IF(MOD(KQ,2).EQ.1) IEVOD=2
Q=QV(KQ)
CALL CHVAL2(M,Q,CHV1,CHV2,J)
CV=CHV1(IORDER)
IF(MOD(KQ,2).EQ.1) CV=CHV2(IORDER+1)
GO TO 64
63 IEVOD=1
IF(MOD(KQ,2).EQ.0) IEVOD=2
Q=QV(KQ)
CALL CHVAL2(M,Q,CHV1,CHV2,J)
CV=CHV1(IORDER)
IF(MOD(KQ,2).EQ.0) CV=CHV2(IORDER+1)
C
C OBTAIN EXPANDING COEFFICIENT, ABXX
C
64 CALL EXPAND(Q,IEVOD,IORDER,CV,3,AB,N)
C
C CALCULATE MATHIEU FUNCTIONS AND DERIVATIVES
C
KQEO=KQ
IF(MOD(NEVOD,2).EQ.1) KQEO=KQ+4
GO TO(71,72,73,74,72,71,74,73), KQEO
71 DO 81 I=1,21
SE1(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),0,AB,N)
81 SE1D(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),1,AB,N)
GO TO 80
73 DO 82 I=1,21
82 SE0(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),0,AB,N)
GO TO 80
72 DO 83 I=1,21
83 CE1(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),0,AB,N)
CE1D(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),1,AB,N)
GO TO 80
74 DO 84 I=1,21
84 CE0(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),0,AB,N)
80 CONTINUE
C

```



```

C      CALCULATE SUM (BETA * NU) AND SUM (ALFA * GAMMA)
C
      DO 86 I=1,21
      CE01(I)=CE0M(I)*CE1(I)
      SE01(I)=SE0M(I)*SE1(I)
      SE1SQ(I)=SE1(I)*SE1(I)
      CE1SQ(I)=CE1(I)*CE1(I)
      CEDSE(I)=CE1D(I)*SE1M(I)
86     SEDCE(I)=SE1D(I)*CE1M(I)
      S1=SIMPSN(CE01,20,H)
      S2=SIMPSN(SE01,20,H)
      S3=SIMPSN(CEDSE,20,H)
      S4=SIMPSN(SEDCE,20,H)
      S5=SIMPSN(SE1SQ,20,H)
      S6=SIMPSN(CE1SQ,20,H)
C
      GO TO (87,88), NEVOD
87     ALFA=S2/S5
      BETA=S1/S6
      GAMMA=S4/S6M
      ANU=S3/S5M
      GO TO 89
88     ALFA=S1/S6
      BETA=S2/S5
      GAMMA=S3/S5M
      ANU=S4/S6M
89     XMSQN1=XMSQN1+BETA*ANU
      XMSQN2=XMSQN2+ALFA*GAMMA
      XMSQ=(XMSQN1*XMSQN2)/XMSQD
92     WRITE(6,103) ALFA,BETA,GAMMA,ANU,XMSQ
9C    CONTINUE
C
C      EVALUATE CHARACTERISTIC EQUATION
C
      Y1=-(XMSQ*(1.0D0-RHC)**2)
      Y2=VAR*VAR
      GO TO(93,94), NEVOD
93     Y3=(SED/SE)-(RHC*GED/GE)
      Y4=(P*CED/CE)-(RHC*FED/FE)
      WRITE(6,101) Y1, Y2, Y3, Y4
      GO TO 95
94     Y3=(CED/CE)-(RHC*FED/FE)
      Y4=(P*SED/SE)-(RHC*GED/GE)
      WRITE(6,101) Y1, Y2, Y3, Y4
95     YX=Y2*Y3*Y4
      YZ=YX/Y1
      Y=YZ-1.0D0
      WRITE(6,203) YX, Y1, YZ, Y, VAR
203    FORMAT(1X,5D12.5)
C
C      DESIDE ON TOLERANCES

```

```

C
  IF(NC.NE.1.AND.DABS(Y).LE.2.0D-3) GO TO 39
  IF(NC.EQ.1) GO TO 32
  IF(NC.EQ.2) GO TO 34
  IF(NS.EQ.0) GO TO 34
  IF(Y*YS1) 36,36,31
31  YS1=Y
  VARS1=VAR
  VAR=(VARS1+VARS2)/2.0D0
  NS=NS+1
  IF(NS.LE.20) GO TO 10
  GO TO 47

C
C  1 ST CALCULATION OF Y, DECREMENT VAR BY 0.C1
C
32  YS1=Y
  VARS1=VAR
33  VAR=VAR-1.0D-2
  NC=NC+1
  YMIN=Y
  VARMIN=VAR
  IF(NC.LE.20) GO TO 10
  GO TO 48
34  IF(Y*YS1) 36,36,35
35  IF(DABS(Y)-DABS(YS1)) 32,33,33
36  YS2=Y
  VARS2=VAR
  VAR=(VARS1+VARS2)/2.0D0
  NS=NS+1
  YMIN=Y
  VARMIN=VAR
  IF(NS.LE.20) GO TO 10
47  WRITE(6,106)
  WRITE(6,108) YMIN, VARMIN
  GO TO 39
48  WRITE(6,102)
  VAR=VARX
  WRITE(6,108) YMIN, VARMIN
39  RES(K)=VAR
70  CONTINUE
  RETURN
101  FORMAT(1X,4D12.5)
102  FORMAT(1X,'ERROR : RESULT HAS SAME SIGN FOR 10 TRIES')
106  FORMAT(1X,'ERROR : 10 TRY FAILED TO OBTAIN RESOLUTION')
107  FORMAT(1X,8D12.5)
108  FORMAT(1X, 'OBTAINED RESOLUTION =', D12.5,
C      'MIN CALCULATED NORMALIZED WAVELENGTH =', D12.5)
  END
SUBROUTINE POWER (NEVOD,P,G,MODE,BOUND,KASE,A)

C
C  PURPOSE : CALCULATE POWER DISTRIBUTION ON ELLIPTICAL

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```

C      GUIDE.
C      INPUT   : NEVOD -(INTEGER) SPECIFIES = 1 FOR ODD MODE
C               = 2 FOR EVEN MODE
C               P -(DOUBLE PRECISION) IS THE RATIO BETWEEN CORE
C               AND CLADDING PERMEABILITY.
C               G -(DOUBLE PRECISION) IS THE ANISOTROPY, EZ/EX.
C               MODE -(INTEGER) IS THE MODE OF CHARACTERISTIC
C               EQUATION.
C               PHI -(DOUBLE PRECISION) IS INDEPENDENT VARIABLE
C               IN MODIFIED MATHIEU FUNCTION.
C               A -(DOUBLE PRECISION) IS THE NORMALIZED
C               WAVELENGTH.
C               KASE -(INTEGER) = 1 FOR NORMALIZED CROSS-SECTION
C               AREA,
C               = 2 FOR NORMALIZED MAJOR AXIS.
C      OUTPUT  : RES -(DOUBLE PRECISION) CONTAINS THE RATIO OF
C               POWER DISTRIBUTION.
C
C      IMPLICIT DOUBLE PRECISION (A-H,O-Z)
C      DIMENSION CHV1(23), CHV2(23), AB(25), QV(4), A(56),
C      ETA(21), PHI(41), SE1(21), SE1D(21), SE0(21),
C      SE0D(21), CE1(21), CE1D(21), CE0(21), CE0D(21),
C      SE(21), SED(21), CE(21), CED(21),
C      FE(41), FED(41), GE(41), GED(41),
C      S1(21), S2(21), S3(21), S4(21), S5(21), S6(21),
C      S7(21), S8(21), S9(21), S10(21), S11(21),
C      S12(21), S13(21), S14(21),
C      S21(21), S22(21), S23(21)
C      DATA PI/3.141592653589793D0/
C
C      ETA : INDEPENDENT VARIABLE IN MATHIEU FUNCTION
C      XI : INDEPENDENT VARIABLE IN MODIFIED MATHIEU FUNCTION
C      IORDER : WAVE MODE NUMBER
C      P : EP1/EPO
C      G : ANISOTROPY EP2/EPX
C
C      M=20
C      EP=(1.0D-9)/(36.0D0*PI)
C      XNU=4.0D0*PI*1.0D-7
C      CONS=DSQRT(EP/XNU)
C      LAST=40
C      IF(KASE.EQ.2) LAST=56
C      PHI0=0.0D0
C      PHI1=90D0
C      PHI2=5.0D0*90D0
C      DT=DTANH(PHI1)
C      DCOS2=DCOSH(PHI1)*DCOSH(PHI1)
C      X2=PI*DT*DCOS2
C      WRITE(6,210) X2
C
C      SUBDIVIDE ETA AND PHI FOR INTEGRATION.

```

```

C      H1=2.000*PI/20.000
      ETA(1)=0.000
      DO 11 I=2,21
11     ETA(I)=ETA(I-1)+H1
      H2=BOUND/20.000
      PHI(1)=0.000
      DO 12 I=2,21
12     PHI(I)=PHI(I-1)+H2
      H3=(PHI2-PHI1)/20.000
      DO 13 I=22,41
13     PHI(I)=PHI(I-1)+H3
C
C      CSA - NORMALIZED CROSS SECTION AREA OR MAJOR AXIS
C
C      CSA=0.000
      XINC=2.50-2
      DO 70 K=1, LAST
      CSA=CSA+XINC
      IF(A(K).EQ.1.000) GO TO 81.
      GO TO (14,15), KASE
14     WRITE(6,201) CSA
201    FORMAT(1X,'NORMALIZED CROSS SECTION AREA = ',D12.5)
C
      X1=(PI*PI*CSA)/(4.000*DT*DCOS2)
      GO TO 16
15     WRITE(6,202) CSA
202    FORMAT(1X,'NORMALIZED MAJOR AXIS = ',D12.5)
C
      X1=(PI*PI*CSA**2)/(4.000*DCOS2)
C
C      CALCULATE CONSTANTS.
C
16     VAR2=1.000/(A(K)*A(K))
C
C      EVALUATE Q : INDEPENDANT VARIABLE
C      GAMMAE = QV(1), GAMMAH = QV(2), GAMMA2 = QV(3)
C
      QV(1)=X1*(P*G-G*VAR2)
      QV(2)=X1*(P-VAR2)
      QV(3)=X1*(1.000-VAR2)
      QV(4)=X1*(1.000-VAR2)
C
      C4=CONS/A(K)
      C1=P*C4
      C2=1.000/(A(K)*CONS)
      C3=P+VAR2
      C5=1.000+VAR2
      C6=(QV(2)*QV(2))/(QV(3)*QV(3))
C
C      CALCULATE MATHIEU AND MODIFIED MATHIEU FUNCTIONS

```

```

C
DO 50 KQ=1,4
GO TO (17,18), NEVOD
17 IEVOD=1
IF(MOD(KQ,2).EQ.1) IEVOD=2
IORDER=MODE
IF(MOD(KQ,2).EQ.1) IORDER=MODE
Q=QV(KQ)
WRITE(6,301) X1,Q
CALL CHVAL2(M,Q,CHV1,CHV2,J)
CV=CHV1(IORDER)
IF(MOD(KQ,2).EQ.1) CV=CHV2(IORDER+1)
GO TO 19
18 IEVOD=1
IF(MOD(KQ,2).EQ.0) IEVOD=2
IORDER=MODE
IF(MOD(KQ,2).EQ.0) IORDER=MODE
Q=QV(KQ)
WRITE(6,301) X1,Q
CALL CHVAL2(M,Q,CHV1,CHV2,J)
CV=CHV1(IORDER)
IF(MOD(KQ,2).EQ.0) CV=CHV2(IORDER+1)
C
C OBTAIN EXPANDING COEFFICIENT, ABXX
C
19 CALL EXPAND(Q,IEVOD,IORDER,CV,3,AB,N)
C
C CALCULATE MATHIEU FUNCTIONS AND DERIVATIVES,
C ORDER = MODE
C
KQEO=KQ
IF(MOD(NEVOD,2).EQ.1) KQEO=KQ+4
GO TO(21,22,23,24,22,21,24,23), KQEO
21 DO 41 I=1,21
SE1(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),0,AB,N)
41 SE1D(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),1,AB,N)
GO TO 45
23 DO 42 I=1,21
SE0D(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),1,AB,N)
42 SE0(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),0,AB,N)
GO TO 45
22 DO 43 I=1,21
CE1(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),0,AB,N)
43 CE1D(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),1,AB,N)
GO TO 45
24 DO 44 I=1,21
CE0D(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),1,AB,N)
44 CE0(I)=ANGMFC(Q,IEVOD,IORDER,ETA(I),0,AB,N)
C
C NORMALIZATION FACTOR FOR MODIFIED MATHIEU FUNCTION
C

```

```

45  CALL FACTOR(IEVOD,IORDER,Q,AB,N,PS)
C
C  CALCULATE MODIFIED MATHIEU FUNCTIONS
C
GO TO (31,32,33,34,32,31,34,33),KQEO
31  DO 51 I=1,21
    CALL STORE(Q,PHI(I),N)
    SE(I)=SERAD(Q,IORDER,0,PS,AB,N)
51  SED(I)=SERAD(Q,IORDER,1,PS,AB,N)
    GO TO 50
33  DO 52 I=21,41
    CALL STORE(Q,PHI(I),N)
    GE(I)=GERAD(Q,IORDER,0,PS,AB,N)
52  GED(I)=GERAD(Q,IORDER,1,PS,AB,N)
    GO TO 50
32  DO 53 I=1,21
    CALL STORE(Q,PHI(I),N)
    CE(I)=CERAD(Q,IORDER,0,PS,AB,N)
53  CED(I)=CERAD(Q,IORDER,1,PS,AB,N)
    GO TO 50
34  DO 54 I=21,41
    CALL STORE(Q,PHI(I),N)
    FE(I)=FERAD(Q,IORDER,0,PS,AB,N)
54  FED(I)=FERAD(Q,IORDER,1,PS,AB,N)
50  CONTINUE
C
C  CALCULATE INTEGRAND
C
DO 56 I=1,21
S21(I)=SE0(I)*SE1(I)
S22(I)=SE1(I)*SE1(I)
S23(I)=CE1D(I)*SE1(I)
S1(I)=SE1D(I)*SE1D(I)
S2(I)=CE1D(I)*CE1D(I)
S3(I)=CE1D(I)*SE1(I)
S4(I)=SE0D(I)*SE0D(I)
S5(I)=CE0D(I)*CE0D(I)
56 S6(I)=CE0D(I)*SE0(I)
C
C
DO 57 I=1,21
S7(I)=SED(I)*SED(I)
S8(I)=SE(I)*SE(I)
S9(I)=CED(I)*CED(I)
57 S10(I)=CE(I)*CE(I)
C
C
DO 58 I=21,41
I1=I-10
S11(I1)=GED(I)*GED(I)
S12(I1)=GE(I)*GE(I)

```

```

S13(11)=FED(1)*FED(1)
S14(11)=FE(1)*FE(1)
58
C
C
C
PERFORM THE INTEGRATION

ST1=SIMPSN(S21,20,H1)
ST2=SIMPSN(S22,20,H1)
ST3=SIMPSN(S23,20,H1)
T11=PI*SIMPSN(S7,10,H2)
T12=SIMPSN(S1,20,H1)*SIMPSN(S8,20,H2)
T21=PI*SIMPSN(S9,20,H2)
T22=SIMPSN(S2,20,H1)*SIMPSN(S10,20,H2)
T31=SIMPSN(S3,20,H1)
T32=(CE(21)*SE(21)-CE(1)*SE(1))
T41=PI*SIMPSN(S11,20,H3)
T42=SIMPSN(S4,20,H1)*SIMPSN(S12,20,H3)
T51=PI*SIMPSN(S13,20,H3)
T52=SIMPSN(S5,20,H1)*SIMPSN(S14,20,H3)
T61=SIMPSN(S6,20,H1)
T62=(FE(41)*GE(41)-FE(21)*GE(21))

C
C
C
CALCULATE THE ARBITRARY CONSTANTS.

A11=(1.000-QV(2)/QV(3))*FE(21)*(ST3/ST2)
A12=CONS*A(K)*ST1/ST2
A13=((P*FE(21)*CED(21))/CE(21)-QV(2)*FED(21)/QV(3))
A21=A11/(A12*A13)
A1=A21*GE(21)/SE(21)
A1=A1*A1
B1=FE(21)*FE(21)
BA1=A21*FE(21)*GE(21)/SE(21)
A2=A21*A21
BA2=A21
T1=C1*A1*(T11+T12)
T2=C2*B1*(T21+T22)
T3=C3*BA1*T31*T32
WRITE(6,302) T1,T2,T3
PCOR=T1+T2-T3
T4=C4*A2*(T41+T42)
T5=C2*1.000*(T51+T52)
T6=C5*BA2*T61*T62
WRITE(6,302) T4,T5,T6
PCLAD=C6*(T4+T5-T6)
RCOR=PCOR/(PCOR+PCLAD)
RCLAD=1.000-RCOR
WRITE(6,211) RCOR, RCLAD, PCOR, PCLAD
GO TO 70
81
WRITE(6,212)
70
CONTINUE
210
FORMAT(1X, D12.5)
211
FORMAT(11X, 4D12.5)

```

```

212  FORMAT(1X,'NORMALIZED WAVELENGTH = 1, NO POWER CALCULATION')
301  FORMAT(1X,2D12.5)
302  FORMAT(1X,3D12.5)
      RETURN
      END
      DOUBLE PRECISION FUNCTION SIMPSN(Q,N,H)
      DOUBLE PRECISION Q(21),H

C
C      INTEGRATION BY SIMPSON'S RULE
C
      SIMPSN=Q(1)+4.0D0*Q(2)+Q(N+1)
      DO 1 I=4,N,2
1      SIMPSN=SIMPSN+2.0D0*Q(I-1)+4.0D0*Q(I)
      SIMPSN=SIMPSN*H/3.0D0
      RETURN
      END
      SUBROUTINE CHVAL2(N,QQ,CHV1,CHV2,J)
C      PURPOSE:   TO COMPUTE THE CHARACTERISTIC VALUES OF ODD
C                AND EVEN MATHIEU FUNCTIONS OF POSITIVE OR
C                NEGATIVE 'Q'
C      INPUT:     N-(INTEGER) SPECIFIES THAT CH. VALUES BE
C                OBTAINED FOR ORDERS 0 THRU N-1 FOR EVEN
C                FUNCTIONS AND FOR ORDERS 1 THRU N-1 FOR
C                ODD FUNCTIONS
C                QQ-(DOUBLE PRECISION) THE PARAMETER 'Q' IN
C                MATHIEU'S DIFFERENTIAL EQUATION
C      OUTPUT:    CHV1-(DOUBLE PRECISION) AN ARRAY OF LENGTH N
C                CONTAINING CH. VALUES OF ODD MATHIEU FUNCTIONS
C                OF ORDERS 1 THRU N-1.
C                CHV1(N) IS A DUMMY VARIABLE.
C                CHV2-(DOUBLE PRECISION) AN ARRAY OF LENGTH N
C                CONTAINING CH. VALUES OF EVEN MATHIEU FUNCTIONS
C                OF ORDERS 0 THRU N-1.
C                J-(INTEGER) MAXIMUM ORDER UPTO WHICH CH. VALUES
C                HAVE BEEN SUCCESSFULLY COMPUTED
      DOUBLE PRECISION CV1(6,25),CV2(6,25),CHV1(N),CHV2(N),QQ
      DOUBLE PRECISION QABS,DABS
131  FORMAT('C',5X,'NOT ALL CH. VALUES AVAILABLE----WARNING')
      IF(QQ.LT.0.0D0) GO TO 20
      IF(N.GT.1) CALL MFCVAL(N-1,N-1,QQ,CV1,J1)
      CALL MFCVAL(N,N-1,QQ,CV2,J2)
      IF(J1.LT.(N-1).OR.J2.LT.N) WRITE(6,101)
      J=MIN0(J1,J2-1)
      DO 10 I=1,J
          IF(N.GT.1) CHV1(I) =CV1(1,I)
          CHV2(I)=CV2(1,I)
10  CONTINUE
      CHV2(J+1)=CV2(1,J+1)
      RETURN
20  QABS=DABS(QQ)
      CALL MFCVAL(N,N-1,QABS,CV2,J2)

```



```

IF(N.NE.1) GO TO 25
CHV2(1)=CV2(1,1)
RETURN
25 CALL MFCVAL(N-1,N-1,QABS,CV1,J1)
IF(J1.NE.(N-1).OR.J2.NE.N) WRITE(6,101)
J=MIN0(J1,J2-1)
DO 30 I=1,J,2
    CHV2(I)=CV2(1,I)
    CHV1(I)=CV2(1,I+1)
    CHV2(I+1)=CV1(1,I)
    IF((I+1).LE.J)CHV1(I+1)=CV1(1,I+1)
30 CONTINUE
IF(MOD((N-1),2).EQ.0)CHV2(N)=CV2(1,N)
RETURN
END
SUBROUTINE MFCVAL(N,R,QQ,CV,J)
C *****
INTEGER J,K,KK,L,M,N,R,TYPE
DOUBLE PRECISION A,CV,DL,DR,DTM,Q,QQ,T,TM,TOL,TOLA
DOUBLE PRECISION FILL(3)
DIMENSION CV(6,N)
EQUIVALENCE (DL,DR,T)
COMMON/MF1/Q,TOL,TYPE,DUMMY(4)
COMMON/MF2/FILL
TOL=1.0D-13
IF(N-R) 10,10,20
10 L=1
GO TO 30
20 L=2
30 Q=QQ
DO 500 K=1,N
    J=K
    IF(Q) 960,490,40
40 KK=MIN0(K,4)
TYPE=2*MOD(L,2)+MOD(K-L+1,2)
C FIRST APPROXIMATION
GO TO(100,200,300,400),KK
100 IF(Q-1.0D0)110,140,140
110 GO TO(120,130),L
120 A=1.0D0-Q-.125D0*Q*Q
GO TO 420
130 A=Q*Q
A=A*(-.5D0+.0546875D0*A)
GO TO 420
140 IF(Q-2.0D0) 150,180,180
150 GO TO(160,170),L
160 A=1.033D0-1.0746D0*Q-.0068D0*Q*Q
GO TO 420
170 A=.23D0-.495D0*Q-.191D0*Q*Q
GO TO 420
180 A=-.25D0-2.0D0*Q+2.0D0*DSQRT(Q)

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```

      GO TO 420
200  DL=L
      IF(Q*DL-6.000) 210,350,350
210  GO TO(220,230),L
220  A=4.0152100-Q*(.04600+.066785700*Q)
      GO TO 420
230  A=1.000+1.0500700*Q-.18014300*Q*Q
      GO TO 420
300  IF(Q-8.000) 310,350,350
310  GO TO(320,330),L
320  A=8.9386700+.17815600*Q-.025213200*Q*Q
      GO TO 420
330  A=3.7001700+.95348500*Q-.047506500*Q*Q
      GO TO 420
350  DR=K-1
      A=CV(1,K-1)-DR+4.000*DSQRT(Q)
      GO TO 420
400  A=CV(1,K-1)-CV(1,K-2)
      A=3.000*A+CV(1,K-3)
420  IF(Q.GE.1.000) GO TO 440
      IF(K.NE.1) GO TO 430
425  TOLA=DMAX1(DMIN1(TOL,DABS(A)),1.00-14)
      GO TO 450
430  TOLA=TOL*DABS(A)
      GO TO 450
440  TOLA=TOL*DMAX1(Q,DABS(A))
445  TOLA=DMAX1(DMIN1(TOLA,DABS(A),.400*DSQRT(Q)),1.00-14)
C    CRUDE UPPER AND LOWER BOUNDS
450  CALL BOUNDS(K,A,TOLA,CV,N,M)
      IF(M.NE.0) IF(M-1) 470,910,900
C    ITERATE
      CALL MFITR8(TOLA,CV(1,K),CV(2,K),M)
      IF(M.GT.0) GO TO 920
C    FINAL BOUNDS AND FUNCTIONS, D
470  T=CV(1,K)-TOLA
      CALL TMOFA(T,TM,DTM,M)
      IF(M.GT.0) GO TO 940
      CV(3,K)=T
      CV(4,K)=-TM/DTM
480  T=CV(1,K)+TOLA
      CALL TMOFA(T,TM,DTM,M)
      IF(M.GT.0) GO TO 950
      CV(5,K)=T
      CV(6,K)=-TM/DTM
      GO TO 500
C    Q EQUALS ZERO
490  CV(1,K)=(K-L+1)**2
      CV(2,K)=0.000
      CV(3,K)=CV(1,K)
      CV(4,K)=0.000
      CV(5,K)=CV(1,K)

```

```

      CV(6,K)=0.000
500  CONTINUE
550  RETURN
C    PRINT ERROR MESSAGES
900  WRITE(6,901) K
901  FORMAT('0','CRUDE BOUNDS CANNOT',' BE LOCATED, NO OUTPUT',
C      ' FOR K=',I2)
      GO TO 930
910  WRITE(6,911) K
911  FORMAT('0','ERROR IN SUBPROGRAM TMOFA, VIA SUBPROGRAM
C      BOUNDS, NO OUTPUT, FOR K=',I2)
      GO TO 930
920  WRITE(6,921) K
921  FORMAT('0','ERROR IN SUBPROGRAM, TMOFA, VIA SUBPROGRAM,
C      MFITR8, NO OUTPUT, FOR K=',I2)
930  J=J-1
      GO TO 550
940  WRITE(6,941) K
941  FORMAT('0','ERROR IN SUBPROGRAM, TMOFA, NO LOWER BOUND,
C      FOR K=',I2)
      CV(3,K)=0.000
      CV(4,K)=0.00
      GO TO 490
950  WRITE(6,951) K
951  FORMAT('0','ERROR IN SUBPROGRAM, TMOFA, NO UPPER BOUND,
C      FOR K=',I2)
      CV(5,K)=0.00
      CV(6,K)=0.00
      GO TO 500
960  WRITE(6,961)
961  FORMAT(2CH0Q GIVEN NEGATIVELY,,20H USED ABSOLUTE VALUE)
      Q=-Q
      GO TO 40
      END
      SUBROUTINE BOUNDS(K,APPROX,TOLA,CV,N,MM)
      INTEGER K,KA,M,MM,N
      DOUBLE PRECISION A,APPROX,A0,A1,CV,DTM,D0,D1,Q,TH,TOLA
      DIMENSION CV(6,N)
      COMMON/MF1/Q,DUMMY(7)
      COMMON/MF2/A0,A,A1
      KA=0
      IF(K.EQ.1) GO TO 20
      IF(APPROX-CV(1,K-1)) 10,10,20
10    A0=CV(1,K-1)+1.000
      GO TO 30
20    A0=APPROX
30    CALL TMOFA(A0,TM,DTM,M)
      IF(M.GT.0) GO TO 250
      D0=-TM/DTM
      IF(D0) 100,300,50
C    A0 IS LOWER BOUND,

```

```

C      SEARCH FOR UPPER BOUND
50     A1=A0+D0+.100
      CALL TMOFA(A1, TM, DTM, M)
      IF(M.GT.0) GO TO 250
      D1=-TM/DTM
      IF(D1) 200,350,60
60     A0=A1
      D0=D1
      KA=KA+1
      IF(KA-4) 50,400,400
C      A1 IS UPPER BOUND, SEARCH FOR LOWER BOUND
100    A1=A0
      D1=D0
      A0=DMAX1(A1+D1-.100,-2.000*Q)
      IF(K.EQ.1) GO TO 110
      IF(A0-CV(1,K-1)) 150,150,110
110    CALL TMOFA(A0, TM, DTM, M)
      IF(M.GT.0) GO TO 250
      DC=-TM/DTM
      IF(D0) 120,300,200
120    KA=KA+1
      IF(KA-4) 100,400,400
150    KA=KA+1
      IF(KA-4) 160,400,400
160    AC=A1+DMAX1(TOLA,DABS(D1))
      GO TO 30
200    A=.500*(A0+DC+A1+D1)
      IF(A.LE.A0.OR.A.GE.A1) A=.500*(A0+A1)
250    MM=M
      RETURN
300    CV(1,K)=A0
310    CV(2,K)=0.00
      M=-1
      GO TO 250
350    CV(1,K)=A1
      GO TO 310
400    M=2
      GO TO 250
      END
      SUBROUTINE MFITRB(TOLA,CV,DCV,MM)
      INTEGER M,MM,N
      DOUBLE PRECISION A,A0,A1,A2,CV,D,DCV,DTM,TM,TOLA
      LOGICAL LAST
      COMMON/MF2/A0,A,A1
      N=0
      LAST=.FALSE.
50     N=N+1
      CALL TMOFA(A, TM, DTM, M)
      IF(M.GT.0) GO TO 400
      D=-TM/DTM
C      IS TOLERANCE MET

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      IF(N.EQ.40.OR.A-AC.LE.TOLA.OR.A1-A.LE.TOLA.OR.
C     DABS(D).LT.TOLA)LAST=.TRUE.
      IF(D) 110,100,120
100  CV=A
      DCV=0.00
      GO TO 320
C     REPLACE UPPER BOUND BY A
110  A1=A
      GO TO 200
C     REPLACE LOWER BOUND BY A
120  A0=A
200  A2=A+D
      IF(LAST) GO TO 300
      IF(A2.GT.A0.AND.A2.LT.A1) GO TO 250
      A=.500*(A0+A1)
      GO TO 50
250  A=A2
      GO TO 50
300  IF(A2.LE.A0.OR.A2.GE.A1) GO TO 350
      CALL TMOFA(A2,TM,DTM,M)
      IF(M.GT.0) GO TO 400
      D=-TM/DTM
      CV=A2
310  DCV=D
320  MM=M
      RETURN
350  CV=A
      GO TO 310
400  CV=0.00
      DCV=0.00
      GO TO 320
      END
      SUBROUTINE TMOFA(ALFA,TM,DTM,ND)
      INTEGER K,KK,KT,L,MF,MO,M1,M2S,ND,TYPE
      DOUBLE PRECISION A,AA,ALFA,B,DG,DTM,DTYPE,
C     F,FL,G,H(200),HP,Q,QINV,
C     Q1,Q2,T,TM,TOL,TT,V
      COMMON G(200,2),DG(200,2),AA,A(3),B(3),DTYPE,QINV,
C     Q1,Q2,T,TT,K,L,KK,KT
      COMMON/MF1/Q,TOL,TYPE,M1,MO,M2S,MF
      EQUIVALENCE (H(1),G(1,1)),(Q1,HP),(Q2,F)
      DATA FL/1.00+30/
C     STATEMENT FUNCTION
      V(K)=(AA-DBLE(FLDAT(K))*2)/Q
      ND=0
      KT=0
      AA=ALFA
      DTYPE=TYPE
      QINV=1.000/Q
      DO 10 L=1,2
        DO 5 K=1,200

```

```

      G(K,L)=0.00
      DG(K,L)=0.00
5      CONTINUE
10     CONTINUE
      IF(MOD(TYPE,2)) 20,30,20
20     M0=3
      GO TO 40
30     M0=TYPE+2
40     K=.500+DSQRT(DMAX1(3.000*Q+AA,C.00))
      M2S=MIN0(2*K+M0+4,398+MOD(M0,2))
C      EVALUATION OF THE TAIL OF A CONTINUED FRACTION
      A(1)=1.000
      A(2)=V(M2S+2)
      B(1)=V(M2S)
      B(2)=A(2)*B(1)-1.000
      Q1=A(2)/B(2)
      DO 50 K=1,200
      MF=M2S+2+2*K
      T=V(MF)
      A(3)=T*A(2)-A(1)
      B(3)=T*B(2)-B(1)
      Q2=A(3)/B(3)
      IF(DABS(Q1-Q2).LT.TOL) GO TO 70
      Q1=Q2
      A(1)=A(2)
      A(2)=A(3)
      B(1)=B(2)
      B(2)=B(3)
50     CONTINUE
      KT=1
70     T=1.000/T
      TT=-T*T*QINV
      L=MF-M2S
      DO 80 K=2,L,2
      T=1.000/(V(MF-K)-T)
      TT=T*T*(TT-QINV)
80     CONTINUE
      KK=M2S/2+1
      IF(KT.EQ.1) Q2=T
      G(KK,2)=.500*(Q2+T)
      DG(KK,2)=TT
C      STAGE 1
      G(2,1)=1.000
      DO 140 K=M0,M2S,2
      KK=K/2+1
      IF(K.LT.5) IF(K-3) 100,110,120
      G(KK,1)=V(K-2)-1.000/G(KK-1,1)
      DG(KK,1)=QINV+DG(KK-1,1)/G(KK-1,1)*2
      GO TO 130
100    G(2,1)=V(0)
      DG(2,1)=QINV

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      GO TO 130
110  G(2,1)=V(1)+DTYPE-2.000
      DG(2,1)=QINV
      GO TO 130
120  G(3,1)=V(2)+(DTYPE-2.00)/G(2,1)
      DG(3,1)=QINV+(2.00-DTYPE)*DG(2,1)/G(2,1)**2
      IF(TYPE.EQ.2) G(2,1)=0.00
130  IF(DABS(G(KK,1)).LT.1.00) GO TO 200
140  CONTINUE
C    BACKTRACK
      TM=G(KK,2)-G(KK,1)
      DTM=DG(KK,2)-DG(KK,1)
      M1=M2S
      KT=M2S-M0
      DO 180 L=2,KT,2
      K=M2S-L
      KK=K/2+1
      G(KK,2)=1.00/(V(K)-G(KK+1,2))
      DG(KK,2)=-G(KK,2)**2*(QINV-DG(KK+1,2))
      IF(K-2) 150,150,160
150  G(2,2)=2.000*G(2,2)
      DG(2,2)=2.00*DG(2,2)
160  T=G(KK,2)-G(KK,1)
      IF(DABS(T)-DABS(TM)) 170,180,180
170  TM=T
      DTM=DG(KK,2)-DG(KK,1)
      M1=K
180  CONTINUE
      GO TO 320
C    STAGE 2
200  M1=K
      K=M2S
      KK=K/2+1
210  IF(K.EQ.M1) IF(K-2) 300,300,310
      K=K-2
      KK=KK-1
      T=V(K)-G(KK+1,2)
      IF(DABS(T)-1.00) 250,220,220
220  G(KK,2)=1.000/T
      DG(KK,2)=(DG(KK+1,2)-QINV)/T**2
      GO TO 210
C    STAGE 3
250  IF(K.EQ.M1) IF(T) 220,290,220
      HP=DG(KK+1,2)-QINV
260  G(KK,2)=FL
      H(KK)=T
      K=K-2
      KK=KK-1
      F=V(K)*T-1.00
      IF(K.EQ.M1) IF(F) 280,290,280
      IF(DABS(F)-DABS(T)) 270,280,280

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270  HP=HP/T**2-QINV
      T=F/T
      GO TO 260
280  G(KK,2)=T/F
      DG(KK,2)=(HP-QINV*T*T)/F**2
      GO TO 210
290  ND=1
      GO TO 320
C    CHAINING M EQUALS 2
300  G(2,2)=2.00*G(2,2)
      DG(2,2)=2.00*DG(2,2)
310  TM=G(KK,2)-G(KK,1)
      DTM=DG(KK,2)-DG(KK,1)
320  RETURN
      END
      SUBROUTINE EXPAND(QD,FNC,R,CV,NORM,CD,N)
C    PURPOSE:  TO GET EXPANDING COEFFICIENTS FROM ROUTINE
C              COEF.  TERMINATE THE TERMS FOR REQD. ACCUTACY
C              AND DO THE NORMALIZATION
      DIMENSION CD(25)
      DOUBLE PRECISION A,CV,QD,Q,TOL,T,AB,ERR,DABS,
C              CD,SUM,T1,DSQRT,SUM1
      INTEGER R,FNC,TYPE,CASE,NORM
      COMMON DUM1(1600),A,T,DUM2(6),AB(200)
      COMMON /MF1/Q,TOL,TYPE,M1,MD,M25,MF
101  FORMAT('0','THE # OF EXPANDING COEFFICIENT REQD. IS MORE
C    THAN 25',5X,'WARNING')
102  FORMAT('0','ERROR IN SUBPROGRAM','TMOFA VIA COEF. VERIFY
C    ARGUMENTS  NO OUTPUT')
      TOL=1.0D-13
C    TO TEST THE LAST VALUE OF ARRAY CD ERR IS USED
      ERR=1.0D-20
      Q=QD
      TYPE=2*MOD(FNC,2)+MOD(R,2)
C    FOR NEGATIVE Q AND ODD ORDERS, EXP. COEFFS. FOR EVEN AND
C    ODD FUNCTIONS ARE INTERCHANGED.
      IF(Q.LT.0.CD0.AND.MOD(R,2).EQ.1)TYPE=2*MOD((FNC-1),2)+MOD(R,2)
      M=0
      A=CV
      Q=DABS(Q)
      CALL COEF(M)
      IF(M.EQ.0) GO TO 5
      WRITE(6,102)
      RETURN
5    TYPE=2*MOD(FNC,2)+MOD(R,2)
      CASE=TYPE+1
C    THE COEFFICIENTS PASSED THRU COMMON ARRAY AB IN DOUBLE
C    PRECISION IS GIVEN TO AN ARRAY CD OF LENGTH 25 FOR
C    FURTHER PROCESSION
      DO 10 I=1,25
      CD(I)=AB(I)

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      IF(DABS(CD(I)).LT.1.D-30) CD(I)=0.00
10  CONTINUE
      IF(CD(25).GT.ERR) WRITE(6,101)
      N1=R/2+1
      DO 20 I=N1,25
      IF(CD(I).EQ.0.00) GO TO 25
      N=I
20  CONTINUE
C   NORMALISING THE CODES. PRESENTLY IN NEUTRAL NORM
25  SUM=0.000
      IF(NORM.EQ.1) GO TO 140
C   GETTING STRATTON NORMALISATION FACTOR
      IF(QD.LT.0.00) GO TO 91
      GO TO (40,40,60,80),CASE
40  DO 50 J=1,N
      SUM=SUM+CD(J)
50  CONTINUE
      GO TO 100
60  DO 70 J=1,N
      SUM=SUM+CD(J)*DBLE(FLOAT(2*(J-1)))
70  CONTINUE
      GO TO 100
80  DO 90 J=1,N
      SUM=SUM+CD(J)*DBLE(FLOAT(2*J-1))
90  CONTINUE
      GO TO 100
C   GOT NEGATIVE Q STRATTON NORMALISATION FACTOR IS DIFFERENT
91  T1=-1.000
      IF(MOD(R/2,2).EQ.1) T1=-T1
      GO TO(92,92,94,96),CASE
92  DO 93 J=1,N
      T1=-T1
      SUM=SUM+T1*CD(J)
93  CONTINUE
      GO TO 100
94  DO 95 J=1,N
      T1=-T1
      SUM=SUM+CD(J)*T1*DBLE(FLOAT(2*(J-1)))
95  CONTINUE
      GO TO 100
96  DO 97 J=1,N
      T1=-T1
      SUM=SUM+CD(J)*T1*DBLE(FLOAT(2*J-1))
97  CONTINUE
100 IF(NORM.EQ.2) GO TO 120
C   GETTING INCE'S NORMALISATION FACTOR
      SUM1=0.00
      DO 110 J=1,N
      SUM1=SUM1+CD(J)*CD(J)
110 CONTINUE
      IF(FNC.EQ.2.AND.MOD(R,2).EQ.0) SUM1=SUM1+CD(1)*CD(1)

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SUM1=DSQRT(SUM1)
IF(NORM.EQ.3) SUM=DSIGN(SUM1,SUM)
C   DIVIDE ALL COEFS. BY NORMALISING FACTOR FOR 2 & 3 ONLY
120  DO 130 I=1,N
    CD(I)=CD(I)/SUM
130  CONTINUE
C   FOR MATHIEU FUNCTIONS OF SE2N+2 TYPE(CASE=3) COEFS. SHOULD
C   BE B2, B4 ETC. BUT THE ROUTINE COEF RETURNS A B0=0 ALSO.
C   THIS IS TO BE DROPPED.
140  IF(CASE.NE.3) RETURN
    DO 150 I=2,N
    CD(I-1)=CD(I)
150  CONTINUE
    CD(N)=0.DO
    RETURN
    END
    SUBROUTINE COEF(M)
    INTEGER K,KA,KB,KK,M,MF,ML,MM,MO,M1,M2S,TYPE
    DOUBLE PRECISION A,AB,FL,G,H(200),Q,T,TOL,V,V2
    COMMON G(200,2),DUM1(800),A,T,K,KA,KB,KK,MM,ML,AB(200)
    COMMON /MF1/Q,TOL,TYPE,M1,MO,M2S,MF
    EQUIVALENCE (H(1),G(1,1))
    DATA FL,V2/1.0+30,1.0-15/
C   STATEMENT FUNCTION
    V(K)=(A-DBLE(FLOAT(K))*2)/Q
    CALL TMOFA(A,T,T,M)
    IF(M.NE.0) GO TO 300
    DO 60 K=1,200
    AB(K)=0.DO
60   CONTINUE
    KA=M1-MO+2
    DO 90 K=2,KA,2
    KK=(M1-K)/2+1
    IF(K-2) 70,70,90
70   AB(KK)=1.DO
    GO TO 90
80   AB(KK)=AB(KK+1)/G(KK+1,1)
90   CONTINUE
    KA=0
    DO 130 K=M1,M2S,2
    KK=K/2+1
    ML=K
    IF(G(KK,2).EQ.FL) GO TO 100
    AB(KK)=AB(KK-1)*G(KK,2)
    GO TO 110
100  T=AB(KK-2)
    IF(K.EQ.4.AND.M1.EQ.2) T=T+T
    AB(KK)=T/(V(K-2)*H(KK)-1.DO)
110  IF(DABS(AB(KK)).GE.1.0-17) KA=C
    IF(KA.EQ.5) GO TO 260
    KA=KA+1

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130  CONTINUE
      T=DLOG(DABS(AB(KK))/V2)/DLOG(1.D0/DABS(G(KK,2)))
      KA=2*IDINT(T)
      ML=KA+2+M2S
      IF(ML.GT.399) GO TO 400
      KB=KA+2+MF
      T=1.D0/V(KB)
      KK=MF-M2S
      DO 150 K=2, KK, 2
          T=1.D0/(V(KB-K)-T)
150  CONTINUE
      KK=ML/2+1
      G(KK,2)=T
      DO 200 K=2, KA, 2
          KK=(ML-K)/2+1
          G(KK,2)=1.D0/(V(ML-K)-G(KK+1,2))
200  CONTINUE
      KA=M2S+2
      DO 250 K=KA, ML, 2
          KK=K/2+1
          AB(KK)=AB(KK-1)*G(KK,2)
250  CONTINUE
C      NEUTRAL NORMALIZATION
260  T=AB(1)
      MM=MOD(TYPE,2)
      KA=MM+2
      DO 280 K=KA, ML, 2
          KK=K/2+1
          IF(DABS(T)-DABS(AB(KK))) 270,280,280
270  T=AB(KK)
      MM=K
280  CONTINUE
      DO 290 K=1, KK
          AB(K)=AB(K)/T
290  CONTINUE
300  RETURN
400  M=-1
      GO TO 300
      END
      DOUBLE PRECISION FUNCTION ANGMFC(QD,FNC,R,XD,DERIV,AR,N)
C      PURPOSE:  TO COMPUTE A PERIODIC MATHIEU FUNCTION,
C                ODD OR EVEN TYPE OR ITS DERIVATIVE
      DOUBLE PRECISION PC,PS,DPC,DPS
      EXTERNAL PC,PS,DPC,DPS
      DIMENSION AR(25),AB(25)
      INTEGER FNC,R,DERIV,TYPE,CASE,P
      DOUBLE PRECISION AR,AB,XD,X,T1,SUM,DCOS,DSIN,QD
      COMMON/NTERM/NLIMIT
      COMMON/ANG/AB,X,P
      NLIMIT=N
      QS=QD

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X=XD
TYPE=2*MOD(FNC,2)+MOD(R,2)
CASE=TYPE+1
DO 1 I=1,N
  AB(I)=AR(I)
1  CONTINUE
C  FOR NEGATIVE Q IN ALL SUMMATIONS ALTERNATE TERMS HAVE
C  A MINUS SIGN
  IF(QS)20,90,35
20  T1=-1.000
  IF(CASE.EQ.3)T1=1.000
  IF(MOD(R/2,2).EQ.1)T1=-T1
  DO 30 I=1,N
    T1=-T1
    AB(I)=T1*AB(I)
30  CONTINUE
35  P=-1
  IF(CASE.EQ.1) P=-2
  IF(CASE.EQ.3) P=0
  IF(DERIV.EQ.1) GO TO 60
  GO TO(40,40,50,50),CASE
40  CALL SIGMA(PC,SUM)
  ANGMFC=SUM
  RETURN
50  CALL SIGMA(PS,SUM)
  ANGMFC=SUM
  RETURN
60  GO TO(70,70,80,80),CASE
70  CALL SIGMA(DPC,SUM)
  ANGMFC=SUM
  RETURN
80  CALL SIGMA(DPS,SUM)
  ANGMFC=SUM
  RETURN
90  IF(DERIV.EQ.1)GO TO 120
  GO TO(100,100,110,110),CASE
100  ANGMFC=DCOS(DBLE(FLOAT(R))*X)
  RETURN
110  ANGMFC=DSIN(DBLE(FLOAT(R))*X)
  RETURN
120  GO TO(130,130,140,140),CASE
130  ANGMFC=-R*DSIN(DBLE(FLOAT(R))*X)
  RETURN
140  ANGMFC=R*DCOS(DBLE(FLOAT(R))*X)
  RETURN
  END
  DOUBLE PRECISION FUNCTION PC(K)
  INTEGER P,K
  DOUBLE PRECISION AB(25),X,DCOS
  COMMON/ANG/AB,X,P
C  EVALUATES ONE TERM OF THE EVEN PERIODIC SOLUTION

```

```

PC=AB(K)*DCOS(DBLE(FLOAT(2*K+P))*X)
RETURN
END
DOUBLE PRECISION FUNCTION PS(K)
INTEGER P,K
DOUBLE PRECISION AB(25),X,DSIN
COMMON/ANG/AB,X,P
C EVALUATES ONE TERM OF THE ODD PERIODIC SOLUTION
PS=AB(K)*DSIN(DBLE(FLOAT(2*K+P))*X)
RETURN
END
DOUBLE PRECISION FUNCTION DPC(K)
INTEGER P,K
DOUBLE PRECISION AB(25),X,T,DSIN
COMMON/ANG/AB,X,P
C EVALUATES ONE TERM OF THE DERIVATIVE OF THE EVEN PERIODIC
C MATHIEU FUNCTION.
T=2*K+P
DPC=-AB(K)*T*DSIN(T*X)
RETURN
END
DOUBLE PRECISION FUNCTION DPS(K)
INTEGER P,K
DOUBLE PRECISION AB(25),X,T,DCOS
COMMON/ANG/AB,X,P
C EVALUATES ONE TERM OF THE DERIVATIVE OF THE ODD PERIODIC
C MATHIEU FUNCTION.
T=2*K+P
DPS=AB(K)*T*DCOS(T*X)
RETURN
END
SUBROUTINE SIGMA(DUM,SUM)
C PURPOSE: TO SUM N TERMS (SPECIFIED BY THE COMMON
C COMMON BLOCK N TERM) OF A FUNCTION.
DOUBLE PRECISION DUM,SUM,ERR,T1,TERM,DABS
COMMON/NTERM/NLIMIT
101 FORMAT('C','CONVERGENCE NOT TO SATISFACTION',10X,'WARNING')
ERR=1.0D-13
T1=DUM(1)
SUM=T1
N=NLIMIT
IF(NLIMIT.GE.22)N=22
IMIN=5
DO 10 I=2,N
TERM=DUM(I)
SUM=SUM+TERM
TERM=DABS(TERM)
IF(I.LT.IMIN) GO TO 10
IF(DABS(SUM).GE.ERR*TERM) RETURN
10 CONTINUE
20 IF(I.EQ.22)WRITE(6,101)

```

```

RETURN
END
SUBROUTINE FACTOR(FNC,R,QD,AB,N,PS)
C   PURPOSE:  TO COMPUTE THE NORMALIZATION FACTOR,P2N,P2N+1,
C             S2N+1,OR S2N+2 FOR MODIFIED MATHIEU FUNCTIONS
C             OF POSITIVE 'Q' AND P2N',P2N+1',S2N+1',
C             OR S2N+2' FOR THOSE OF NEGATIVE 'Q'
C
DIMENSION A9(25)
INTEGER FNC,R,CASE
DOUBLE PRECISION AB,PS,PSP,DABS,DSQRT,RQ,ODEV,
C             SUM1,SUM2,T1,T2,QD
RQ=DSQRT(DABS(QD))
CASE=2*MOD(FNC,2)+MOD(R,2)+1
IF(QD.LT.0.0D0)CASE=CASE+4.
ODEV=1.0D0
IF(MOD(R/2,2).EQ.1)ODEV=-1.0D0
SUM1=0.0D0
SUM2=0.0D0
T1=-1.0D0
GO TO (10,30,50,70,10,70,50,30),CASE
C   FOR ALL Q AND EVEN ORDER ---- P2N AND P2N' ---- FNC=2
10  DO 20 I=1,N
    SUM1=SUM1+AB(I)
    T1=-T1
    SUM2=SUM2+T1*AB(I)
20  CONTINUE
    PS=SUM1*SUM2/A9(1)
    PSP=PS*ODEV
    IF(QD.LT.0.0D0)PS=PSP
    RETURN
C   FOR POSITIVE Q AND ODD ORDERS P2N+1, P2N+1' IF FNC=2
C   NEG Q AND ODD ORDER IF FNC=1
30  DO 40 I=1,N
    T1=-T1
    SUM1=SUM1+AB(I)
    SUM2=SUM2+T1*AB(I)*DBLE(FLOAT(2*I-1))
40  CONTINUE
    PS=SUM1*SUM2/(RQ*AB(1))
    PSP=PS*ODEV
    IF(QD.LT.0.0D0) PS=PSP
    RETURN
C   FOR ALL Q AND EVEN ORDER IF FNC=1 ---- S2N+2, S2N+2'
50  T1=1.0D0
    DO 60 I=1,N
    T2=AB(I)*DBLE(FLOAT(2*I))
    SUM1=SUM1+T2
    T1=-T1
    SUM2=SUM2+T1*T2
60  CONTINUE
    PS=SUM1*SUM2/(RQ*RQ*AB(1))
    PSP=PS*ODEV

```

```

      IF(QD.LT.0.0D0)PS=PSP
      RETURN
C     FOR POSITIVE Q AND ODD ORDER ---- IF FNC=1
C     FOR NEGATIVE Q AND ODD ORDERS S2N+1, S2N+1°   IF FNC=2
70    DO 80 I=1,N
      SUM1=SUM1+AB(I)*DBLE(FLOAT(2*I-1))
      T1=-T1
      SUM2=SUM2+T1*AB(I)
80    CONTINUE
      PS=SUM1*SUM2/(RQ*AB(1))
      PSP=PS*ODEV
      IF(QD.LT.0.0D0)PS=PSP
      RETURN
      END
      SUBROUTINE STORE(QD,XI,NMAX)
C     PURPOSE:  TO COMPUTE AND STORE THE VALUES OF BESSEL
C               FUNCTIONS AND DERIVATIVES REQUIRED IN MATHIEU
C               FUNCTION CALCULATION.
      DOUBLE PRECISION QD,QABS,RQ,DABS,DSQRT,DEXP,BS1V1,BS1V2,
*          BS2V2,DBS1V1,DBS1V2,DBS2V2,XI,V1,V2
      COMMON/LOCAL/DUMMY1(4),V1,V2,DUMMY2(50)
      COMMON/RADIAL/BS1V1(25),BS1V2(25),BS2V2(25),DBS1V1(25),
*          DBS1V2(25),DBS2V2(25)
      N=NMAX+3
      IF(N.GE.25)N=25
      N1=N-1
      QABS=DABS(QD)
      RQ=DSQRT(QABS)
      V1=RQ*DEXP(-XI)
      V2=QABS/V1
      IF(QD.LT.0.0D0) GO TO 20
      CALL BESSEL(1,V1,BS1V1,N)
      CALL BESSEL(1,V2,BS1V2,N)
      CALL BESSEL(2,V2,BS2V2,N)
      DBS1V1(1)=-BS1V1(2)
      DBS1V2(1)=-BS1V2(2)
      DBS2V2(1)=-BS2V2(2)
      DO 10 I=2,N1
      DBS1V1(I)=(BS1V1(I-1)-BS1V1(I+1))*0.5D0
      DBS1V2(I)=(BS1V2(I-1)-BS1V2(I+1))*0.5D0
      DBS2V2(I)=(BS2V2(I-1)-BS2V2(I+1))*0.5D0
10    CONTINUE
      RETURN
20    CALL BSL2(1,V1,BS1V1,N)
      CALL BSL2(1,V2,BS1V2,N)
      CALL BSL2(2,V2,BS2V2,N)
      DBS1V1(1)=BS1V1(2)
      DBS1V2(1)=BS1V2(2)
      DBS2V2(1)=-BS2V2(2)
      DO 30 I=2,N1
      DBS1V1(I)=(BS1V1(I-1)+BS1V1(I+1))*0.5D0

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DBS1V2(I)=(BS1V2(I-1)+BS1V2(I+1))*0.5DC
DBS2V2(I)=-(BS2V2(I-1)+BS2V2(I+1))*0.5DC
30 CONTINUE
RETURN
END
SUBROUTINE BESSEL(SOL,U,BSJY,N)
INTEGER N,NN,SOL
DOUBLE PRECISION BSJY(N),U
NN=N-1
IF(U.EQ.0.DC.AND.SOL.EQ.2) GO TO 80
IF(U.GE.8.D0) GO TO 30
GO TO(10,20),SOL
10 CALL JOJ1(U,BSJY)
GO TO 40
20 CALL YOY1(U,BSJY)
GO TO 40
30 CALL LUKE(U,SOL,BSJY)
40 IF(N.LT.2) GO TO 100
GO TO(50,60),SOL
50 CALL JNS(BSJY,U,N)
GO TO 100
C RECURRENCE FORMULAR
60 DO 70 K=2,NN
BSJY(K+1)=2.DC*DBLE(FLOAT(K-1))*BSJY(K)/U-BSJY(K-1)
70 CONTINUE
GO TO 100
80 NN=NN+1
DO 90 K=1,NN
BSJY(K)=-1.D+37
90 CONTINUE
100 RETURN
END
SUBROUTINE JOJ1(X,BJ)
DOUBLE PRECISION BJ(2),T(5),X
T(1)=X/2.D0
BJ(1)=1.D0
BJ(2)=T(1)
T(2)=-T(1)**2
T(3)=1.D0
T(4)=1.D0
10 T(4)=T(4)*T(2)/T(3)**2
BJ(1)=BJ(1)+T(4)
T(5)=T(4)*T(1)/(T(3)+1.D0)
BJ(2)=BJ(2)+T(5)
IF(DMAX1(DABS(T(4)),DABS(T(5))).LT.1.D-15) RETURN
T(3)=T(3)+1.D0
GO TO 10
END
SUBROUTINE YCY1(X,BY)
DOUBLE PRECISION T(10),X,BY(2)
T(1)=X/2.D0

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      T(2)=-T(1)**2
      BY(1)=1.00
      BY(2)=T(1)
      T(7)=0.00
      T(10)=-T(1)
      T(3)=0.00
      T(4)=0.00
      T(5)=1.00
1C    T(3)=T(3)+1.00
      T(4)=T(4)+1.00/T(3)
      T(5)=T(5)*T(2)/T(3)**2
      BY(1)=BY(1)+T(5)
      T(6)=-T(5)*T(4)
      T(7)=T(7)+T(6)
      T(8)=T(5)*T(1)/(T(3)+1.00)
      BY(2)=BY(2)+T(8)
      T(9)=-T(8)*(2.00*T(4)+1.00/(T(3)+1.00))
      T(10)=T(10)+T(9)
      IF(DMAX1(DABS(T(6)),DABS(T(9))),GE.1.0-15) GO TO 10
      T(2)=.5772156649015328600+DLOG(T(1))
      BY(1)=.6366197723675813400*(BY(1)*T(2)+T(7))
      BY(2)=.6366197723675813400*(BY(2)*T(2)-1.00/X)+T(10)/
C      3.141592653589793200
      RETURN
      END
      SUBROUTINE LUKE(U,KIND,BSJY)
      INTEGER K,KIND
      DOUBLE PRECISION A(19),B(19),CS,C(19),D(19),G(3),BSJY(2),
C      R(2),S(2),SN,T,U,X
C    WARNING - THE FOLLOWING DATA STATEMENTS ARE NOT IN ASA
C    STANDARD FORTRAN
      DATA A/.9995950647696728741600,
*      -.538079561396069130-3,
*      -.131796771233615700-3,
*      .1514224970486440-5,
*      .158468617920630-6,
*      -.8560695539460-8,
*      -.295723433550-9,
*      .65735562540-10,
*      -.2237497030-11,
*      -.448211400-12,
*      .69548270-13,
*      -.1513400-14,
*      -.924220-15,
*      .155580-15,
*      -.4760-17,
*      -.2740-17,
*      .610-18,
*      -.40-19,
*      -.10-19/
      DATA B/-.7769355694205321360-2,

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*      -.7748032309654476700-2,
*      .25365411654307960-4,
*      .3942735983997110-5,
*      -.107234982991290-6,
*      -.7213897993280-8,
*      .737646028930-9,
*      .1506878110-11,
*      -.5745895370-11,
*      .459965740-12,
*      .22703230-13,
*      -.8878900-14,
*      .744970-15,
*      .58470-16,
*      -.24100-16,
*      .2650-17,
*      .130-18,
*      -.100-18,
*      .20-19/
DATA C/.10006775358659134623400,
*      .901007251959081830-3,
*      .221724349185994540-3,
*      -.1965759463191040-5,
*      -.20889531143270-6,
*      .10281443508940-7,
*      .375970547890-9,
*      -.76388913580-10,
*      .2387346700-11,
*      .518254890-12,
*      -.76939690-13,
*      .1440080-14,
*      .1032940-14,
*      -.168210-15,
*      .4590-17,
*      .3020-17,
*      -.650-18,
*      .40-19,
*      .10-19/
DATA D/.23376829936285803280-1,
*      .23346801223545575330-1,
*      -.35760105909013820-4,
*      -.5608631494926270-5,
*      .132738940843400-6,
*      .9169758450660-8,
*      -.868388803710-9,
*      -.3780730050-11,
*      .6631455860-11,
*      -.505843900-12,
*      -.27207820-13,
*      .9853810-14,
*      -.793980-15,
*      -.67570-16,

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*      .2625D-16,
*      -.280D-17,
*      -.15D-18,
*      .10D-18,
*      -.2D-19/
X=8.D0/U
G(1)=1.D0
G(2)=2.D0*X-1.D0
R(1)=A(1)+A(2)*G(2)
S(1)=B(1)+B(2)*G(2)
R(2)=C(1)+C(2)*G(2)
S(2)=D(1)+D(2)*G(2)
DO 10 K=3,19
G(3)=(4.D0*X-2.D0)*G(2)-G(1)
R(1)=R(1)+A(K)*G(3)
S(1)=S(1)+B(K)*G(3)
R(2)=R(2)+C(K)*G(3)
S(2)=S(2)+D(K)*G(3)
G(1)=G(2)
G(2)=G(3)
10 CONTINUE
T=.79788456C8028654D0/DSQRT(U)
SN=DSIN(U-.7853981633974483DC)
CS=DCOS(U-.7853981633974483DC)
GO TO(20,30),KIND
20 BSJY(1)=T*(R(1)*CS-S(1)*SN)
BSJY(2)=T*(R(2)*SN+S(2)*CS)
GO TO 40
30 BSJY(1)=T*(S(1)*CS+R(1)*SN)
BSJY(2)=T*(S(2)*SN-R(2)*CS)
40 RETURN
END
SUBROUTINE JNS(JJ,U,M)
INTEGER K,KA,KK,M
DOUBLE PRECISION A,B,D(2),DM,G(25),JJ(M),
C      P(3),Q(3),U
DM=2*M
P(1)=0.D0
Q(1)=1.D0
P(2)=1.D0
Q(2)=DM/U
D(1)=P(2)/Q(2)
A=2.D0
10 B=(DM+A)/U
P(3)=B*P(2)-P(1)
Q(3)=B*Q(2)-Q(1)
D(2)=P(3)/Q(3)
IF(DABS(D(1)-D(2)).LT.1.D-15) GO TO 20
P(1)=P(2)
P(2)=P(3)
Q(1)=Q(2)

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      Q(2)=Q(3)
      D(1)=D(2)
      A=A+2.00
      GO TO 10
20    G(M)=D(2)
      KA=M-2
      DO 30 K=1,KA
      KK=M-K
      A=2*KK
      G(KK)=U/(A-U*G(KK+1))
      IF(G(KK).EQ.0.00) G(KK)=1.0-35
30    CONTINUE
      DO 40 K=2,M
      JJ(K+1)=G(K)*JJ(K)
40    CONTINUE
      RETURN
      END
      SUBROUTINE BSL2(SOL,U,BSIK,N)
      PURPOSE:  TO COMPUTE MODIFIED BESSEL FUNCTIONS, 'I' OR 'K'
      C          TYPE FOR ORDERS 0 THRU N-1 IN DOUBLE PRECISION.
      C
      INTEGER N,SOL
      DOUBLE PRECISION BSIK(N),U
      IF(SOL.EQ.2)GO TO 30
      IF(U.GE.8.000) GO TO 10
      CALL IOI1(U,BSIK)
      GO TO 20
10    CALL LUKE2(U,SOL,BSIK)
20    IF(N.LT.2)RETURN
      CALL INS(BSIK,U,N)
      RETURN
30    IF(U.EQ.0.000) GO TO 70
      IF(U.GE.5.000) GO TO 40
      CALL KOK1(U,BSIK)
      GO TO 50
40    CALL LUKE2(U,SOL,BSIK)
50    IF(N.LT.2)RETURN
      C      RECURRENCE FORMULA
      NN=N-1
      DO 60 K=2,NN
      BSIK(K+1)=2.000*DBLE(FLOAT(K-1))*BSIK(K)/U+BSIK(K-1)
60    CONTINUE
      RETURN
70    DO 80 K=1,N
      BSIK(K)=1.00+75
80    CONTINUE
      RETURN
      END
      SUBROUTINE IOI1(X,BI)
      C      PURPOSE:  TO EVALUATE 'I0' AND 'I1' BESSEL FUNCTIONS
      C      BY SUMMING THE SERIES.
      DOUBLE PRECISION BI(2),T(5),X,DMAX1,DABS

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      T(1)=X/2.000
      BI(1)=1.000
      BI(2)=T(1)
      T(2)=T(1)**2
      T(3)=1.000
      T(4)=1.000
10    T(4)=T(4)*T(2)/T(3)**2
      BI(1)=BI(1)+T(4)
      T(5)=T(4)*T(1)/(T(3)+1.000)
      BI(2)=BI(2)+T(5)
      IF(DMAX1(DABS(T(4)),DABS(T(5))),LT.1.00-15)RETURN
      T(3)=T(3)+1.000
      GO TO 10
      END
      SUBROUTINE KOK1(X,BK)
C      PURPOSE:  TO EVALUATE 'K0' AND 'K1' BESSEL FUNCTIONS
C                BY SUMMING THE SERIES.
      DOUBLE PRECISION T(10),X,BK(2),DMAX1,DABS,DLOG
      T(1)=X/2.000
      T(2)=T(1)**2
      BK(1)=1.000
      BK(2)=T(1)
      T(7)=0.000
      T(10)=-T(1)
      T(3)=0.000
      T(4)=0.000
      T(5)=1.000
10    T(3)=T(3)+1.000
      T(4)=T(4)+1.000/T(3)
      T(5)=T(5)*T(2)/T(3)**2
      BK(1)=BK(1)+T(5)
      T(6)=T(5)*T(4)
      T(7)=T(7)+T(6)
      T(8)=T(5)*T(1)/(T(3)+1.000)
      BK(2)=BK(2)+T(8)
      T(9)=-T(8)*(2.000*T(4)+1.000/(T(3)+1.000))
      T(10)=T(10)+T(9)
      IF(DMAX1(DABS(T(6)),DABS(T(9))),GE.1.00-15) GO TO 10
      T(2)=.577215664901532900+DLOG(T(1))
      BK(1)=-BK(1)*T(2)+T(7)
      BK(2)=BK(2)*T(2)+1.000/X+T(10)/2.000
      RETURN
      END
      SUBROUTINE LUKE2(U,KIND,BSIK)
C      PURPOSE:  TO EVALUATE MODIFIED BESSEL FUNCTIONS, I0 AND I1
C                OR K0 AND K1 FROM SHIFTED CHEBYSHEV SERIES.
      DOUBLE PRECISION A(34),B(21),C(34),D(21),U,BSIK(2),X,G(34),
C                R(2),S(2),DEXP,DSQRT
      DATA A/1.0082792054587400,.8445122624920943D-2,
      *.1727006307775665D-3,.724759109995896D-5,.51358772687802D-6,
      *.5681696580912D-7,.851309122285D-8,.1238425364D-8,

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*.293016723D-10,-.7895669832D-10,-.3312712763D-10,
*-.449733864D-11,.17997903D-11,.96574832D-12,.3860424D-13,
*-.10403934D-12,-.2395045D-13,.955447D-14,.444315D-14,
*-.85864D-15,-.70878D-15,.8676D-16,.11194D-15,-.1211D-16,
*-.1813D-16,.249D-17,.299D-17,-.62D-18,-.49D-18,.16D-18,
*.7D-19,-.4D-19,-.1D-19,.1D-19/
  DATA B/.98840917423C8258D0,-.1131050446469282D-1,
*.2695326127627237D-3,-.1110668519666535D-4,
*.63257510850049D-6,-.450473376411D-7,.379299645568D-8,
*-.36454717921D-9,.3904375576D-10,-.457993622D-11,
*.59081063D-12,-.7883236D-13,.1136042D-13,
*-.172697D-14,.27545D-15,-.4589D-16,.796D-17,-.143D-17,
*.27D-18,-.5D-19,.1D-19/
  DATA C/.9758006023262859D0,-.2446744296327638D-1,
*-.2772053607638289D-3,-.973214672802013D-5,
*-.62972423863981D-6,-.6596114215424D-7,-.96138729194D-8,
*-.140114090103D-8,-.4756316654D-10,.8153068107D-10,
*.3540814832D-10,.510256407D-11,-.180440934D-11,
*-.102359447D-11,-.5267784D-13,.107C9419D-12,.2611976D-13,
*-.956129D-14,-.471335D-14,.82924D-15,.74262D-15,-.8045D-16,
*-.11657D-15,.1107D-16,.1884D-16,-.233D-17,-.311D-17,
*.61D-18,.51D-18,-.16D-18,-.8D-19,.4D-19,.1D-19,-.1D-19/
  DATA D/1.03595C858772358D0,.354652912433311D-1,
*-.4684750281668886D-3,.16185C6381005343D-4,
*-.84517204812368D-6,.5713221810284D-7,-.464555460661D-8,
*.43541733857D-9,-.4575729704D-10,.528813281D-11,
*-.66261293D-12,.8904792D-13,-.1272607D-13,.192086D-14,
*-.3045D-15,.5045D-16,-.871D-17,.156D-17,-.29D-18,
*.6D-19,-.1D-19/
  IF(KIND.EQ.2) GO TO 20
  X=8.00D/U
  G(1)=1.00D
  G(2)=2.00D*X-1.00D
  N=34
  DO 10 K=3,N
  G(K)=(4.00D*X-2.00D)*G(K-1)-G(K-2)
10  CONTINUE
  R(1)=0.00D
  R(2)=0.00D
  DO 15 K=1,N
  I=N+1-K
  R(1)=R(1)+A(I)*G(I)
  R(2)=R(2)+C(I)*G(I)
15  CONTINUE
  BS(K(1))=.39894229C4014327DC*R(1)*DEXP(U)/DSQRT(U)
  BS(K(2))=.39894229C4014327DC*R(2)*DEXP(U)/DSQRT(U)
  RETURN
20  X=5.00D/U
  G(1)=1.00D
  G(2)=2.00D*X-1.00D
  N=21

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DO 25 K=3,N
G(K)=(4.0D0*X-2.0D0)*G(K-1)-G(K-2)
25 CONTINUE
S(1)=0.0D0
S(2)=0.0D0
DO 30 K=1,N
I=N+1-K
S(1)=S(1)+B(I)*G(I)
S(2)=S(2)+D(I)*G(I)
30 CONTINUE
BSIK(1)=1.2533141373155D0*S(1)*DEXP(-U)/DSQRT(U)
BSIK(2)=1.2533141373155D0*S(2)*DEXP(-U)/DSQRT(U)
RETURN
END
SUBROUTINE INS(II,U,M)
C PURPOSE: TO EVALUATE 'I' BESSEL FUNCTIONS OF HIGHER
C ORDERS BY A CONTINUED FRACTION EXPANSION
C METHODS.
INTEGER K,KA,KK,M
DOUBLE PRECISION A,B,D(2),DM,G(25),II(M),P(3),Q(3),U
IF(U.EQ.0.0D0) GO TO 50
DM=2*M
P(1)=1.0D0
Q(1)=1.0D0
P(2)=1.0D0
Q(2)=DM/U
D(1)=P(2)/Q(2)
A=2.0D0
10 B=(DM+A)/U
P(3)=3*P(2)+P(1)
Q(3)=8*Q(2)+Q(1)
D(2)=P(3)/Q(3)
IF(DABS(D(1)-D(2)).LT.1.0D-15) GO TO 20
P(1)=P(2)
P(2)=P(3)
Q(1)=Q(2)
Q(2)=Q(3)
D(1)=D(2)
A=A+2.0D0
GO TO 10
20 G(M)=D(2)
KA=M-2
DO 30 K=1,KA
KK=M-K
A=2*KK
G(KK)=U/(A+U*G(KK+1))
IF(DABS(G(KK)).LE.1.0D-35) G(KK)=1.0D-35
30 CONTINUE
DO 40 K=2,M
IF(DABS(II(K)).LT.1.0D-35) GO TO 35
II(K+1)=G(K)*II(K)

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```

GO TO 40
35  II(K+1)=0.000
40  CONTINUE
    RETURN
50  DO 60 I=3,M
60  II(I)=0.000
    RETURN
    END
    DOUBLE PRECISION FUNCTION CERAD(QD,R,DERIV,PS,AR,N)
C    PURPOSE:  TO COMPUTE A MODIFIED MATHIEU FUNCTION
C              (OR DERIVATIVE) OF FIRST KIND CORRESPONDING
C              TO EVEN MATHIEU FUNCTION ( CE FUNCTIONS)
    EXTERNAL C2NP,C2NIP,C2NN,C2NIN,DC2NP,DC2NIP,DC2NN,DC2NIN
    DOUBLE PRECISION AB(25),QD,PS,PSP,OUTPUT,AR(25)
    INTEGER R,CASE,DERIV,FNC
    COMMON/NTerm/N1
    COMMON/LOCAL/DUMMY(8),AB
    N1=N
    PSP=PS
    DO 5 I=1,N
    AB(I)=AR(I)
5    CONTINUE
    CASE=MOD(R,2)+1
    IF(QD.LT.0.000) CASE=CASE+2
    IF(DERIV.EQ.1) GO TO 50
    GO TO(10,20,30,40),CASE
C    THE VALUE OF CE2N(Z,Q)
10   CALL SIGMA(C2NP,OUTPUT)
    CERAD=PS*OUTPUT/AB(1)
    RETURN
C    THE VALUE OF CE2N+1(Z,Q)
20   CALL SIGMA(C2NIP,OUTPUT)
    CERAD=PS*OUTPUT/AB(1)
    RETURN
C    THE VALUE OF CE2N(Z,-Q)
30   CALL SIGMA(C2NN,OUTPUT)
    CERAD=PSP*OUTPUT/AB(1)
    RETURN
C    THE VALUE OF CE2N+1(Z,-Q)
40   CALL SIGMA(C2NIN,OUTPUT)
    CERAD=PSP*OUTPUT/AB(1)
    RETURN
C    FOLLOWING ARE DERIVATIVES OF FUNCTIONS
50   GO TO(60,70,80,90),CASE
C    THE VALUE OF CE2N'(Z,Q)
60   CALL SIGMA(DC2NP,OUTPUT)
    CERAD=PS*OUTPUT/AB(1)
    RETURN
C    THE VALUE OF CE2N+1'(Z,Q)
70   CALL SIGMA(DC2NIP,OUTPUT)
    CERAD=PS*OUTPUT/AB(1)

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      RETURN
C     THE VALUE OF CE2N*(Z,-Q)
80    CALL SIGMA(DC2NN,OUTPUT)
      CERAD=PSP*OUTPUT/AB(1)
      RETURN
C     THE VALUE OF CE2N+1*(Z,-Q)
90    CALL SIGMA(DC2NIN,OUTPUT)
      CERAD=PSP*OUTPUT/AB(1)
      RETURN
      END
      DOUBLE PRECISION FUNCTION C2NP(K)
C     PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C               FOR CE2N(Z,Q).
      DOUBLE PRECISION AB(25),BSJV1,BSJV2,BSYV2,DBSJV1,DBSJV2,
*              DBSYV2
      COMMON/LOCAL/DUMMY(8),AB
      COMMON/RADIAL/BSJV1(25),BSJV2(25),BSYV2(25),DBSJV1(25),
C              DBSJV2(25),DBSYV2(25)
      C2NP=AB(K)*BSJV1(K)*BSJV2(K)
      IF(MOD(K,2).EQ.0)C2NP=-C2NP
      RETURN
      END
      DOUBLE PRECISION FUNCTION C2NIP(K)
C     PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C               FOR CE2N+1(Z,Q).
      DOUBLE PRECISION AB(25),BSJV1,BSJV2,BSYV2,DBSJV1,DBSJV2,
*              DBSYV2
      COMMON/LOCAL/DUMMY(8),AB
      COMMON/RADIAL/BSJV1(25),BSJV2(25),BSYV2(25),DBSJV1(25),
C              DBSJV2(25),DBSYV2(25)
      C2NIP=AB(K)*(BSJV1(K)*BSJV2(K+1)+BSJV1(K+1)*BSJV2(K))
      IF(MOD(K,2).EQ.0)C2NIP=-C2NIP
      RETURN
      END
      DOUBLE PRECISION FUNCTION C2NN(K)
C     PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C               FOR CE2N(Z,-Q).
      DOUBLE PRECISION AB(25),BSIV1,BSIV2,BSKV2,DBSIV1,DBSIV2,
*              DBSKV2
      COMMON/LOCAL/DUMMY(8),AB
      COMMON/RADIAL/BSIV1(25),BSIV2(25),BSKV2(25),DBSIV1(25),
*              DBSIV2(25),DBSKV2(25)
      C2NN=AB(K)*BSIV1(K)*BSIV2(K)
      IF(MOD(K,2).EQ.0)C2NN=-C2NN
      RETURN
      END
      DOUBLE PRECISION FUNCTION C2NIN(K)
C     PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES FOR CE2N+1(Z,-Q)
      DOUBLE PRECISION AB(25),BSIV1,BSIV2,BSKV2,DBSIV1,DBSIV2,DBSKV2
      COMMON/LOCAL/DUMMY(8),AB
      COMMON/RADIAL/BSIV1(25),BSIV2(25),BSKV2(25),DBSIV1(25),DBSIV2(25),

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C          DBSKV2(25)
C2NIN=AB(K)*(BSIV1(K)*BSIV2(K+1)+BSIV1(K+1)*BSIV2(K))
IF(MOD(K,2).EQ.0)C2NIN=-C2NIN
RETURN
END
C  DOUBLE PRECISION FUNCTION DC2NP(K)
C  PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C           FOR CE2N*(Z,Q).
C  DOUBLE PRECISION AB(25),BSJV1,BSJV2,BSYV2,DBSJV1,DBSJV2,
C  *          DBSYV2,V1,V2
C  COMMON/LOCAL/DUMMY1(4),V1,V2,AB
C  COMMON/RADIAL/BSJV1(25),BSJV2(25),BSYV2(25),DBSJV1(25),
C  *          DBSJV2(25),DBSYV2(25)
C  DC2NP=AB(K)*(-DBSJV1(K)*BSJV2(K)*V1+BSJV1(K)*DBSJV2(K)*V2)
C  IF(MOD(K,2).EQ.0)DC2NP=-DC2NP
C  RETURN
C  END
C  DOUBLE PRECISION FUNCTION DC2NIP(K)
C  PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C           FOR CE2N+1*(Z,Q).
C  DOUBLE PRECISION AB(25),BSJV1,BSJV2,BSYV2,DBSJV1,DBSJV2,
C  *          DBSYV2,V1,V2
C  COMMON/LOCAL/DUMMY1(4),V1,V2,AB
C  COMMON/RADIAL/BSJV1(25),BSJV2(25),BSYV2(25),DBSJV1(25),
C  *          DBSJV2(25),DBSYV2(25)
C  DC2NIP=AB(K)*(-DBSJV1(K)*BSJV2(K+1)*V1+
C  *          BSJV1(K)*DBSJV2(K+1)*V2-
C  *          DBSJV1(K+1)*BSJV2(K)*V1+BSJV1(K+1)*DBSJV2(K)*V2)
C  IF(MOD(K,2).EQ.0)DC2NIP=-DC2NIP
C  RETURN
C  END
C  DOUBLE PRECISION FUNCTION DC2NN(K)
C  PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C           FOR CE2N*(Z,Q).
C  DOUBLE PRECISION AB(25),BSIV1,BSIV2,BSKV2,DBSIV1,DBSIV2,
C  *          DBSKV2,V1,V2
C  COMMON/LOCAL/DUMMY1(4),V1,V2,AB
C  COMMON/RADIAL/BSIV1(25),BSIV2(25),BSKV2(25),DBSIV1(25),
C  *          DBSIV2(25),DBSKV2(25)
C  DC2NN=AB(K)*(-DBSIV1(K)*BSIV2(K)*V1+BSIV1(K)*DBSIV2(K)*V2)
C  IF(MOD(K,2).EQ.0)DC2NN=-DC2NN
C  RETURN
C  END
C  DOUBLE PRECISION FUNCTION DC2NIN(K)
C  PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C           FOR CE2N+1*(Z,-Q).
C  DOUBLE PRECISION AB(25),BSIV1,BSIV2,BSKV2,DBSIV1,DBSIV2,
C  *          DBSKV2,V1,V2
C  COMMON/LOCAL/DUMMY1(4),V1,V2,AB
C  COMMON/RADIAL/BSIV1(25),BSIV2(25),BSKV2(25),DBSIV1(25),
C  *          DBSIV2(25),DBSKV2(25)

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      DC2NIN=AB(K)*(-DBSIV1(K)*BSIV2(K+1)*V1+
      *          BSIV1(K)*DBSIV2(K+1)*V2-
      *          OBSIV1(K+1)*BSIV2(K)*V1+BSIV1(K+1)*OBSIV2(K)*V2)
      IF(MOD(K,2).EQ.0)DC2NIN=-DC2NIN
      RETURN
      END
      DOUBLE PRECISION FUNCTION SERAD(QD,R,DERIV,PS,AR,N)
      C      PURPOSE:  TO COMPUTE A MODIFIED MATHIEU FUNCTION
      C      (OR DERIVATIVE) OF FIRST KIND CORRESPONDING TO
      C      ODD MATHIEU FUNCTION ( SE FUNCTIONS)
      EXTERNAL S2N2P,S2N1P,S2N2N,S2N1N,DS2N2P,DS2N1P,DS2N2N,
      *          DS2N1N
      DOUBLE PRECISION AB(25),QD,PS,PSP,OUTPUT,AR(25)
      INTEGER R,CASE,DERIV
      COMMON/NTerm/N1
      COMMON/LOCAL/DUMMY(8),AB
      N1=N
      PSP=PS
      DO 5 I=1,N
      AB(I)=AR(I)
      C      5      CONTINUE
      CASE=MOD(R,2)+1
      IF(QD.LT.0.000) CASE=CASE+2
      IF(DERIV.EQ.1) GO TO 50
      GO TO(10,20,30,40),CASE
      C      THE VALUE OF SE2N+2(Z,Q)
      C      10      CALL SIGMA(S2N2P,OUTPUT)
      SERAD=-PS*OUTPUT/AB(1)
      RETURN
      C      THE VALUE OF SE2N+1(Z,Q)
      C      20      CALL SIGMA(S2N1P,OUTPUT)
      SERAD=PS*OUTPUT/AB(1)
      RETURN
      C      THE VALUE OF SE2N+2(Z,-Q)
      C      30      CALL SIGMA(S2N2N,OUTPUT)
      SERAD=PSP*OUTPUT/AB(1)
      RETURN
      C      THE VALUE OF S2EN+1(Z,-Q)
      C      40      CALL SIGMA(S2N1N,OUTPUT)
      SERAD=PSP*OUTPUT/AB(1)
      RETURN
      C      FOLLOWING ARE DERIVATIVES OF FUNCTIONS
      C      50      GO TO(60,70,80,90),CASE
      C      THE VALUE OF SE2N+2'(Z,Q)
      C      60      CALL SIGMA(DS2N2P,OUTPUT)
      SERAD=-PS*OUTPUT/AB(1)
      RETURN
      C      THE VALUE OF SE2N+1'(Z,Q)
      C      70      CALL SIGMA(DS2N1P,OUTPUT)
      SERAD=PS*OUTPUT/AB(1)
      RETURN

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C      THE VALUE OF SE2N+2*(Z,-Q)
80    CALL SIGMA(DS2N2N,OUTPUT)
      SERAD=PSP*OUTPUT/AB(1)
      RETURN
C      THE VALUE OF SE2N+1*(Z,-Q)
90    CALL SIGMA(DS2N1N,OUTPUT)
      SERAD=PSP*OUTPUT/AB(1)
      RETURN
      END
      DOUBLE PRECISION FUNCTION S2N2P(K)
      PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
                FOR SE2N+2(Z,Q).
      DOUBLE PRECISION AB(25),BSJV1,BSJV2,BSYV2,DBSJV1,DBSJV2,
*              DBSYV2
      COMMON/LOCAL/DUMMY1(8),AB
      COMMON/RADIAL/BSJV1(25),BSJV2(25),BSYV2(25),DBSJV1(25),
*              DBSJV2(25),DBSYV2(25)
      S2N2P=AB(K)*(BSJV1(K)*BSJV2(K+2)-BSJV1(K+2)*BSJV2(K))
      IF(MOD(K,2).EQ.0)S2N2P=-S2N2P
      RETURN
      END
      DOUBLE PRECISION FUNCTION S2N1P(K)
      PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
                FOR SE2N+1(Z,Q).
      DOUBLE PRECISION AB(25),BSJV1,BSJV2,BSYV2,DBSJV1,DBSJV2,
*              DBSYV2
      COMMON/LOCAL/DUMMY1(8),AB
      COMMON/RADIAL/BSJV1(25),BSJV2(25),BSYV2(25),DBSJV1(25),
*              DBSJV2(25),DBSYV2(25)
      S2N1P=AB(K)*(BSJV1(K)*BSJV2(K+1)-BSJV1(K+1)*BSJV2(K))
      IF(MOD(K,2).EQ.0)S2N1P=-S2N1P
      RETURN
      END
      DOUBLE PRECISION FUNCTION S2N2N(K)
      PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
                FOR SE2N(Z,-Q).
      DOUBLE PRECISION AB(25),BSIV1,BSIV2,BSKV2,DBSIV1,DBSIV2,
*              DBSKV2
      COMMON/LOCAL/DUMMY1(8),AB
      COMMON/RADIAL/BSIV1(25),BSIV2(25),BSKV2(25),DBSIV1(25),
*              DBSIV2(25),DBSKV2(25)
      S2N2N=AB(K)*(BSIV1(K)*BSIV2(K+2)-BSIV1(K+2)*BSIV2(K))
      IF(MOD(K,2).EQ.0)S2N2N=-S2N2N
      RETURN
      END
      DOUBLE PRECISION FUNCTION S2N1N(K)
      PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
                FOR SE2N+1(Z,-Q).
      DOUBLE PRECISION AB(25),BSIV1,BSIV2,BSKV2,DBSIV1,DBSIV2,
*              DBSKV2
      COMMON/LOCAL/DUMMY1(8),AB

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COMMON/RADIAL/BSIV1(25),BSIV2(25),BSKV2(25),DBSIV1(25),
*      DBSIV2(25),DBSKV2(25)
S2N1N=AB(K)*(BSIV1(K)*BSIV2(K+1)-BSIV1(K+1)*BSIV2(K))
IF(MOD(K,2).EQ.0)S2N1N=-S2N1N
RETURN
END
DOUBLE PRECISION FUNCTION DS2N2P(K)
C      PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C      FOR SE2N+2*(Z,Q).
DOUBLE PRECISION AB(25),BSJV1,BSJV2,BSYV2,DBSJV1,DBSJV2,
*      DBSYV2,V1,V2
COMMON/LOCAL/DUMMY1(4),V1,V2,AB
COMMON/RADIAL/BSJV1(25),BSJV2(25),BSYV2(25),DBSJV1(25),
*      DBSJV2(25),DBSYV2(25)
DS2N2P=AB(K)*(-DBSJV1(K)*BSJV2(K+2)*V1+
*      BSJV1(K)*DBSJV2(K+2)*V2+
*      DBSJV1(K+2)*BSJV2(K)*V1-BSJV1(K+2)*DBSJV2(K)*V2)
IF(MOD(K,2).EQ.0)DS2N2P=-DS2N2P
RETURN
END
DOUBLE PRECISION FUNCTION DS2N1P(K)
C      PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C      FOR SE2N+1*(Z,Q).
DOUBLE PRECISION AB(25),BSJV1,BSJV2,BSYV2,DBSJV1,DBSJV2,
*      DBSYV2,V1,V2
COMMON/LOCAL/DUMMY1(4),V1,V2,AB
COMMON/RADIAL/BSJV1(25),BSJV2(25),BSYV2(25),DBSJV1(25),
*      DBSJV2(25),DBSYV2(25)
DS2N1P=AB(K)*(-DBSJV1(K)*BSJV2(K+1)*V1+
*      BSJV1(K)*DBSJV2(K+1)*V2+
*      DBSJV1(K+1)*BSJV2(K)*V1-BSJV1(K+1)*DBSJV2(K)*V2)
IF(MOD(K,2).EQ.0)DS2N1P=-DS2N1P
RETURN
END
DOUBLE PRECISION FUNCTION DS2N2N(K)
C      PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C      FOR SE2N+2*(Z,-Q).
DOUBLE PRECISION AB(25),BSIV1,BSIV2,BSKV2,DBSIV1,DBSIV2,
*      DBSKV2,V1,V2
COMMON/LOCAL/DUMMY1(4),V1,V2,AB
COMMON/RADIAL/BSIV1(25),BSIV2(25),BSKV2(25),DBSIV1(25),
*      DBSIV2(25),DBSKV2(25)
DS2N2N=AB(K)*(-DBSIV1(K)*BSIV2(K+2)*V1+
*      BSIV1(K)*DBSIV2(K+2)*V2+
*      DBSIV1(K+2)*BSIV2(K)*V1-BSIV1(K+2)*DBSIV2(K)*V2)
IF(MOD(K,2).EQ.0)DS2N2N=-DS2N2N
RETURN
END
DOUBLE PRECISION FUNCTION DS2N1N(K)
C      PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C      FOR SE2N+1*(Z,-Q).

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      DOUBLE PRECISION AB(25),BSIV1,BSIV2,BSKV2,DBSIV1,DBSIV2,
      *          DBSKV2,V1,V2
      COMMON/LOCAL/DUMMY1(4),V1,V2,AB
      COMMON/RADIAL/BSIV1(25),BSIV2(25),BSKV2(25),DBSIV1(25),
      *          DBSIV2(25),DBSKV2(25)
      DS2N1N=AB(K)*(-DBSIV1(K)*BSIV2(K+1)*V1+
      *          BSIV1(K)*DBSIV2(K+1)*V2+
      *          DBSIV1(K+1)*BSIV2(K)*V1-BSIV1(K+1)*DBSIV2(K)*V2)
      IF(MOD(K,2).EQ.0)DS2N1N=-DS2N1N
      RETURN
      END
      DOUBLE PRECISION FUNCTION FERAD(QD,R,DERIV,PS,AR,N)
      PURPOSE:  TO COMPUTE A MODIFIED MATHIEU FUNCTION
      (OR DERIVATIVE) OF SECOND KIND CORRESPONDING TO
      C          EVEN MATHIEU FUNCTION ( CE FUNCTIONS)
      C
      C
      EXTERNAL FY2N,FY2N1,FK2N,FK2N1,DFY2N,DFY2N1,DFK2N,DFK2N1
      DOUBLE PRECISION AB(25),QD,PS,PSP,OUTPUT,AR(25),PI
      INTEGER R,CASE,DERIV
      COMMON/NTERM/N1
      COMMON/LOCAL/DUMMY(8),AB
      DATA PI/3.141592653589793D0/
      N1=N
      PSP=PS
      DO 5 I=1,N
      AB(I)=AR(I)
      C
      C
      5  CONTINUE
      CASE=MOD(R,2)+1
      IF(QD.LT.0.0D0) CASE=CASE+2
      IF(DERIV.EQ.1) GO TO 50
      GO TO(10,20,30,40),CASE
      C
      C
      10  THE VALUE OF FEY2N(Z,Q)
      CALL SIGMA(FY2N,OUTPUT)
      FERAD=PS*OUTPUT/AB(1)
      RETURN
      C
      C
      20  THE VALUE OF FEY2N+1(Z,Q)
      CALL SIGMA(FY2N1,OUTPUT)
      FERAD=PS*OUTPUT/AB(1)
      RETURN
      C
      C
      30  THE VALUE OF FEK2N(Z,-Q)
      CALL SIGMA(FK2N,OUTPUT)
      FERAD=PSP*OUTPUT/(AB(1)*PI)
      RETURN
      C
      C
      40  THE VALUE OF FEK2N+1(Z,-Q)
      CALL SIGMA(FK2N1,OUTPUT)
      FERAD=PSP*OUTPUT/(AB(1)*PI)
      RETURN
      C
      C
      FOLLOWING ARE DERIVATIVES OF FUNCTIONS
      50  GO TO(60,70,80,90),CASE
      C
      C
      60  THE VALUE OF FEY2N'(Z,Q)
      CALL SIGMA(DFY2N,OUTPUT)
      FERAD=PS*OUTPUT/AB(1)

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      RETURN
C     THE VALUE OF FEY2N+1*(Z,Q)
70    CALL SIGMA(DFY2N1,OUTPUT)
      FERAD=PS*OUTPUT/AB(1)
      RETURN
C     THE VALUE OF FEK2N*(Z,-Q)
80    CALL SIGMA(DFK2N,OUTPUT)
      FERAD=PSP*OUTPUT/(AB(1)*PI)
      RETURN
C     THE VALUE OF FEK2N+1*(Z,-Q)
90    CALL SIGMA(DFK2N1,OUTPUT)
      FERAD=PSP*OUTPUT/(AB(1)*PI)
      RETURN
      END
      DOUBLE PRECISION FUNCTION FY2N(K)
C     PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C               FOR FEY2N(Z,Q).
      DOUBLE PRECISION AB(25),BSJV1,BSJV2,BSYV2,DBSJV1,DBSJV2,
*                DBSYV2
      COMMON/LOCAL/DUMMY1(8),AB
      COMMON/RADIAL/BSJV1(25),BSJV2(25),BSYV2(25),DBSJV1(25),
*                DBSJV2(25),DBSYV2(25)
      FY2N=AB(K)*BSJV1(K)*BSYV2(K)
      IF(MOD(K,2).EQ.0)FY2N=-FY2N
      RETURN
      END
      DOUBLE PRECISION FUNCTION FY2N1(K)
C     PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C               FEY2N+1(Z,Q).
      DOUBLE PRECISION AB(25),BSJV1,BSJV2,BSYV2,DBSJV1,DBSJV2,
*                DBSYV2
      COMMON/LOCAL/DUMMY1(8),AB
      COMMON/RADIAL/BSJV1(25),BSJV2(25),BSYV2(25),DBSJV1(25),
*                DBSJV2(25),DBSYV2(25)
      FY2N1=AB(K)*((BSJV1(K)*BSYV2(K+1)+BSJV1(K+1)*BSYV2(K))
      IF(MOD(K,2).EQ.0)FY2N1=-FY2N1
      RETURN
      END
      DOUBLE PRECISION FUNCTION FK2N(K)
C     PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C               FOR FEK2N(Z,Q).
      DOUBLE PRECISION AB(25),BSIV1,BSIV2,BSKV2,DBSIV1,DBSIV2,
*                DBSKV2
      COMMON/LOCAL/DUMMY1(8),AB
      COMMON/RADIAL/BSIV1(25),BSIV2(25),BSKV2(25),DBSIV1(25),
*                DBSIV2(25),DBSKV2(25)
      FK2N=AB(K)*BSIV1(K)*BSKV2(K)
      RETURN
      END
      DOUBLE PRECISION FUNCTION FK2N1(K)
C     PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES

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```

C          FOR FEK2N+1(Z,Q).
DOUBLE PRECISION AB(25),BSIV1,BSIV2,BSKV2,DBSIV1,DBSIV2,
*          DBSKV2
COMMON/LOCAL/DUMMY1(8),AB
COMMON/RADIAL/BSIV1(25),BSIV2(25),BSKV2(25),DBSIV1(25),
*          DBSIV2(25),DBSKV2(25)
FK2N1=AB(K)*(BSIV1(K)*BSKV2(K+1)-BSIV1(K+1)*BSKV2(K))
RETURN
END
DOUBLE PRECISION FUNCTION DFY2N(K)
PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C          FOR FEY2N*(Z,Q).
C          DOUBLE PRECISION AB(25),BSJV1,BSJV2,BSYV2,DBSJV1,DBSJV2,
*          DBSYV2,V1,V2
COMMON/LOCAL/DUMMY1(4),V1,V2,AB
COMMON/RADIAL/BSJV1(25),BSJV2(25),BSYV2(25),DBSJV1(25),
*          DBSJV2(25),DBSYV2(25)
DFY2N=AB(K)*(-DBSJV1(K)*BSYV2(K)*V1+BSJV1(K)*DBSYV2(K)*V2)
IF(MOD(K,2).EQ.0)DFY2N=-DFY2N
RETURN
END
DOUBLE PRECISION FUNCTION DFY2N1(K)
PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C          FOR FEY2N+1*(Z,Q).
C          DOUBLE PRECISION AB(25),BSJV1,BSJV2,BSYV2,DBSJV1,DBSJV2,
*          DBSYV2,V1,V2
COMMON/LOCAL/DUMMY1(4),V1,V2,AB
COMMON/RADIAL/BSJV1(25),BSJV2(25),BSYV2(25),DBSJV1(25),
*          DBSJV2(25),DBSYV2(25)
DFY2N1=AB(K)*(-DBSJV1(K)*BSYV2(K+1)*V1+
*          BSJV1(K)*DBSYV2(K+1)*V2-
*          DBSJV1(K+1)*BSYV2(K)*V1+BSJV1(K+1)*DBSYV2(K)*V2)
IF(MOD(K,2).EQ.0)DFY2N1=-DFY2N1
RETURN
END
DOUBLE PRECISION FUNCTION DFK2N(K)
PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C          FOR FEK2N*(Z,Q).
C          DOUBLE PRECISION AB(25),BSIV1,BSIV2,BSKV2,DBSIV1,DBSIV2,
*          DBSKV2,V1,V2
COMMON/LOCAL/DUMMY1(4),V1,V2,AB
COMMON/RADIAL/BSIV1(25),BSIV2(25),BSKV2(25),DBSIV1(25),
*          DBSIV2(25),DBSKV2(25)
DFK2N=AB(K)*(-DBSIV1(K)*BSKV2(K)*V1+BSIV1(K)*DBSKV2(K)*V2)
RETURN
END
DOUBLE PRECISION FUNCTION DFK2N1(K)
PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C          FOR FEK2N+1*(Z,-Q).
C          DOUBLE PRECISION AB(25),BSIV1,BSIV2,BSKV2,DBSIV1,DBSIV2,
*          DBSKV2,V1,V2

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COMMON/LOCAL/DUMMY1(4),V1,V2,AB
COMMON/RADIAL/BSIV1(25),BSIV2(25),BSKV2(25),DBSIV1(25),
* DBSIV2(25),DBSKV2(25)
DFK2N1=AB(K)*(-DBSIV1(K)*BSKV2(K+1)*V1+
* BSIV1(K)*DBSKV2(K+1)*V2+
* DBSIV1(K+1)*BSKV2(K)*V1-BSIV1(K+1)*DBSKV2(K)*V2)
RETURN
END
DOUBLE PRECISION FUNCTION GERAD(QD,R,DERIV,PS,AR,N)
C PURPOSE: TO COMPUTE A MODIFIED MATHIEU FUNCTION
C (OR DERIVATIVE) OF SECOND KIND CORRESPONDING TO
C ODD MATHIEU FUNCTION ( CE FUNCTIONS)
EXTERNAL GY2N2,GY2N1,GK2N2,GK2N1,DGY2N2,DGY2N1,DGK2N2,DGK2N1
DOUBLE PRECISION AB(25),QD,PS,PSP,OUTPUT,AR(25),PI
INTEGER R,CASE,DERIV
COMMON/NTERM/N1
COMMON/LOCAL/DUMMY(8),AB
DATA PI/3.141592653589793D0/
N1=N
PSP=PS
DO 5 I=1,N
AB(I)=AR(I)
5 CONTINUE
CASE=MOD(R,2)+1
IF(QD.LT.0.0D0) CASE=CASE+2
IF(DERIV.EQ.1) GO TO 50
GO TO(10,20,30,40),CASE
C THE VALUE OF GEY2N+2(Z,Q)
10 CALL SIGMA(GY2N2,OUTPUT)
GERAD=-PS*OUTPUT/AB(1)
RETURN
C THE VALUE OF GEY2N+1(Z,Q)
20 CALL SIGMA(GY2N1,OUTPUT)
GERAD=PS*OUTPUT/AB(1)
RETURN
C THE VALUE OF GEK2N+2(Z,-Q)
30 CALL SIGMA(GK2N2,OUTPUT)
GERAD=PSP*OUTPUT/(AB(1)*PI)
RETURN
C THE VALUE OF GEK2N+1(Z,-Q)
40 CALL SIGMA(GK2N1,OUTPUT)
GERAD=PSP*OUTPUT/(AB(1)*PI)
RETURN
C FOLLOWING ARE DERIVATIVES OF FUNCTIONS
50 GO TO(60,70,80,90),CASE
C THE VALUE OF GEY2N+2'(Z,Q)
60 CALL SIGMA(DGY2N2,OUTPUT)
GERAD=-PS*OUTPUT/AB(1)
RETURN
C THE VALUE OF GEY2N+1'(Z,Q)
70 CALL SIGMA(DGY2N1,OUTPUT)

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GERAD=PS*OUTPUT/AB(1)
RETURN
C THE VALUE OF GEK2N+2(Z,-Q)
80 CALL SIGMA(DGK2N2,OUTPUT)
GERAD=PSP*OUTPUT/(AB(1)*PI)
RETURN
C THE VALUE OF GEK2N+1(Z,-Q)
90 CALL SIGMA(DGK2N1,OUTPUT)
GERAD=PSP*OUTPUT/(AB(1)*PI)
RETURN
END
DOUBLE PRECISION FUNCTION GY2N2(K)
PURPOSE: TO CALCULATE K TH TERM IN SUM OF SERIES
          FOR GEY2N+2(Z,Q).
C DOUBLE PRECISION AB(25),BSJV1,BSJV2,BSYV2,DBSJV1,DBSJV2,
C DBSYV2
* COMMON/LOCAL/DUMMY1(8),AB
COMMON/RADIAL/BSJV1(25),BSJV2(25),BSYV2(25),DBSJV1(25),
DBSJV2(25),DBSYV2(25)
* GY2N2=AB(K)*(BSJV1(K)*BSYV2(K+2)-BSJV1(K+2)*BSYV2(K))
IF(MOD(K,2).EQ.0)GY2N2=-GY2N2
RETURN
END
DOUBLE PRECISION FUNCTION GY2N1(K)
PURPOSE: TO CALCULATE K TH TERM IN SUM OF SERIES
          FOR GEY2N+1(Z,Q).
C DOUBLE PRECISION AB(25),BSJV1,BSJV2,BSYV2,DBSJV1,DBSJV2,
C DBSYV2
* COMMON/LOCAL/DUMMY1(8),AB
COMMON/RADIAL/BSJV1(25),BSJV2(25),BSYV2(25),DBSJV1(25),
DBSJV2(25),DBSYV2(25)
* GY2N1=AB(K)*(BSJV1(K)*BSYV2(K+1)-BSJV1(K+1)*BSYV2(K))
IF(MOD(K,2).EQ.0)GY2N1=-GY2N1
RETURN
END
DOUBLE PRECISION FUNCTION GK2N2(K)
PURPOSE: TO CALCULATE K TH TERM IN SUM OF SERIES
          FOR GEK2N+2(Z,Q).
C DOUBLE PRECISION AB(25),BSIV1,BSIV2,BSKV2,DBSIV1,DBSIV2,
C DBSKV2
* COMMON/LOCAL/DUMMY1(8),AB
COMMON/RADIAL/BSIV1(25),BSIV2(25),BSKV2(25),DBSIV1(25),
DBSIV2(25),DBSKV2(25)
* GK2N2=AB(K)*(BSIV1(K)*BSKV2(K+2)-BSIV1(K+2)*BSKV2(K))
RETURN
END
DOUBLE PRECISION FUNCTION GK2N1(K)
PURPOSE: TO CALCULATE K TH TERM IN SUM OF SERIES
          FOR GEK2N+1(Z,-Q).
C DOUBLE PRECISION AB(25),BSIV1,BSIV2,BSKV2,DBSIV1,DBSIV2,
C DBSKV2
*

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COMMON/LOCAL/DUMMY1(8),AB
COMMON/RADIAL/BSIV1(25),BSIV2(25),BSKV2(25),DBSIV1(25),
*      DBSIV2(25),DBSKV2(25)
GK2N1=AB(K)*(BSIV1(K)*BSKV2(K+1)+BSIV1(K+1)*BSKV2(K))
RETURN
END
DOUBLE PRECISION FUNCTION DGY2N2(K)
C  PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C  FOR GEY2N+2*(Z,Q).
DOUBLE PRECISION AB(25),BSJV1,BSJV2,BSYV2,DBSJV1,DBSJV2,
*      DBSYV2,V1,V2
COMMON/LOCAL/DUMMY1(4),V1,V2,AB
COMMON/RADIAL/BSJV1(25),BSJV2(25),BSYV2(25),DBSJV1(25),
*      DBSJV2(25),DBSYV2(25)
DGY2N2=AB(K)*(-DBSJV1(K)*BSYV2(K+2)*V1+
*      BSJV1(K)*DBSYV2(K+2)*V2+
*      DBSJV1(K+2)*BSYV2(K)*V1-BSJV1(K+2)*DBSYV2(K)*V2)
IF(MOD(K,2).EQ.0)DGY2N2=-DGY2N2
RETURN
END
DOUBLE PRECISION FUNCTION DGY2N1(K)
C  PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C  FOR GEY2N+1*(Z,Q).
DOUBLE PRECISION AB(25),BSJV1,BSJV2,BSYV2,DBSJV1,DBSJV2,
*      DBSYV2,V1,V2
COMMON/LOCAL/DUMMY1(4),V1,V2,AB
COMMON/RADIAL/BSJV1(25),BSJV2(25),BSYV2(25),DBSJV1(25),
*      DBSJV2(25),DBSYV2(25)
DGY2N1=AB(K)*(-DBSJV1(K)*BSYV2(K+1)*V1+
*      BSJV1(K)*DBSYV2(K+1)*V2+
*      DBSJV1(K+1)*BSYV2(K)*V1-BSJV1(K+1)*DBSYV2(K)*V2)
IF(MOD(K,2).EQ.0)DGY2N1=-DGY2N1
RETURN
END
DOUBLE PRECISION FUNCTION DGK2N2(K)
C  PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C  FOR GEK2N+2*(Z,-Q).
DOUBLE PRECISION AB(25),BSIV1,BSIV2,BSKV2,DBSIV1,DBSIV2,
*      DBSKV2,V1,V2
COMMON/LOCAL/DUMMY1(4),V1,V2,AB
COMMON/RADIAL/BSIV1(25),BSIV2(25),BSKV2(25),DBSIV1(25),
*      DBSIV2(25),DBSKV2(25)
DGK2N2=AB(K)*(-DBSIV1(K)*BSKV2(K+2)*V1+
*      BSIV1(K)*DBSKV2(K+2)*V2+
*      DBSIV1(K+2)*BSKV2(K)*V1-BSIV1(K+2)*DBSKV2(K)*V2)
RETURN
END
DOUBLE PRECISION FUNCTION DGK2N1(K)
C  PURPOSE:  TO CALCULATE K TH TERM IN SUM OF SERIES
C  FOR GEK2N+1*(Z,-Q).
DOUBLE PRECISION AB(25),BSIV1,BSIV2,BSKV2,DBSIV1,DBSIV2,

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*          DBSKV2,V1,V2
COMMON/LOCAL/DUMMY1(4),V1,V2,AB
COMMON/RADIAL/BSIV1(25),BSIV2(25),BSKV2(25),DBSIV1(25),
*          DBSIV2(25),DBSKV2(25)
DGK2N1=AB(K)*(-DBSIV1(K)*BSKV2(K+1)*V1+
*          BSIV1(K)*DBSKV2(K+1)*V2-
*          DBSIV1(K+1)*BSKV2(K)*V1+BSIV1(K+1)*DBSKV2(K)*V2)
RETURN
END
```

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## ANISOTROPIC ELLIPTIC OPTICAL FIBERS

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The exact characteristic equation for an anisotropic elliptic optical fiber is obtained for odd and even hybrid modes in terms of infinite determinants utilizing Mathieu and modified Mathieu functions. A simplified characteristic equation is obtained by applying the weakly guiding approximation such that the difference in the refractive indices of the core and the cladding is small.

The simplified characteristic equation is used to compute the normalized guide wavelength for an elliptical fiber. When the anisotropic parameter is equal to unity, the results are compared with the previous research and they are in close agreement.

For a fixed value of normalized cross-section area or major axis, the normalized guide wavelength  $\lambda/\lambda_0$  for an anisotropic elliptic fiber is small for larger the value of anisotropy. This condition indicates that more energy is carried inside of the fiber. However, the geometry and anisotropy of the fiber have a smaller effect when the normalized cross-section area is very small or very large.